Cosmic rays and air shower physics I

Ralph Engel *Karlsruhe Institute of Technology (KIT)*

(photograph by S. Saffi)

Discovery of cosmic rays (Pacini, Hess, Kolhörster)

ViktorF.Hess(Wien), Über Beobachtungen der durchdringenden Strahlung bei sieben Freiballonfahrten.

Im Vorjahre habe ich bereits Gelegenheit gehabt, zwei Ballonfahrten zur Erforschung der durchdringenden Strahlung zu unternehmen; über die erste Fahrt wurde schon auf der Naturforscherversammlung in Karlsruhe von mir berichtet¹). Bei beiden Fahrten ergab sich keine wesentliche Änderung der Strahlung gegenüber der am Erdboden beobach-

Nobel Prize 1936

(1911 – 1913)

Cosmic rays as new energy frontier Marketing cosmic rays at dividirontie high altitudes and the set

REVIEWS OF

(but merit stolen) *Robert Millikan*

Volume IX

Arthur Compton

The birth of high-energy particle physics Early Early Evidence of Museum and The Company of Museum and Museum and Museum and Museum and Museum and Museu
And Museum and Museum
A

1936: discovery of muon (initially mistaken as pion) by Anderson and Neddermeyer Discovery of the pion (1947) 1947: discovery of pion by Lattes, Murihead, Occhialini Powell (and "Mrs. I. Roberts"?) and Powell pi -> mu + number -> n 1 1 11 11 11 11 11 1 111 1

Fig. 1. OBSERVATION BY MRS. I. ROBERTS. PHOTOMICROGRAPH WITH COOKE × 45 'FLUORITE' OBJECTIVE. ILFORD 'NUCLEAR RESEARCH', BORON-LOADED C2 EMULSION. m_1 IS THE PRIMARY AND m_2 THE SECONDARY MESON. THE ARROWS, IN THIS AND THE FOLLOWING THE PHOTOGRAPHS ARE COMPLETELY UNRETOUCHED

Cloud chamber observations of cosmic rays 1932: discovery of positron by Anderson

1911: cloud chamber by Wilson

1947: discovery of kaon, the first strange particle, by Rochester & Butler

1951: discovery of Λ, the first strange baryon, by Armenteros et al.

Extensive cosmic ray showers (Auger et al. 1936)

Primary particle energies exceeding 10^6 GeV = 10^{15} eV

Fig. 12-1 Air-shower experiments by Auger and his collaborators in 1938 were made with the counter arrangement shown here. Counters G_1 and G_2 were placed one above the other, with their axes 22 cm apart. The third counter G_3 was moved horizontally to various distances d ranging from 15 centimeters to 75 meters.

Fig. 12-2 Air-shower data obtained by Auger with the counter arrangement shown in Fig. $12-1$. The horizontal scale gives the horizontal distance d between counter G_3 and the pair of counters G_1 and G_2 ; the vertical scale, the number of coincidences per hour. (From a paper in Le Journal de Physique et Le Radium, vol. 10, p. 39, 1939.)

Cascade in a cloud chamber at 3027 m altitude (~10 GeV)

Particle accelerator development

Cyclotron (1931) Ernest Orlando Lawrence (Nobel prize 1939)

Acceleration by 2kV max. energy 80 keV

1932: 27-inch cyclotron: 3 MeV protons

1939: 60-inch cyclotron: 20 MeV protons

1952: Cosmotron: 3 GeV (Brookhaven) **1954:** Bevatron: 6 GeV (Berkeley) **1957:** Synchrophasotron: 10 GeV (Dubna) **1960:** Brookhaven AGS: 33 GeV **1967:** U-70: 70 GeV (Serpukhov)

Cosmic rays of ultra-high energy – 1020 eV

(Received 10 January 1963)

Energy conservation: 1.6 GeV / charged particle, overall estimate robust

Cascade of secondary particles: extensive air shower

AGASA event, *E* **~1.7 – 2.6 1020 eV** *(AGASA, PRL 73 (1994) 3491)*

OMG event, *E* **~3.2 1020 eV** *(Fly's Eye, ApJ 441 (1995) 144)*

Amaterasu event, *E* **~2.4 1020 eV** *(TA, Science 382 (2023) 903)*

Ultra-high-energy cosmic rays – 1020 eV

 Need accelerator of size of the orbit of the planet Mercury to reach 10²⁰ eV with LHC technology

