Asymptotic safety and the road towards running relational observables

Renata Ferrero

based on work in collaboration with Kevin Falls and Martin Reuter

February 22nd 2024

11th TUX Workshop in Quantum Gravity



Plan of the talk

Plan of the talk

Asymptotic Safety

Effective Average Action
Functional Renormalization Group
Quantum Einstein Gravity

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Asymptotic Safety

Effective Average Action
Functional Renormalization Group
Quantum Einstein Gravity

Running relational observables

Composite operators

Covariant formulation

Flow of the observables

Effective Average Action & Functional Renormalization Group

Asymptotic Safety

The RG

RG in perturbation theory

only the finitely many beta functions that are related to the relevant couplings are considered

Callan (1970), Symanzik (1970)

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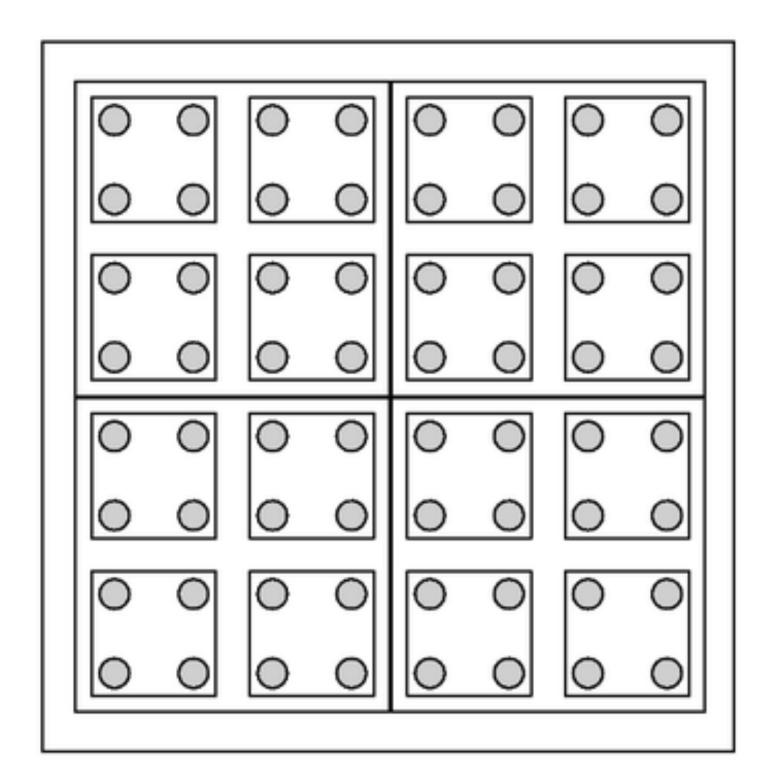
The RG

Exact RGE

Wilsonian Exact RG

the quantum fluctuations in the path integral can be integrated out progressively

Wilson (1971), Kadanoff (1966)



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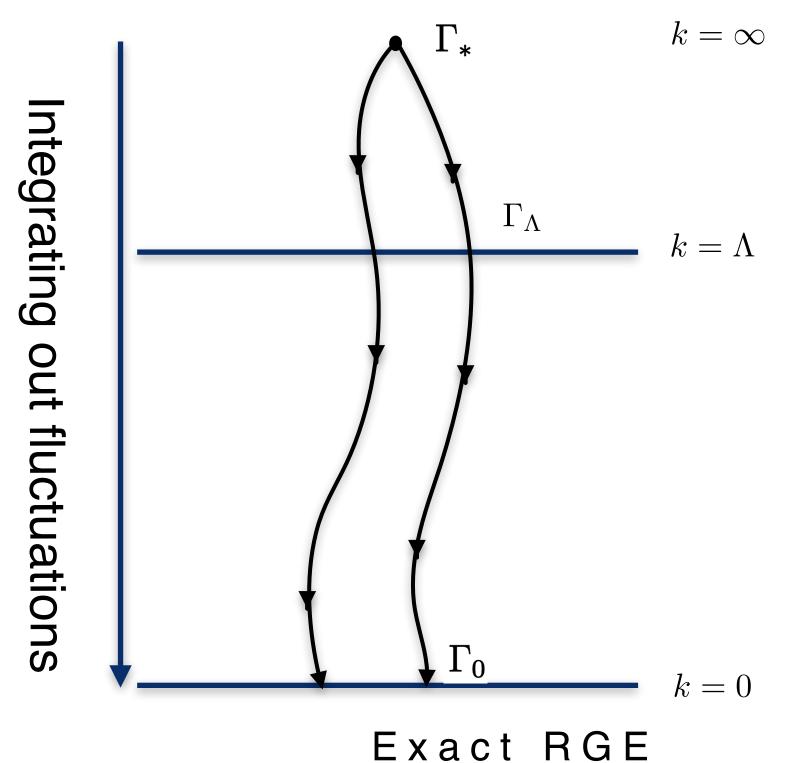
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UV fixed point Microscopic

model: Asymptotic Safety, Dynamical Triangulations

Effective action

Initial condition



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Integrating out fluctuations

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Functional RG

Functional RGE

scale-dependent version of the effective action, the Effective Average Action

> Wetterich (1991) Reuter and Wetterich (1994)

Alternative manipulation of the path integral

• Implement the underlying RG idea already at the level of the EAA, fully independently of the bare action.

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Generating functional

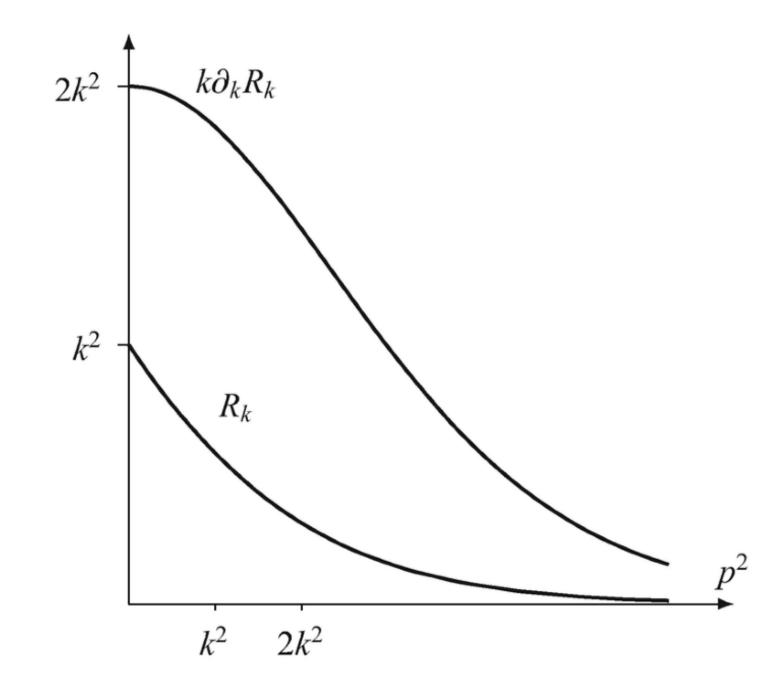
$$W_k[J] = \log \int \mathcal{D}\hat{\phi} \exp\left(-S[\hat{\phi}] - \Delta S_k[\hat{\phi}] + \int d^d x J(x)\hat{\phi}(x)\right)$$

Smooth cutoff

$$\Delta S_k[\hat{\phi}] = \frac{1}{2} \int d^d x \, \hat{\phi}(x) \mathcal{R}_k(-\Box) \hat{\phi}(x)$$

RG kernel: mass-like IR regulator

$$\mathcal{R}_k(p^2) \approx \begin{cases} k^2 & \text{for } p^2 \ll k^2 \\ 0 & \text{for } p^2 \gg k^2 \end{cases}$$



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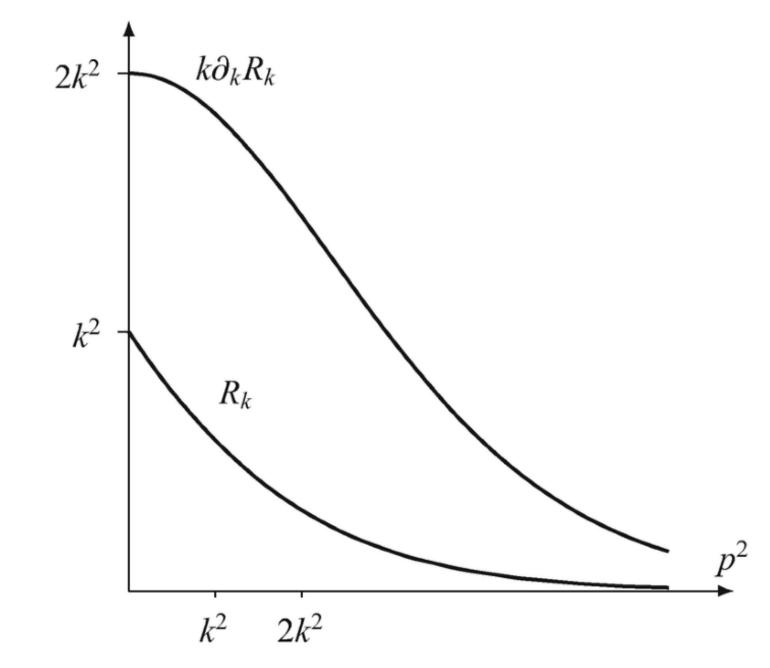
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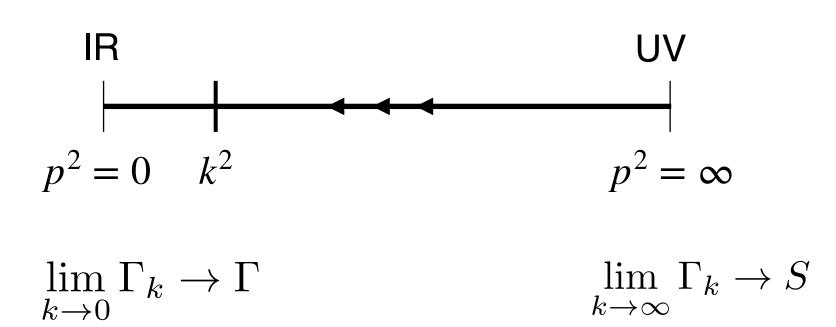


Legendre transform

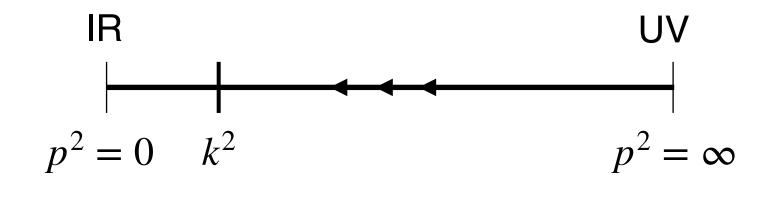
Effective Average Action

$$\tilde{\Gamma}_k[\phi] = \int d^d x \, J_k[\phi](x)\phi(x) + W_k[J_k[\phi]] \longrightarrow \Gamma_k[\phi] = \tilde{\Gamma}_k[\phi] - \Delta S_k[\phi] = \int d^d x \, J(x)\phi(x) + W_k[\phi] - \Delta S_k[\phi]$$

It represents the scaledependent version of the standard effective action.



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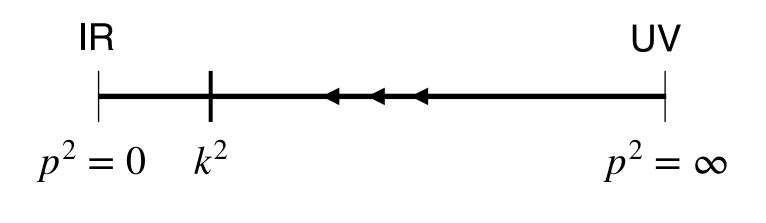
$$\lim_{k\to\infty}\Gamma_k\to S$$

It satisfies the Functional Renormalization Group Equation:

$$k\partial_k \Gamma_k[\phi] = \frac{1}{2} \operatorname{Tr} \left[(\Gamma_k^{(2)} + \mathcal{R}_k)^{-1} k\partial_k \mathcal{R}_k \right]$$

- UV- and IR finite
- Fully nonperturbative or exact

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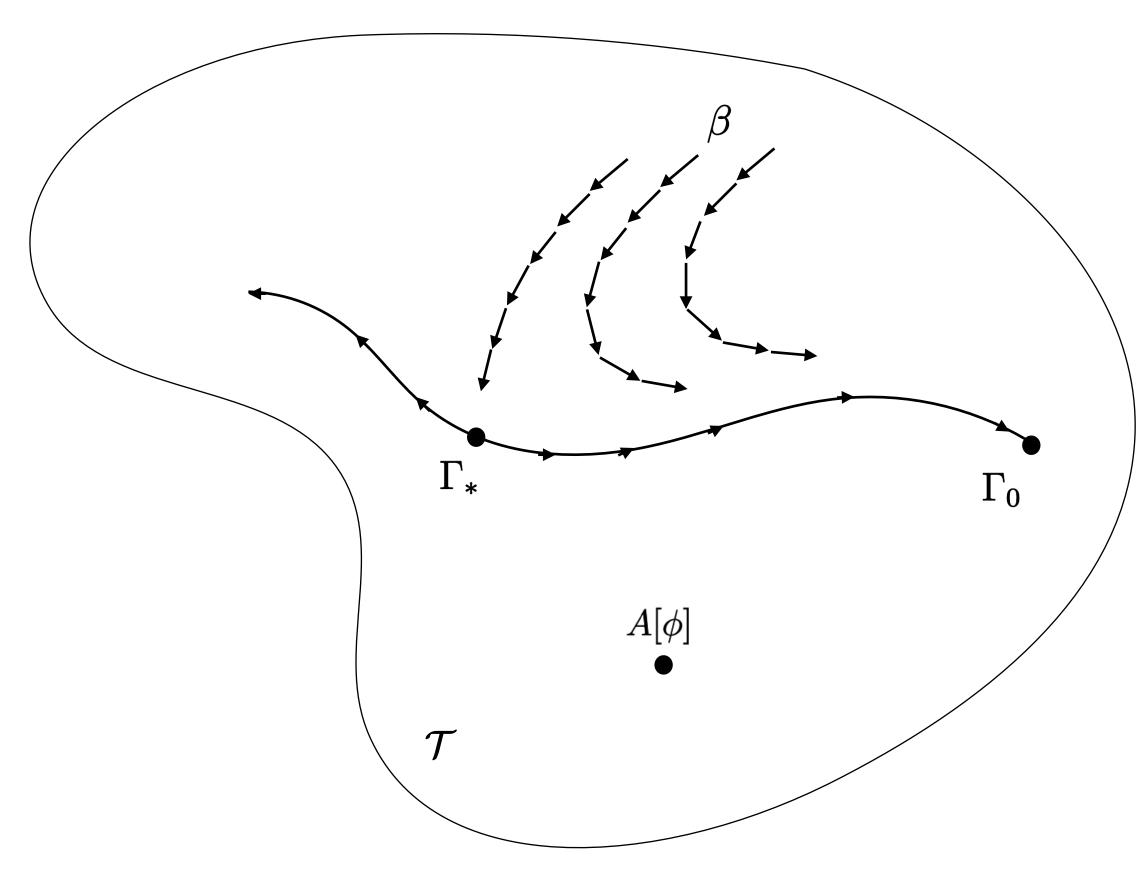
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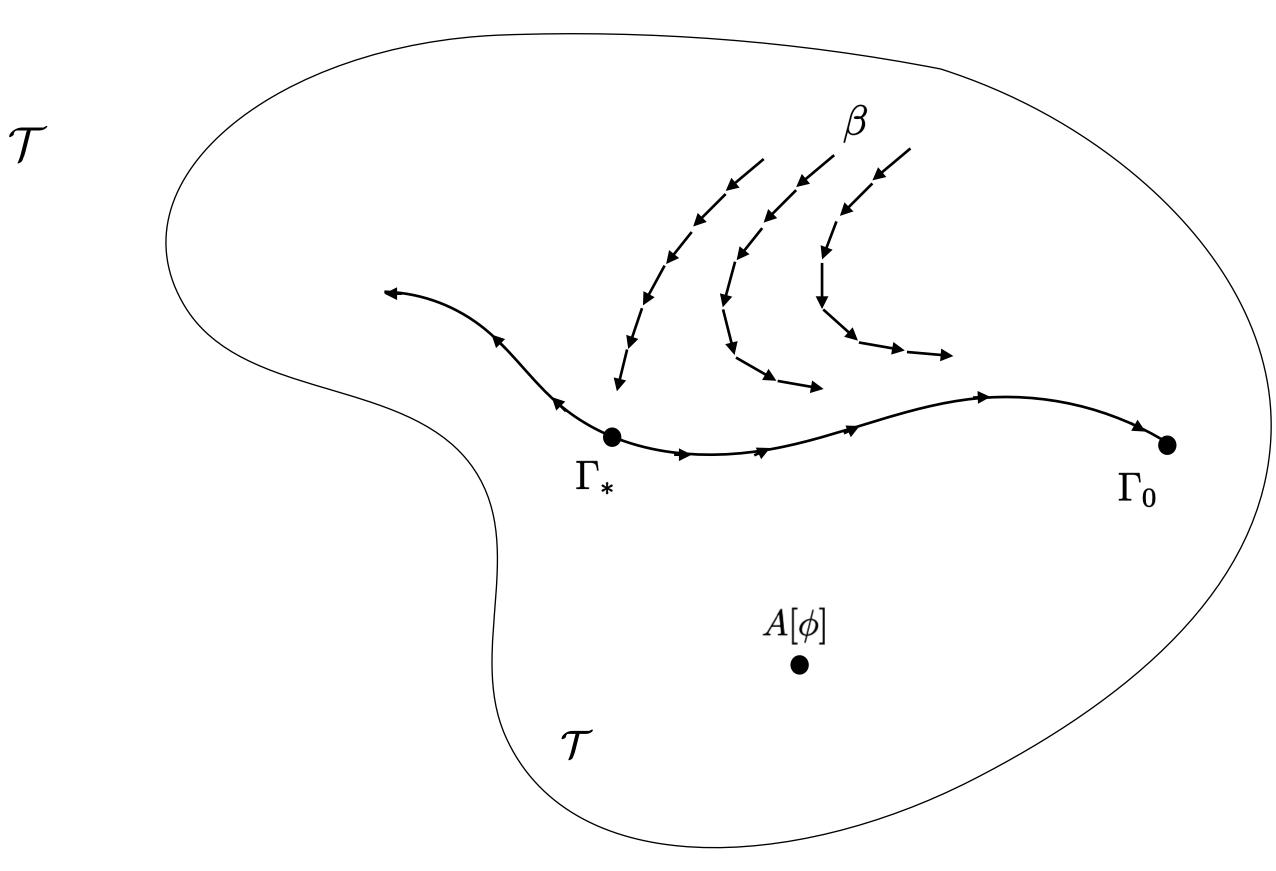
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Predictive solutions do exist in theories that are otherwise perturbatively non-renormalizable.