Large Hadron Collider (LHC), 27 km circumference, superconducting magnets

System.	$E_{\text{lab}} \approx \frac{E_{\text{c.m.}}^2}{2m_p}$
---------	---

\blacksquare Synchrotron onorgy loss the $2m_p$ physics (with units of c ω , $\frac{1}{2}$. The energy ross of c $\frac{1}{2}$. The energy momentum 4-vector of beam 1-vector of beam 1-vec **Synchrotron energy loss**

$$
U_0 = \frac{e^2}{3\epsilon_0}\,\frac{\beta^3\gamma^4}{\rho}\propto \frac{1}{\rho}\,\frac{E^4}{m^4}
$$

Fixed target: Forward direction (beam fragmentation) covered by detectors particle energy was direction moneturements very shallonging (if not impossible) comments begin an ection incasarements very shanenging (in not impossible).
19.

Fixed target vs. collider setup is including the voltation to his equation to the acceleration to the second to the second to the second to th The former of the final and continued acceleration acceleration and the first acceleration and a beam of particle is the end of the acceleration to the second to the second to the second to the second t is the collider, in which two beams of high energy particles are brought into the collider are brought into th
The collider are brought into the collider are brought into the collider are brought into the collider are bro ad 1 $\overline{}$ n.
1 $\overline{\mathbf{G}}$

Flux of cosmic rays

40 km \sim 30 - 40 km \blacksquare 30

Composition

- 85% H nuclei (protons)
- 12% He nuclei
- 1% heavier nuclei
- 2% electrons and positions

0.01 - 0.001% antiprotons

Flux of cosmic rays

Composition

- 85% H nuclei (protons)
- 12% He nuclei
- 1% heavier nuclei
- 2% electrons and positions

0.01 - 0.001% antiprotons

11

Flux approximately power law

$$
\frac{dN}{dEd\Omega dAdt} \propto E^{-\gamma}
$$

Energy spectrum of all-particle flux (energy per particle)

$$
\gamma \approx 2.7
$$
 $10^{11} \text{eV} < E < 10^{15.5} \text{eV}$
\n ≈ 3.1 $10^{15.5} \text{eV} < E < 10^{18.5} \text{eV}$

Power laws are common in nature (i)

Department of Physics and Center for the Study of Complex Systems, University of Michigan, MI 48109. U.S.A.

Power laws, Pareto distributions and Zipf's law

M. E. J. Newman

When the probability of measuring a particular value of some quantity varies inversely as a power of that value, the quantity is said to follow a power law, also known variously as Zipf's law or the Pareto distribution. Power laws appear widely in physics, biology, earth and planetary sciences, economics and finance, computer science, demography and the social sciences. For instance, the distributions of the sizes of cities, earthquakes, solar flares, moon craters, wars and people's personal fortunes all appear to follow power laws. The origin of power-law behaviour has been a topic of debate in the scientific community for more than a century. Here we review some of the empirical evidence for the existence of power-law forms and the theories proposed to explain them.

> *(M. Newman cond-mat/0412004)* from an even for α for α β

Power laws are common in nature (ii) 102 103 104 105 $\overline{\mathbf{1}}$ 1

FIG. 4 Cumulative distributions or "rank/frequency plots" of twelve quantities reputed to follow power laws. The distributions were computed as described in Appendix A. Data in the shaded regions were excluded from the calculations of the exponents in Table I. Source references for the data are given in the text. (a) Numbers of occurrences of words in the novel Moby Dick by Hermann Melville. (b) Numbers of citations to scientific papers published in 1981, from time of publication until June 1997. (c) Numbers of hits on web sites by 60 000 users of the America Online Internet service for the day of 1 December 1997. (d) Numbers of copies of bestselling books sold in the US between 1895 and 1965. (e) Number of calls received by AT&T telephone customers in the US for a single day. (f) Magnitude of earthquakes in California between January 1910 and May 1992. Magnitude is proportional to the logarithm of the maximum amplitude of the earthquake, and hence the distribution obeys a power law even though the horizontal axis is linear. (g) Diameter of craters on the moon. Vertical axis is measured per square kilometre. (h) Peak gamma-ray intensity of solar flares in counts per second, measured from Earth orbit between February 1980 and November 1989. (i) Intensity of wars from 1816 to 1980, measured as battle deaths per 10 000 of the population of the participating countries. (j) Aggregate net worth in dollars of the richest individuals in the US in October 2003. (k) Frequency of occurrence of family names in the US in the year 1990. (l) Populations of US cities in the year 2000.