THEORY SPACE

space of functionals over which the EAA is supposed to be defined.

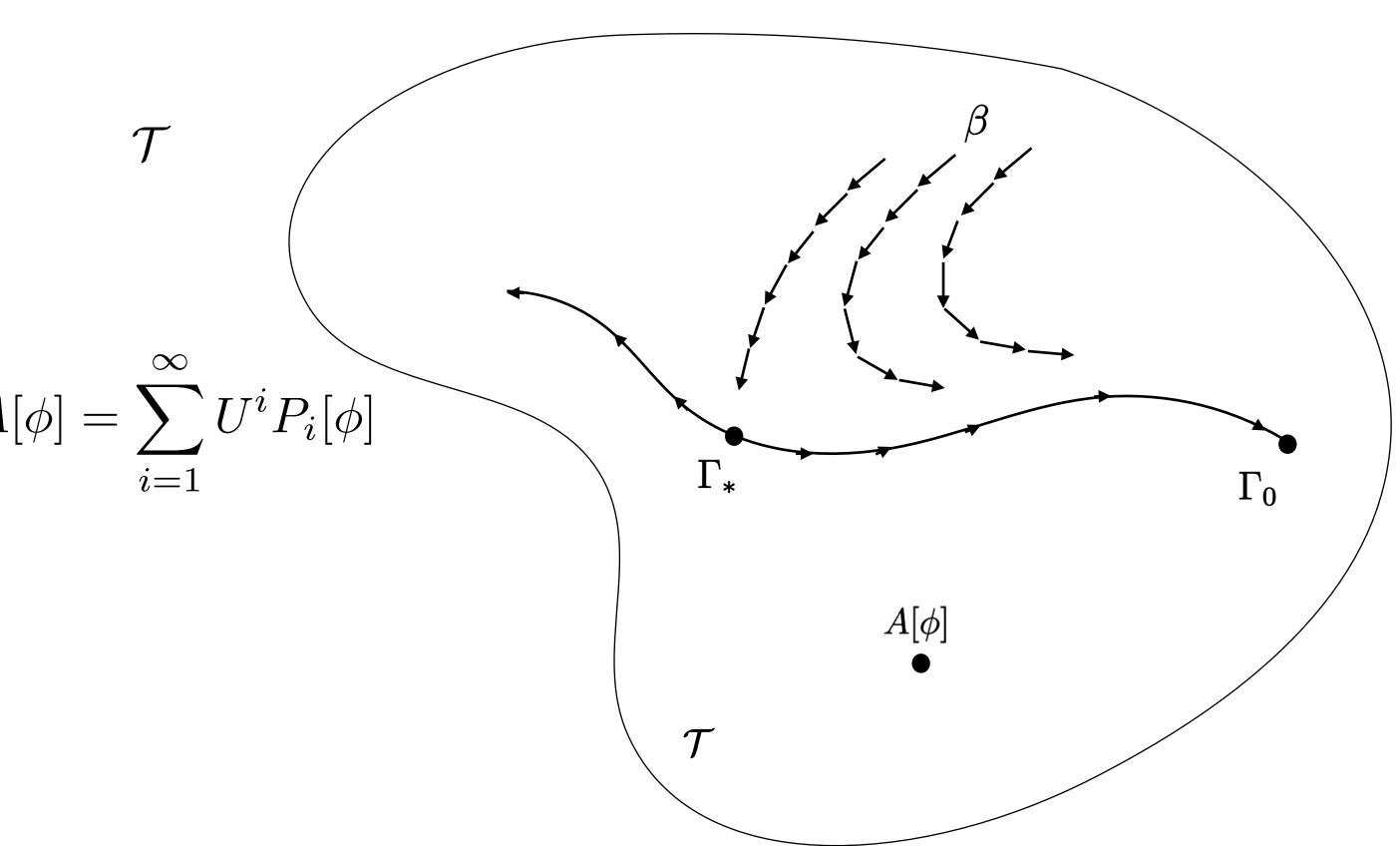


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BASIS ON THEORY SPACE

expansion of the elements of theory space in basis functionals and coupling constants.



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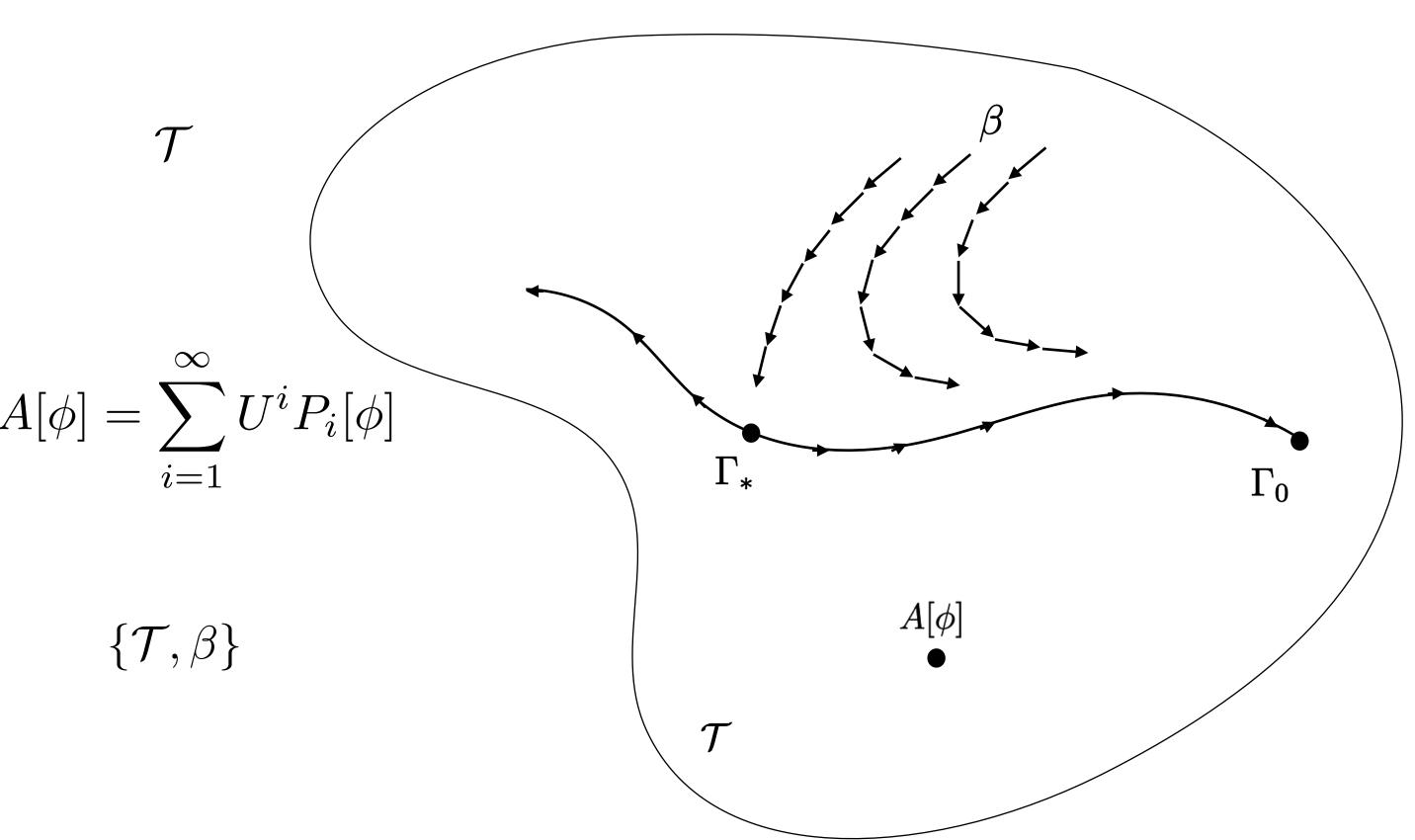
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vector field on theory space.