Power laws are common in nature (ii) 0.01 0.1 1 102 103 104 105 $\overline{\mathbf{1}}$ 1

shed in 1981, from time of publication until June ne Internet service for the day of 1 December 1997. the earthquake, and hence the distribution obeys a ond, measured from Earth orbit between February as battle deaths per 10 000 of the population of the lividuals in the US in October 2003. (k) Frequency of US cities in the year 2000.

Cosmic ray flux and interaction energies

Outline of lectures

- Cosmic rays below the knee direct measurements
- Physics of extensive air showers
- Discussion and exercises *(topics to be decided)*
-

- Cosmic rays of very high energy – indirect measurements

Flight altitude 30-35 km

Long-duration balloon flights

Geomagnetic cutoff and East-West effect

Earth's magnetic field

Radius of curvature smaller than radius of Earth

 $R =$ *p Ze* \approx *E Ze*

 $R_L = 3 \times 10^3$

Vicinity of poles: $B \approx 60 \mu T$ Equator: $B \approx 30 \mu T$

 $q = Ze$ Lorentz force $F_L = qvB$ Inertial force **Charge** $R_L =$ *p ZeB* Lamor radius $F_F = m$ v^2 *RL* $=$ ν *p RL*

$$
\left(\frac{E}{\text{GeV}}\right) \left(\frac{\mu \text{T}}{ZB}\right) \text{ km}
$$

Rigidity

Solar modulation of cosmic flux

Solar modulation – observations UIT UNSCHVALI

Galaxy and galactic magnetic fields

Magnetic field not well known,

 $B = 3 \mu G = 30 \text{ nT}$ close to Solar System

([astronomy.stackexchange.com\)](http://astronomy.stackexchange.com)

$$
R_L \simeq 1 \,\text{pc} \times \left(\frac{E}{10^{15}\,\text{eV}}\right) \left(\frac{\mu\text{G}}{ZB}\right)
$$

Coherent and turbulent magnetic fields (of similar strength)

Galaxy and galactic magnetic fields

Magnetic field not well known, $B = 3 \mu G = 30 \text{ nT}$ close to Solar System

Coherent and turbulent magnetic fields $\left(\text{of similar strength}\right)$ and $\text{B} = 3 \mu\text{G} = 1$ orbit with a gyroradius rg = P /Bc. (**b**) When the field is non-uniform the particle drifts away from a field **line due to due to the gradient and curbulent magnetic fields. Coherent and turbulent magnetic field that has a meeting field that has a scale meet field that has a scale that has a scale that has a scale that has a scale** \mathbf{u} is similar such \mathbf{g} uri \mathbf{u} , all pass through it with \mathbf{u} with \mathbf{u} without being a finite much.

 $1 pc = 3.26 ly = 3.08 10¹⁶ m$

Diffusion: random walk, distance scales ~ (time)2

Non-Extragalactic Sources

Power needed to maintain cosmic ray flux

Assumption: entire galaxy homogeneously filled with cosmic rays

Density of particles for given flux

$$
\frac{dN}{dEdV} = \frac{4\pi}{c} \frac{dN}{dEd\Omega d}
$$

Isotropy
$$
\int d\Omega = 4\pi
$$

$$
P_{\rm SNR} \approx 10^{42} \text{erg}
$$

Kinetic energy released in SN expl

$$
E_{\text{tot}} = \int dV \int dE \ E \cdot \frac{dN}{dEdV}
$$

Mean escape time $\tau_{\rm esc} \approx 10^7 a$

Total cosmic ray energy

Required power of cosmic ray sources
$$
P_{src} = E_{tot}/\tau_{esc} \approx 10^{41} \text{erg/s}
$$
 (1 erg = 0.1 µJ)

ASTRONOMY: COSMIC RAYS from supernova remnants of super-novae are available, and it has not thus far been possible to inter-not thus far been possible to inte
- and it has not thus far been possible to inter-not the inter-not thus far been possible to inter-not the int

COSMIC RAYS FROM SUPER-NOVAE

20 pc Distance ~ 2.2 kpc

Observed galactic SN explosions:

~3 SN explosions / 100 yrs Kinetic energy of ejecta: \sim 10⁵¹ erg

1604 (Kepler) 1572 (Tycho) 1181 (Chinese astronomers) 1054 (Crab nebula) 1006 (Chinese and Arabian records)

Estimates:

If interest in these questions still prevails at that future time, science will therefore be able to test the correctness of our hypothesis some time during the next thousand years or so, as the occurrence of a super-nova in our own system would multiply the intensity of the cosmic rays by a factor one thousand or more. It also seems quite possible to observe with cosmic-ray electroscopes the flare-up of a super-nova in one of the nearer extragalactic nebulae, as for them $r = 1000 n$, and (1) the cosmic representation of the time. The cosmic rays is the cosmic rays of the cosmic rays in the time. \mathcal{L} factor indicates that the origin of the origin of the sound be sound before rays can be sound in the sound the sun normal control in any of the sun of th

MOUNT WILSON OBSERVATORY, CARNEGIE INSTITUTION OF WASHINGTON AND CALI-FORNIA INSTITUTE OF TECHNOLOGY, PASADENA

since, according to the preceding paper, 1047 ergs. See also seems a probable paper, 1934 ergs. See also seems a probable paper. The probable paper of the probable paper. The probable paper of the probable paper. The proba

value of LT. Suppose of LT. Suppose of LT. Suppose of LT. Suppose of the neighborhood \mathcal{A} A. Introduction.—Two important facts support the view that cosmic
normalize the system of our model of the view that we disposed the we we we electrostatic potential with respect to interstellar space. rays are of extragalactic origin, if, for the moment, we disregard the possibility that the earth may possess a very high and self-renewing

$$
\Delta \sigma = 0.01/n^2 \text{ ergs/cm.}^2 \text{ sec.}, \qquad (10)
$$

where n is a number of the order one. It might in this connection be of interest to follow up the causes for Regener's⁴ curious balloon observation of March 29, 1933. \blacksquare charged particles heavier than electrons have been observed in any con- α is a number of the order one. It might in this connection

By W. BAADE AND F. ZWICKY

field to produce the observed dip in intensity at the equator.

$$
P_{SNR} \approx 10^{42} \text{erg/s}
$$

Kinetic energy released in SN explosions

Diffusive shock (Fermi) acceleration Diffusive shock (Fermi) acceleration

13-Dec. THE UNIVERSITY OF heary **s** a) since against in extitution against exament **second order acceleration**
 Ferministic function of the second of the Non reta lic case MV $=$ morri **CLCH** eire $40V$ **TELES** $+2$ $2r$ \mathbf{m} wes euer $\mathcal{L} = \mathcal{L}$ **1 1 1 1 1 1 1** gain order werase PHYSICAL REVIEW MV^2 order FU 2 W ł۳ FU h i z

Fermi's original work:

Particles scatter on moving magnetic clouds

VOLUME 75, NUMBER 8

APRIL 15, 1949

On the Origin of the Cosmic Radiation

ENRICO FERMI Institute for Nuclear Studies, University of Chicago, Chicago, Illinois (Received January 3, 1949)

A theory of the origin of cosmic radiation is proposed according to which cosmic rays are originated and accelerated primarily in the interstellar space of the galaxy by collisions against moving magmetic fields. One of the features of the theory is that it yields naturally an inverse power law for the spectral distribution of the cosmic rays. The chief difficulty is that it fails to explain in a straightforward way the heavy nuclei observed in the primary radiation.