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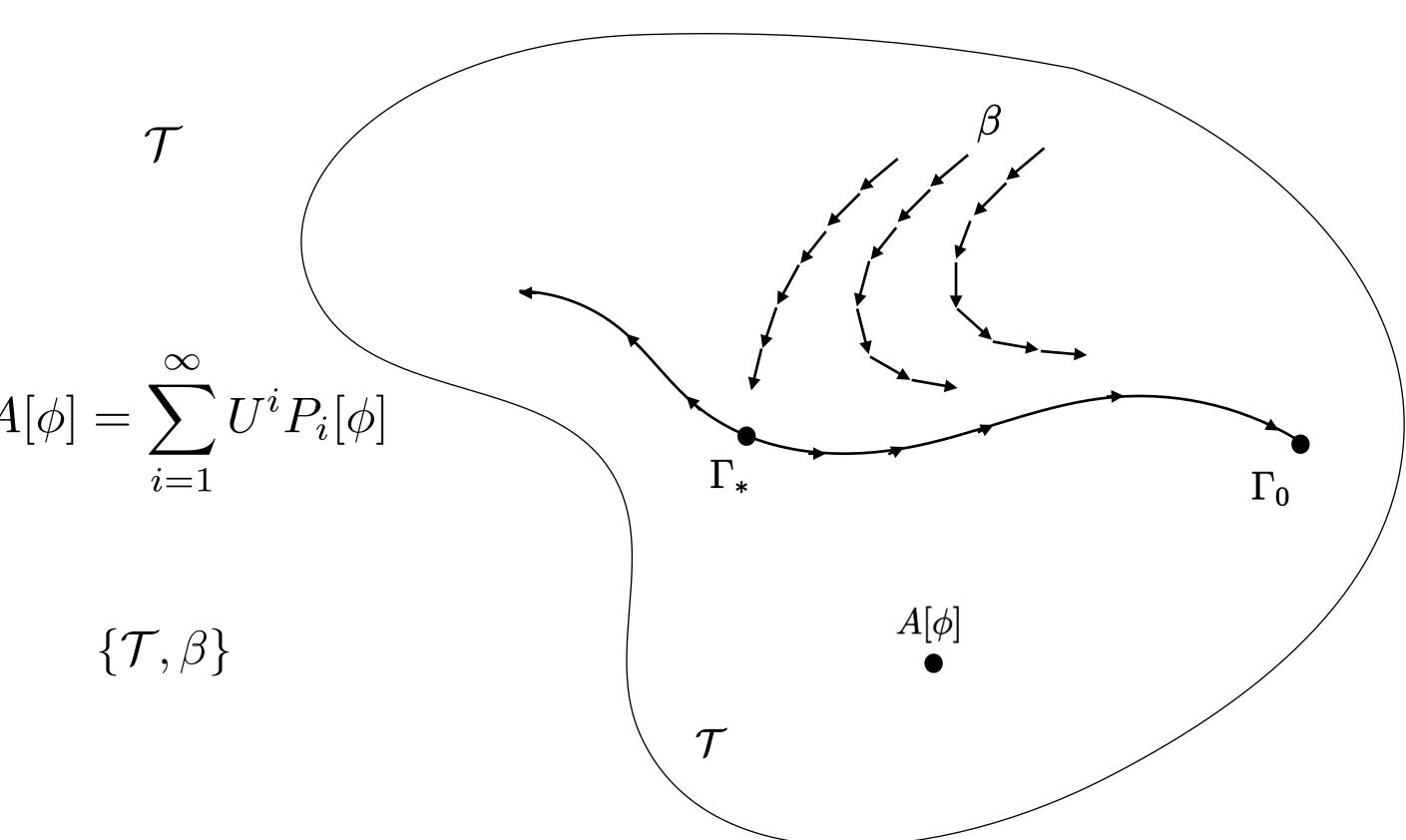
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pair of theory space and RG trajectories.



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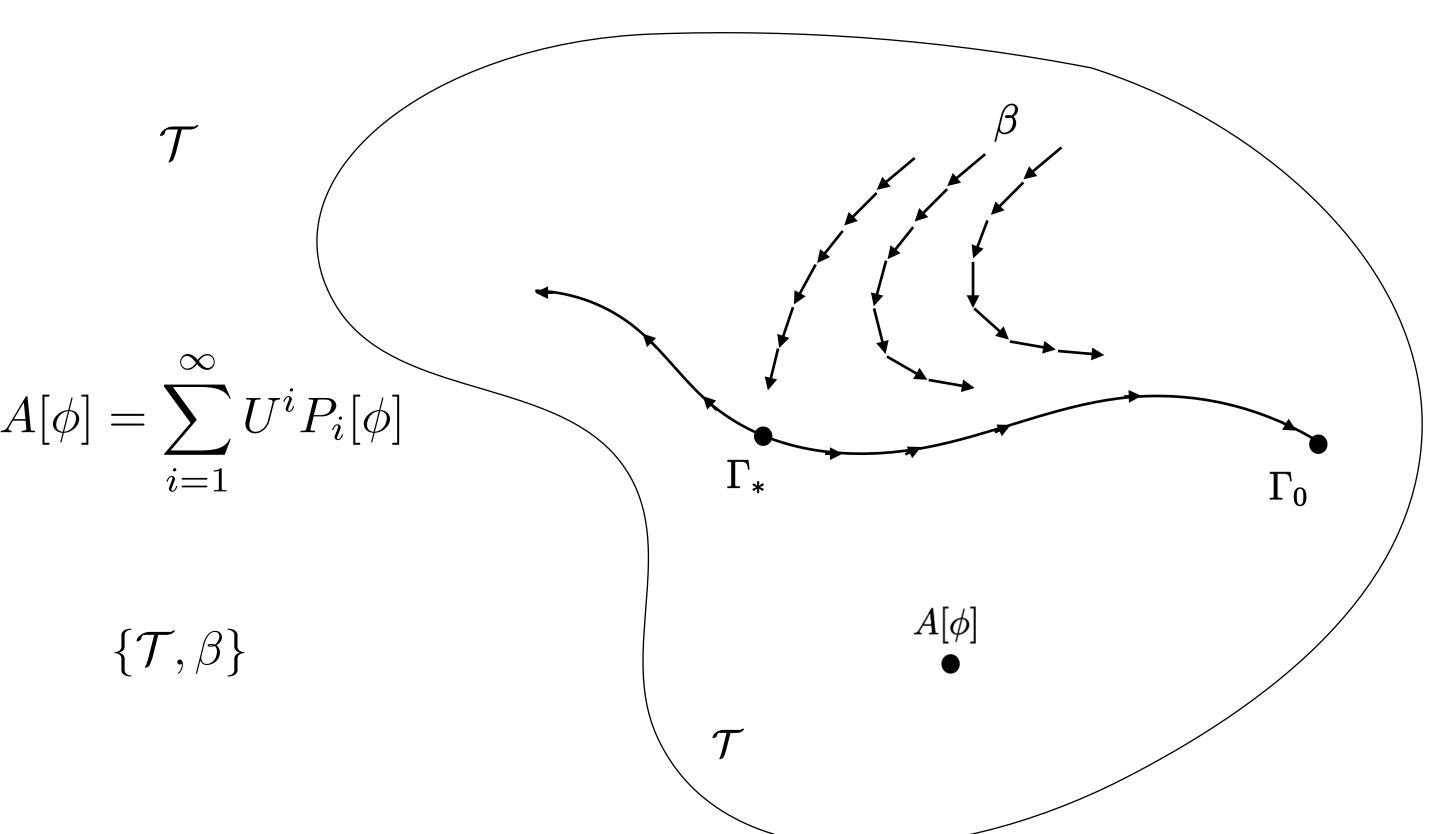
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$$\Gamma_k[\phi] = \sum_{i=1}^{\infty} U^i(k) P_i[\phi] \longrightarrow \frac{1}{2} \text{Tr}[\cdots] = \sum_{i=1}^{\infty} b^i(U^1, U^2, \cdots; k) P_i[\phi] \longrightarrow k \partial_k u^i = \beta^i(u^1, u^2, \cdots; k)$$

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$$\beta^i(u_*) = 0$$

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Gaussian FP: the scaling exponent agrees with the canonical mass dimension (generally $u_*^i = 0$)

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LINEARIZATION/ STABILITY MATRIX

$$k\partial_k u^i(k) = \sum_j \mathcal{B}^i{}_j(u^i(k) - u^i_*), \qquad \mathcal{B}^i{}_j(u_*) \equiv \partial_j \beta^i(u_*)$$

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SCALING EXPONENTS

$$u^{i}(k) = u_{*}^{i} + \sum_{I} C_{I} V_{I}^{i} \left(\frac{k_{0}}{k}\right)^{\theta_{I}}$$

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Gaussian FP: the scaling exponent

Non-Gaussian FP (interacting, UV FP):

at least one of the scaling exponents

differs from the canonical mass

dimension $(u_*^l \neq 0)$

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Gaussian FP: the scaling exponent agrees with the canonical mass dimension (generally $u_*^i = 0$)

Non-Gaussian FP (interacting, UV FP): at least one of the scaling exponents differs from the canonical mass dimension ($u_*^i \neq 0$)

They encode physical information about the universality class of the system and its scaling (observable) behavior.

Asymptotic Safety

Asymptotic Safety

a mechanism which renders physical scattering amplitudes finite (but non-vanishing) at energy scales exceeding the Planck scale.