Stochastic acceleration on SN shock fronts

First order Fermi acceleration

Assumption:

particles scatter elastically on turbulent mag. fields

ent relative energy gain

all angles

$$
\Delta E = \frac{1}{2}m(v + (u_1 - u_2))^2 - \frac{1}{2}mv^2
$$

vertical crossing, non-relativistic shock speed

$$
\frac{\Delta E}{E} = \frac{4}{3} \frac{(u_1 - u_2)}{c}
$$
Energy-independent
Factor from averaging over a

$$
\frac{\Delta E}{E} \approx 2 \frac{(u_1 - u_2)}{v}
$$

Expected energy distribution

Assumption: energy-independent escape probability *Pesc* per cycle

Energy gain per complete cycle of crossings

Energy after *k* cycles

Number of particles available for further acceleration

$$
N = N_0 \ (1 - P_{\rm esc})^k
$$

Flux of particles

 $N(>E) =$

Numerical values depend on many details

Ideal diffusive shock acceleration yields dN/dE ~ *E–2*

$$
\frac{\Delta E}{E} = \xi
$$

$$
E=E_0\ \mathbf{\mathcal{E}}^k
$$

$$
α = -\ln(1-Pesc)/\ln\xi
$$

α ≈ 1 (see Longair's textbook)

Comparison to particle acceleration in heliosphere 106 r Solar /^iu, ; lfind / ^#^ $\overline{}$ second r Separa *'I* 1 in haliaenhara 30 and MeV/nucleon measured SIS. this case 27-day average fluxes were integrated. MeV/nucleon time-intensity plot high-energy events o barticle acceleration is [10]; Figure 3 illustrates how these events contribute $\overline{}$ **Integrated Oxygen Fluences**

cine
Cine
cine
cine aa
Q

4J

Maximum particle energy

Astrophysical plasmas: no static electric fields (charges are free to travel)

Faraday's law: changing magnetic fields produce electric field

Hillas criterion

Lorentz transformation

$$
E_{p,\max} = qL\frac{u}{c}B = \beta_s qLB
$$

(after S. Gabici, KSETA 2024)

Maximum particle energy

Astrophysical plasmas: no static electric fields (charges are free to travel)

Faraday's law: changing magnetic fields produce electric field

Hillas criterion

Lorentz transformation

$$
E_{p,\max} = qL\frac{u}{c}B = \beta_s qLB
$$

 $\overline{\overline{u}}$ \vec{u}

(after S. Gabici, KSETA 2024)

Very general result, broad application, further restrictions apply (energy losses)

(Hillas, Ann. Rev. Astron. Astrophys. 1984)

Leaky Box model **MANAI** Zusammenfassung

proportional to number of particles in box

Sources uniformly distributed

Effect of cosmic ray confinement in galaxy

Simplification: only one particle type considered, no energy losses

Assumption: equilibrium reached, flux independent of time

$$
0 = -\frac{1}{\tau_{\rm esc}} N(E) + Q(E)
$$

$$
\frac{dN(E)}{dt} = -\frac{1}{\tau_{\rm esc}}N(E) + Q(E)
$$

Cross check of model with secondary elements

Interstellar medium in galaxy: \sim 1 atom /cm³

Spallation of nuclei

Re-write in terms of grammage and add interactions

$$
Q_{\rm C \rightarrow B}(E) = f_{\rm C \rightarrow B} \times \frac{1}{\lambda_{\rm C,int}} N_{\rm C}(E)
$$

$$
0 = -\frac{1}{\lambda_{B,esc}(E)}N_B - \frac{1}{\lambda_{B,int}}N_B + Q_0
$$

Caution: here always energy per nucleon used

$$
\frac{N_{\rm B}}{N_{\rm C}} = \frac{f_{\rm C \to B}}{\lambda_{\rm C,int}} \frac{\lambda_{\rm B, esc}(E)}{1 + \lambda_{\rm B, esc}(E)/\lambda_{\rm B,int}}
$$

 \mathcal{C} \rightarrow B (E)

 $\lambda_{\rm esc}(E) = \beta c \rho_{\rm ISM} \tau_{\rm esc}(E)$

New detectors – CALET and DAMPE

CALET – Calorimetric Electron Telescope DAMPE – Dark Matter Particle Explorer

Aperture \sim 0.1 m² sr

Cross check of model with secondary elements

(DAMPE, Science Bulletin 67 (2022) 2162) and the shaded bands representing the statistical uncertainties are the sum in 43

$$
\frac{N_{\rm B}}{N_{\rm C}} = \frac{f_{\rm C \to B}}{\lambda_{\rm C,int}} \frac{\lambda_{\rm B, esc}(E)}{1 + \lambda_{\rm B,esc}(E)/\lambda_{\rm B,int}}
$$