Asymptotic Safety

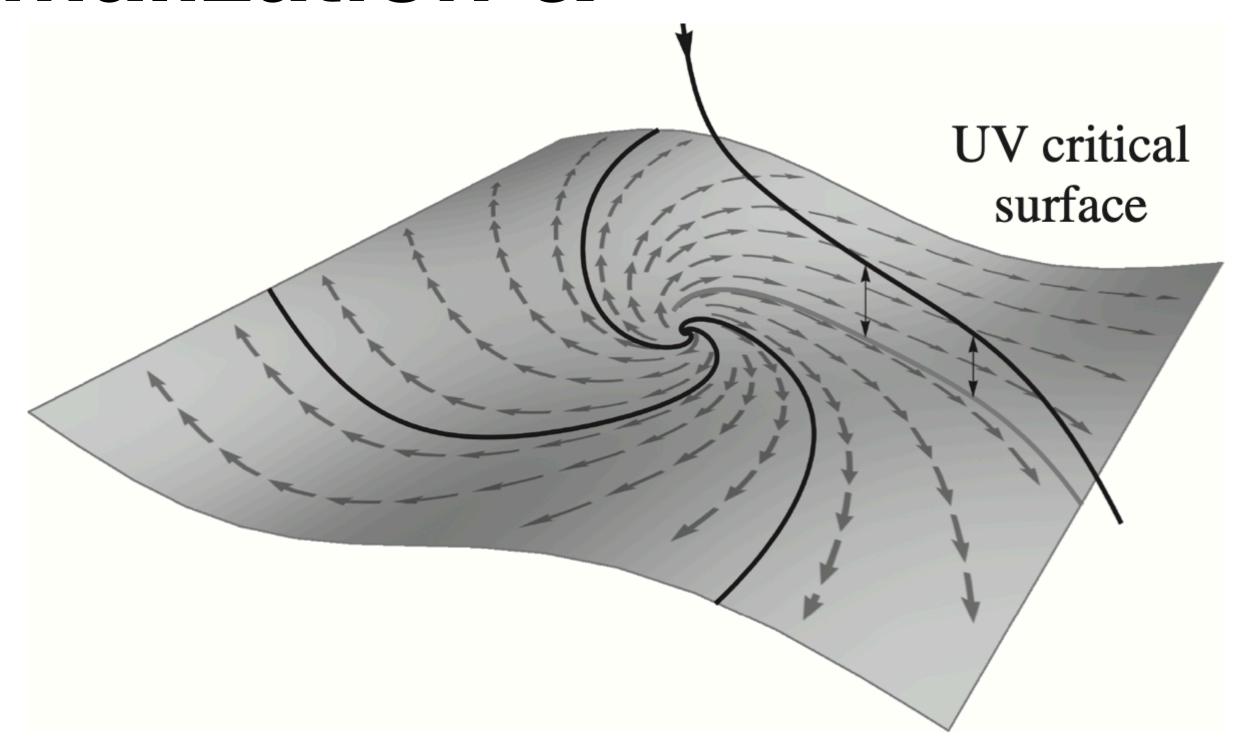
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À LA EFFECTIVE AVERAGE ACTION

nonperturbative renormalization &

predictivity

Renormalization consists in constructing a complete trajectory, a trajectory which lies entirely within theory space.



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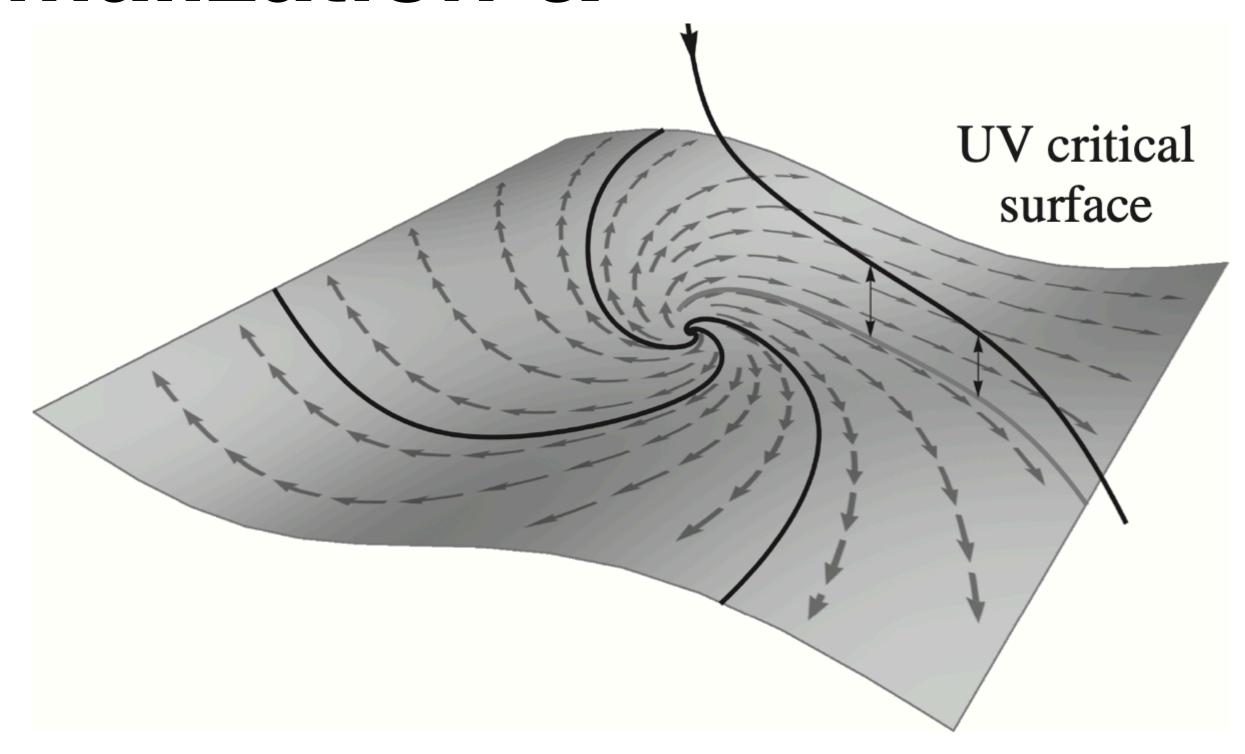
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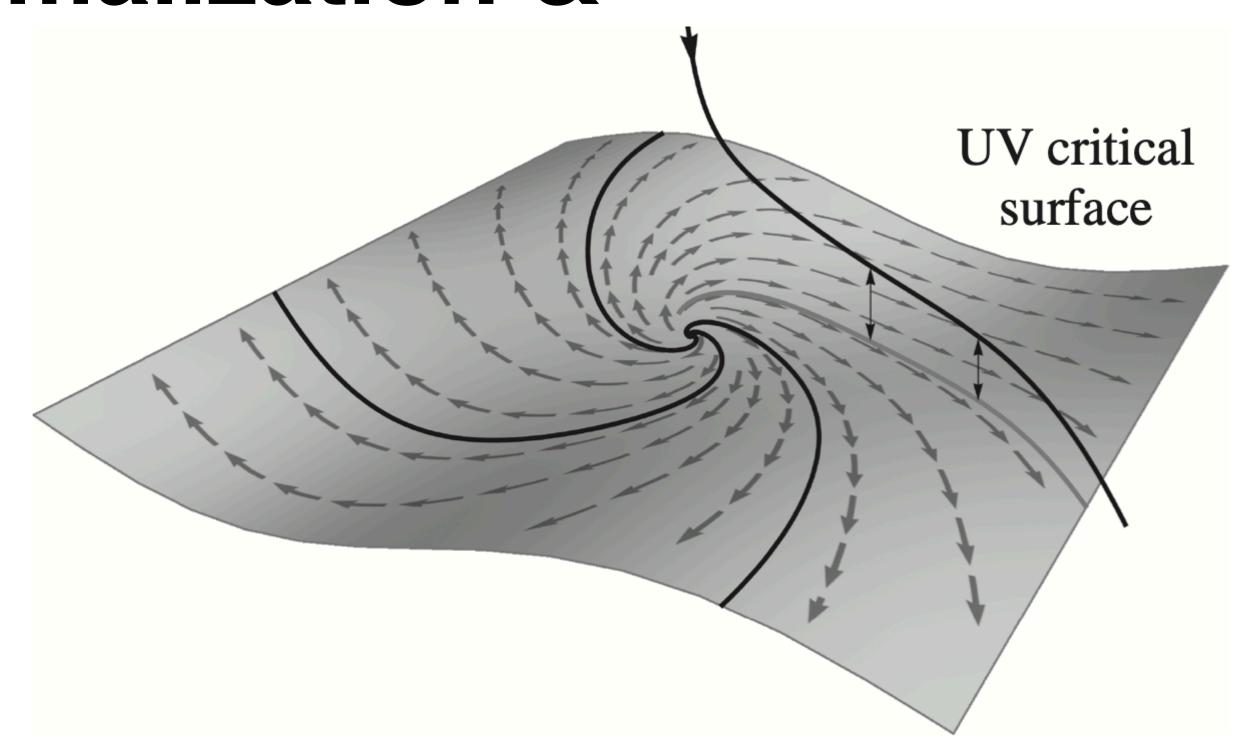
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Couplings:

Relevant: grow in the IR $\operatorname{Re} \theta_i > 0$

Irrelevant: shrink in the IR $\operatorname{Re} \theta_i < 0$



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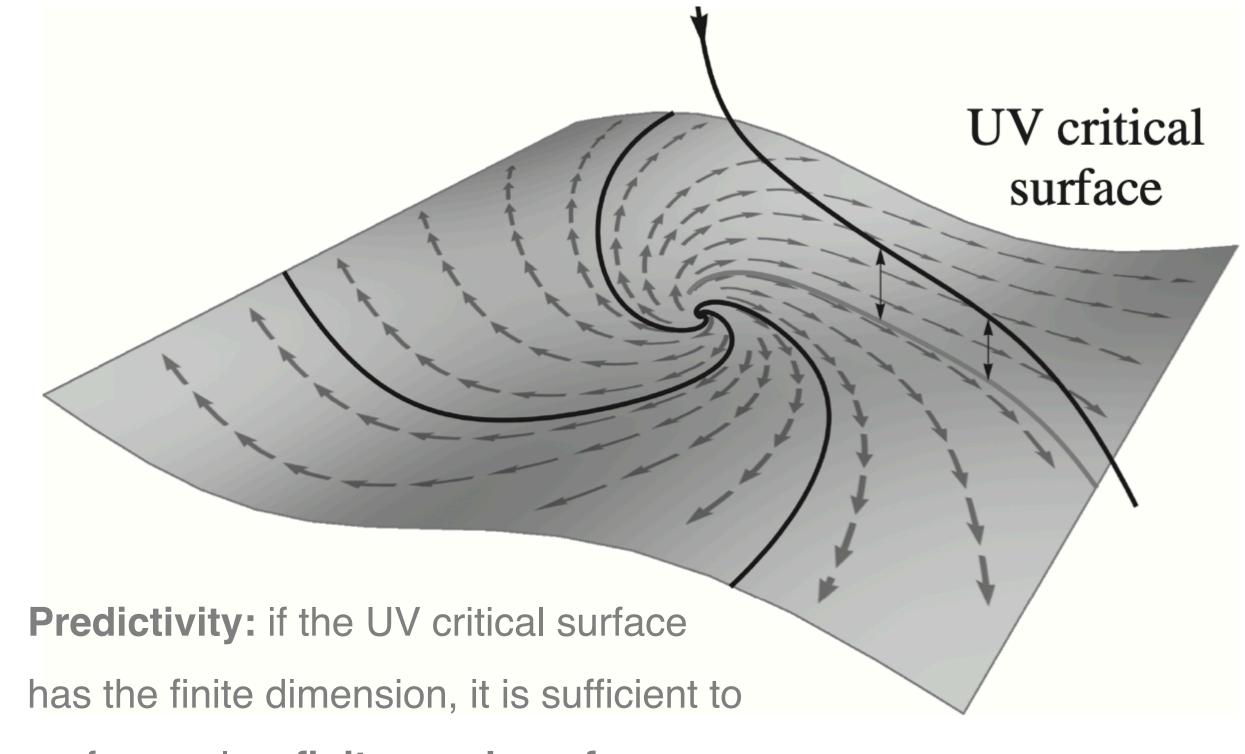
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perform only a finite number of

measurements in order to uniquely

identify Nature's RG trajectory.

Nonperturbative approximate solution

$$\Gamma_k[\phi] = \sum_{i=1}^N U^i(k) P_i[\phi]$$

$$k\partial_k \sum_{i=1}^N U^i(k) P_i[\phi] = \frac{1}{2} \operatorname{Tr} \left[\left(\sum_{i=1}^N U^i(k) P_i^{(2)}[\phi] + \mathcal{R}_k \right)^{-1} k \partial_k \mathcal{R}_k \right]$$

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LHS: beta functions

RHS: Evaluate the trace using heat kernel techniques

- Compare term by term the coefficients multiplying the same operator
- Solve analytically/numerically the differential equations

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A natural ordering principle for the interaction terms entering into the EAA is provided by their number of derivatives.

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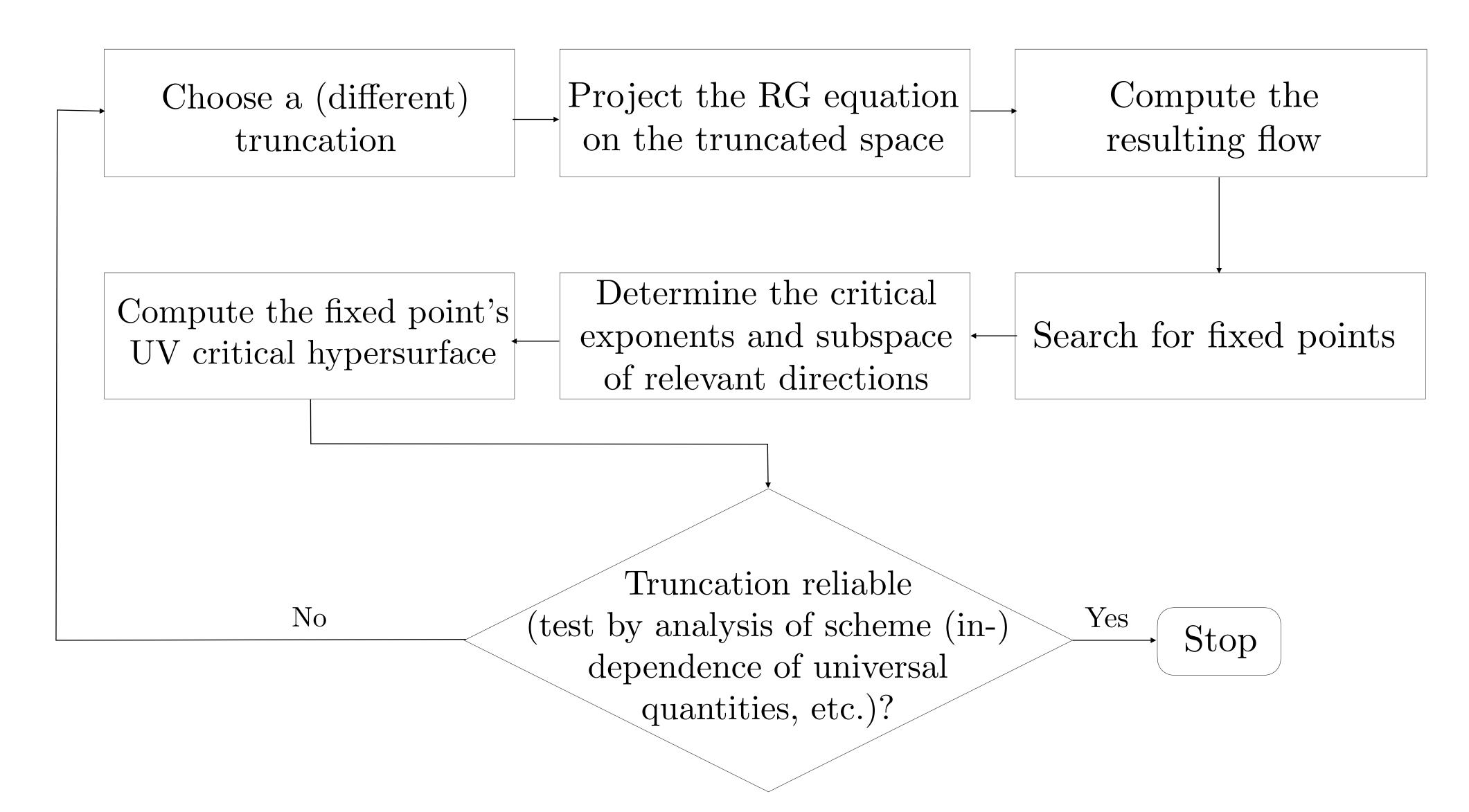
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The Program



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Background field method

$$g_{\mu\nu} = \langle \hat{g}_{\mu\nu} \rangle = \bar{g}_{\mu\nu} + h_{\mu\nu}$$

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The dynamical degrees of freedom are quantized on this background.

$$\langle \hat{h}_{\mu\nu} \rangle_{\bar{g}} = 0 \iff \langle \hat{g}_{\mu\nu} \rangle_{\bar{g}} = \bar{g}_{\mu\nu} \quad \text{for} \quad \bar{g} = \bar{g}(k)^{\text{sc}}$$

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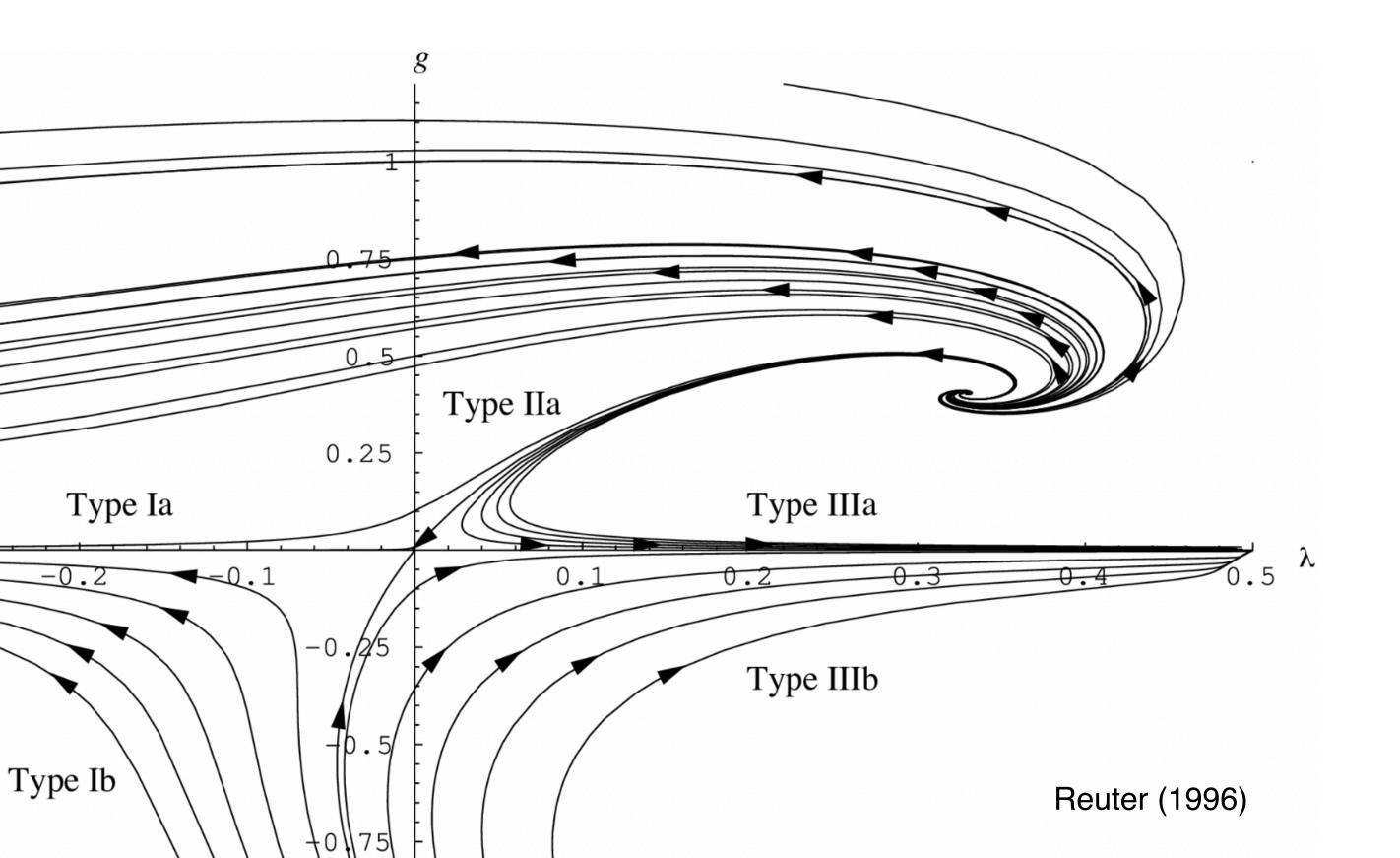
The dynamical degrees of freedom are quantized on this background.