Cosmic ray clocks – energy-dependent escape time 10Be measurements (129) up to 2 GeV (and hence longer decay lifetime) and hence longer decay lifetime (129) up to 2 GeV (and hence longer decay lifetime) and hence longer decay lifetime (and hence longer decay lifetime) an with the fit to the fit to the fit to the other data, although the statistics are not very constraining. The s
The statistics are not very constraint in the statistics are not very constraint in the statistics are not very ks – energy-dependent escape time the analysis (stat. only error). In the right panel final merged result with complete error). In the right pan
In the right panel final merged result with complete error evaluation. In the right panel final merged result

Required by observations model interpretation, see Reference 132. Luckily, the leaky-box-model surviving frac $t = \frac{1}{2}$ for a beginning meaningful $\frac{1}{2}$ for a given model.

Only rigidity, i.e. to derive the the halo size. The halo size size size size sense important and at 3 √ 5 × 1028 cm2 s−1 (at 3 μ diffusion coefficiently with stability, i.e. the combined with stability, i.e. $\alpha \sim \sqrt{\frac{1}{\sigma}}$

 $\tau_{\rm esc}$ \propto *E Z* ⇥0*.*⁷ For example, in a simple diffusive halo model, the surviving fraction determines the *zh = 4 kpc. We can then compare then compare then compare then compare then compare time of ≈107 yr with the lea
In the leaky-box model then compare then compare then compare then compare then compare then compare then c*

 2×10^7 vr **Figure 6:** Finergy dependence not confirmed by new AMS data **Energy dependence not confirmed by new AMS data**

wese • *Cose* • *Cose* • *Cose* • *Cose Cose Cose Cose Cose Cose Cose*

experiments.

Non-linear diffusive shock acceleration

$$
q = \frac{3r}{r - 1 - (v_A/u_2)}
$$
Correction t
larger slope:

$$
f(p) \sim p^{-q}
$$
; $q \approx \frac{3r}{r-1}$

Unexpected structure of energy spectrum *2.2. The proton spectrum*

⁽P. Mertsch, DPG 2023)

The DAMPE detector can measure the most abundant CR nuclei up to an energy of Escape time: diffusion coefficient

Source

$$
N(E) = \tau_{\rm esc}(E) Q(E)
$$

Spectral hardening first found by PAMELA

Unexpected structure of energy spectrum *2.2. The proton spectrum*

diffusion coefficient

Source

(P. Mertsch, DPG 2023)

Source effect

• Break in source spectrum Stanev, Biermann and Gaisser (1993) Parizot (2004) Ptuskin, Zirakashvili and Seo (2013)

Transport effect

• Break in diffusion coefficient

uncertainty in the second
The second s $N(E) = \tau_{\rm esc}(E) Q(E)$

Unexpected structure – check with secondaries *2.2. The proton spectrum* $\mathbf{1}_{\mathbf{1}_{\mathbf{1}}}$. So the results are presented in h secondaries with the spectral indices of \sim and oxygen in the magnitude and the magnitude and the magnitude and the rigidity of the rigidity of the rigidi **-1 sr -1 15 5.4**

value above 7 GV for $A^{\text{L}}=A^{\text{L}}$ and above 30 GV for Be=B. $10\qquad 20\quad 30$ 10^2 2×10^2 10^3 2×10^3 and Best 20.36 \overline{R} **Rigidity** \overline{R} **[GV]** these fits we note that the Lie and th
Be ratio is 2.0 \times 0.1 \times — स्टब्स्ट का स्टब्स्ट का अन्य का अन्य
दिखा प्रथम संस्कृति का अन्य **30 ² 10 ² 2**×**10 ³ 10 ³ 2**×**10] 1.7 (GV) -1 sr -1 s-2 [m 2.7 ~ R** × **Flux 0 1 2 3 4 Helium** $\sum_{\mathbf{C} \subseteq \mathcal{C}}$ addition, above ∞200 $\sum_{\mathbf{C} \subseteq \mathcal{C}}$ and B all $\sum_{\mathbf{C} \subseteq \mathcal{C}}$ and B all $\sum_{\mathbf{C} \subseteq \mathcal{C}}$ $\sum_{i,j}$ a $\left[\begin{array}{ccc} 1 & 0 \end{array}\right]$ **Oxygen** \times 28 **Lithium**×**200** and Tables II and III of Ref. [14], and are reported in **Beryllium×400**Table Supplemental Material Tables T