The background is re-adjusted in such a way that it becomes self-consistent, meaning that the expectation value of the fluctuation vanishes.

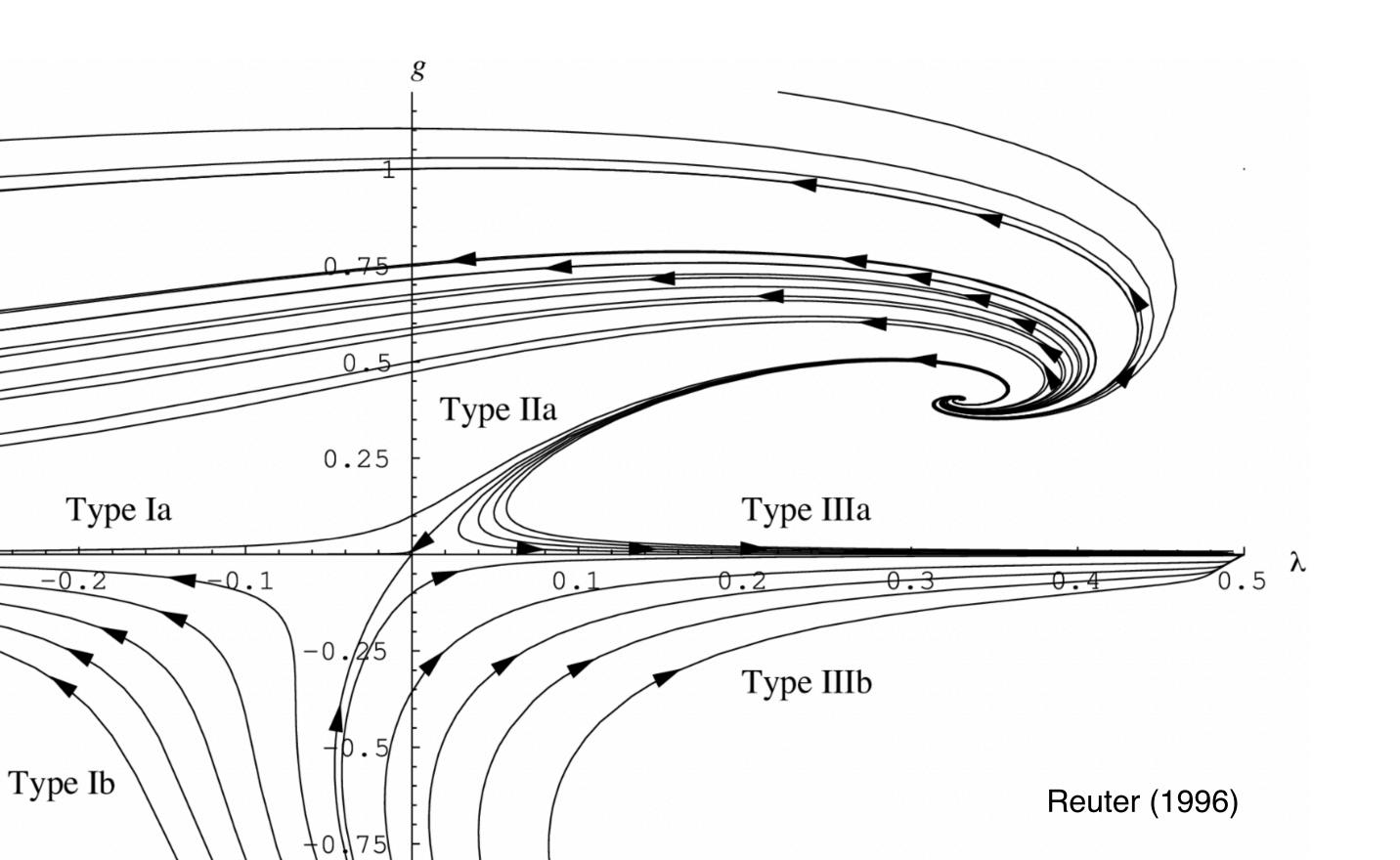
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$$\frac{\delta}{\delta h_{\mu\nu}(x)} \Gamma_k [h; \bar{g}] \bigg|_{h=0, \ \bar{g}=\bar{g}(k)^{\text{sc}}} = 0$$

$$\Gamma_k = \frac{1}{16\pi G(k)} \int d^d x \sqrt{g} (R - 2\Lambda(k)) + \text{gauge fixing} + \text{ghosts}$$