Tables Table **Boron**×**145** $F = \frac{1}{2} \int_{0}^{2\pi} \frac{1}{\sqrt{2}} \, e^{i \omega t} \, e^{-i \omega t} \,$ with the $\begin{bmatrix} 1 & \cdots & \cdots & 1 \end{bmatrix}$ and $\begin{bmatrix} 1 & \cdots & \cdots & 1 \end{bmatrix}$ multiplied by R $\begin{bmatrix} 1 & \cdots & \cdots & 1 \end{bmatrix}$ $\begin{bmatrix} \mathbb{A} & \mathbb{A}^T &$ display purposes only, the C, O, Li, Be, and B fluxes were $\begin{matrix} \times \\ \times \end{matrix}$ indicated. For clarity, the He, O, Li, and B data points $\begin{matrix} \wedge \\ \wedge \end{matrix}$ $\frac{a}{2}$ and the three displaced horizontal $\frac{a}{2}$ $\sum_{i=1}^{\infty}$ is the angle and Δ in the above Δ $\frac{30}{20}$ GV, as do the three primary fluxes above $\frac{1}{20}$ $\frac{d}{d}$ primary cosmic rays fluxes and of secondary cosmic rays fluxes and of secondary cosmic rays $\frac{1}{2}$ cosmic rays fluxes are distinctly different. indices are nearly identical, but distinctly different from the three parameters \mathbf{r} $\begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$ dependence of $\begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix}$ **AQQQQQQQQQQQQQQQ**d $\overline{\mathcal{P}}$ of $\overline{\mathcal{P}}$ in details in deta $r_{\rm E}$, and both $r_{\rm E}$ the littlium, and boron fluxes to the little state to the little state of the little state o $\begin{bmatrix} -1 & -1 & -1 \\ 0 & -1 & 0 \end{bmatrix}$ α Γ and \overline{A} \overline{A} \overline{B} \overline{C} and \overline{C} and \overline{C} and \overline{C} statistical and systematic errors. The detailed variation \mathbb{R} and \mathbb{R} and \mathbb{R} are detailed variations. The detailed variation of \mathbb{R} and \mathbb{R} are detailed variations. The detailed variation of \mathbb $\begin{array}{@{}lllllll@{}} 0 & \text{if} & \text{if}$ **10 20 30 10**² **2×10**² **10**³ **2×10**³ \succ **Spectral Index -3 -2.5 Lithium • Beryllium Boron Helium Carbon Oxygen**

 $\overline{}$

 \times **10³**

Be/C 1.8 *(AMS, PRL 120 (2018) 021101)*

- Can be distinguished by secondaries Vladimirov et al. (2012) \mathbf{E} , \mathbf{S} and III of the Supplemental Material \mathbf{S} as a set of the Supplemental Material \mathbf{S} as a set of the Supplemental Material \mathbf{S}
- break in source spectrum: break in secondaries similar $\frac{1}{2}$ $\frac{$

Unexpected structure – check with secondaries *2.2. The proton spectrum* $\mathbf{1}_{\mathbf{1}_{\mathbf{1}}}$. So the results are presented in h secondaries with the spectral indices of \sim and oxygen in the magnitude and the magnitude and the magnitude and the rigidity of the rigidity of the rigidi **-1 sr -1 15 5.4**

ratio from our previous publication [24] which is consistent with

Be/C 1.8 *(AMS, PRL 120 (2018) 021101)*

Electron and positron spectra

Backup slides

Neutral (secondary) particles as messengers

and similar interactions of nuclei, as well as dissociation of nuclei

Two source classes

- **- Interaction in dense source regions with photon field or gas**
- **- Interaction during propagation with photons of CMB and other backgrounds**