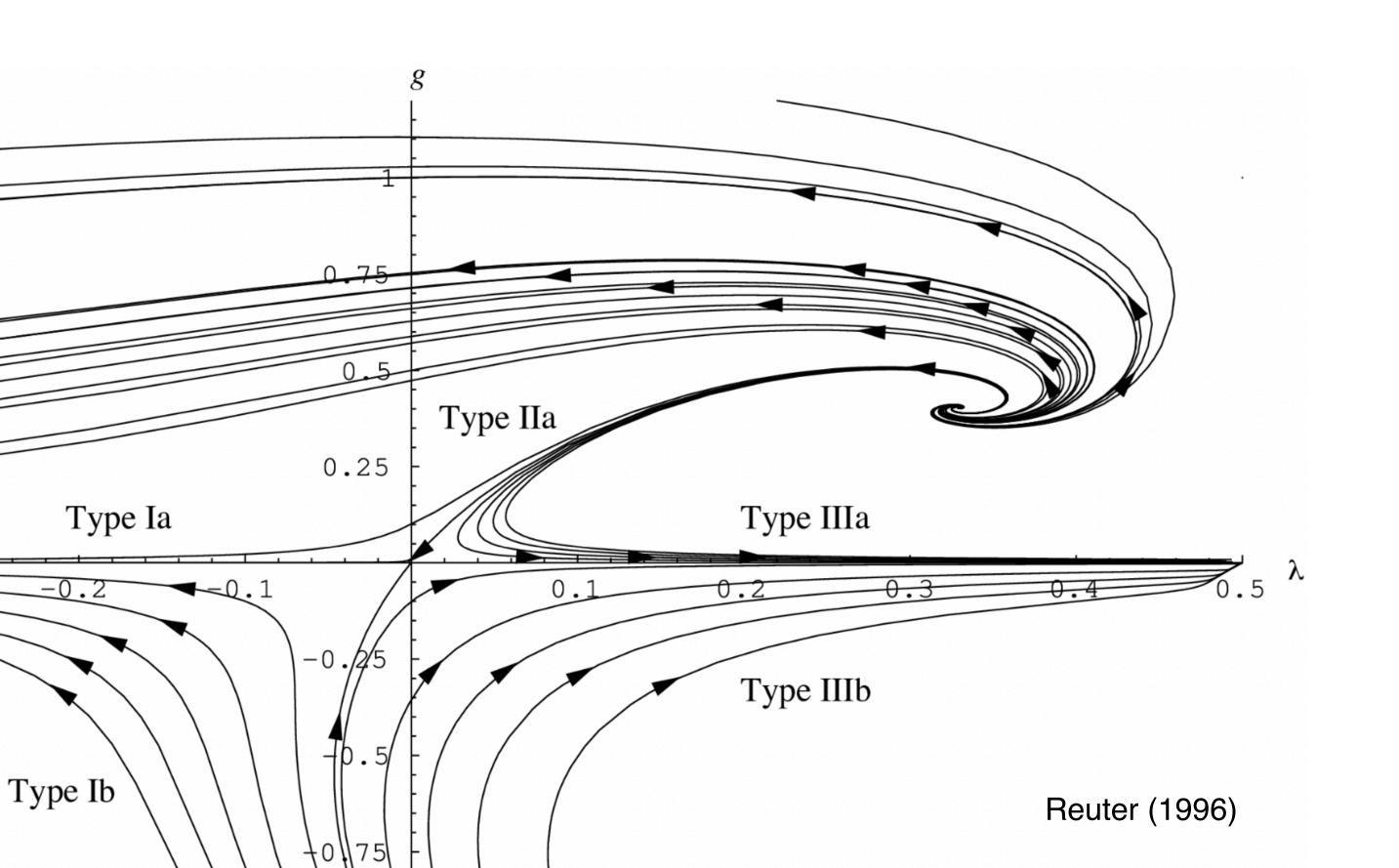


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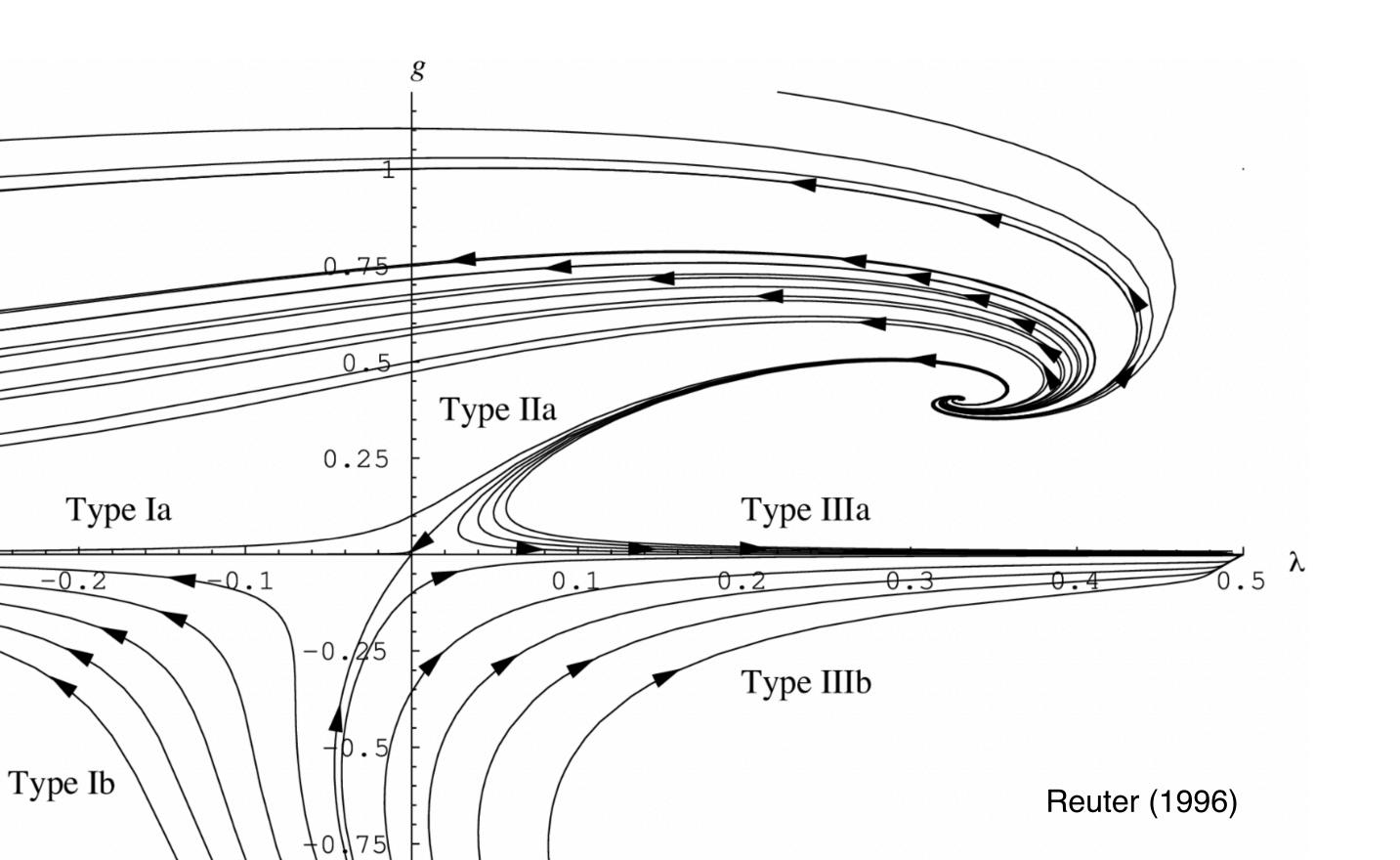
$$k\partial_k g_k = \beta_g(g,\lambda)$$

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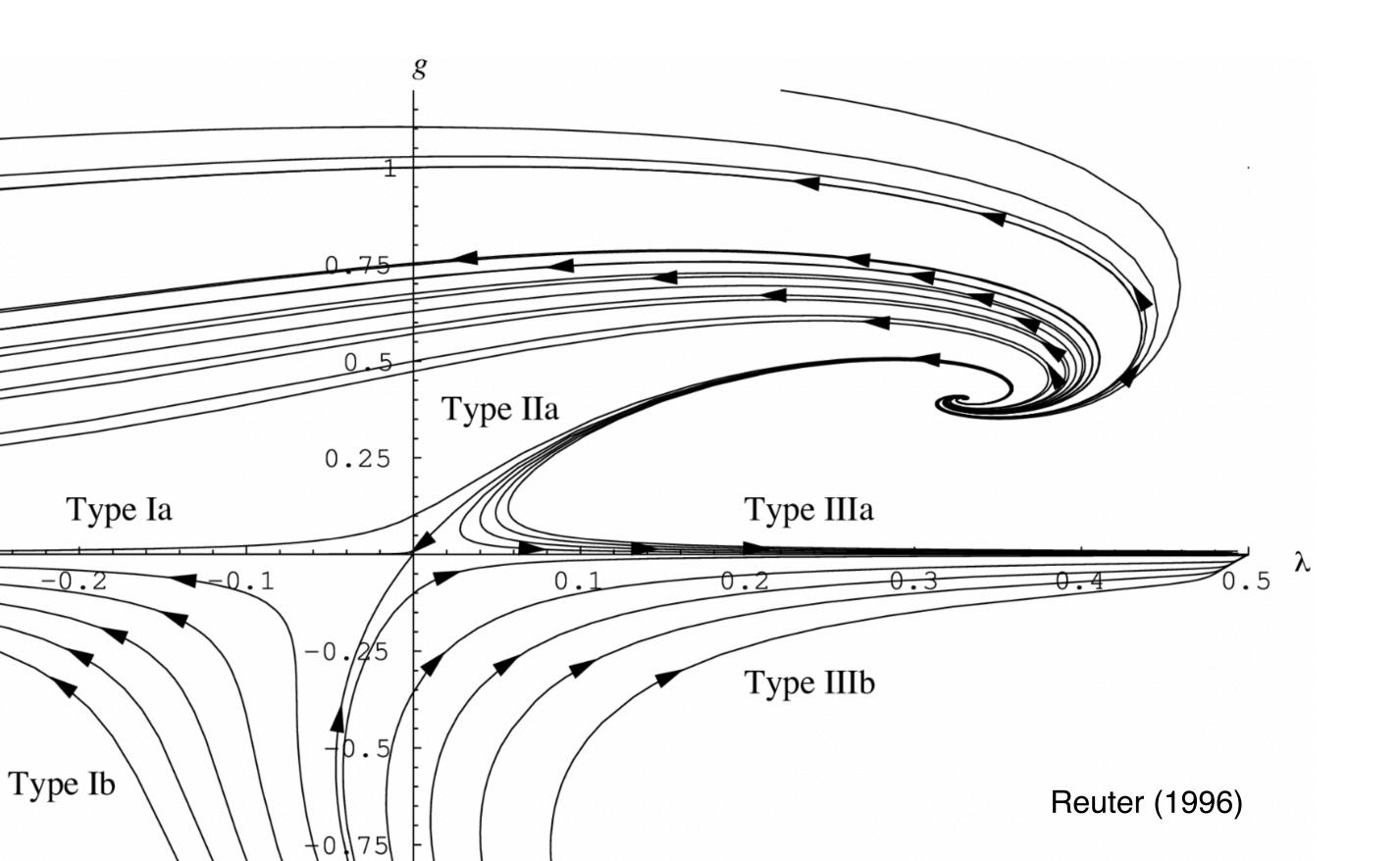
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Non-trivial fixed point.

Complete trajectories.

Finite number of relevant directions.

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ASYMPTOTIC SAFE

Finite number of relevant directions.

Integrating out dofs Timelike-spacelike: no distinguished ordering of the modes with a standard canonical status is given.

State dependence - Observer dependence

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Problem set 1

Problem set 2

Integrating out dofs

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Problem set 1

Obtain RG trajectories on a theory space which is constituted of functionals that are constructed on Lorentzian metrics.

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The FRG as a machinery

One can work out a Lorentzian heat kernel proper time regularization.

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Problem set 2

Analyze the flows of hyperbolic kinetic operators, typically of the d'Alembertian in the background of the running self-consistent metrics.

RF, Reuter (2022)

Which physical information can we extract from an asymptotic safe theory of gravity?

Observables?

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Observables?

Relational observables.

Two options.

"Physical gauge fixing"

via material reference frame ~ reduced phase space quantisation

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Two options.

Construct relational observables as additional operators

after gauge-fixing include matter (physical reference frame) in the EAA

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Kevin Falls, RF (2022)

Composite operators

Running relational observables

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Observables

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Pagani, Reuter (2016)

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$$\lim_{k \to \infty} \mathcal{O}_k = \hat{\mathcal{O}}|_{\hat{\phi} \to \phi}$$

$$\mathcal{O}_{k=0} = \langle \hat{\mathcal{O}} \rangle$$

Covariant formulation

Running relational observables

Covariant formulation

Construct a physical coordinate frame e.g. by adding matter fields

s.t.

perform a diffeomorphism transformation

transform the tensor

transform the physical frame

composed transformation leaves the tensor invariant

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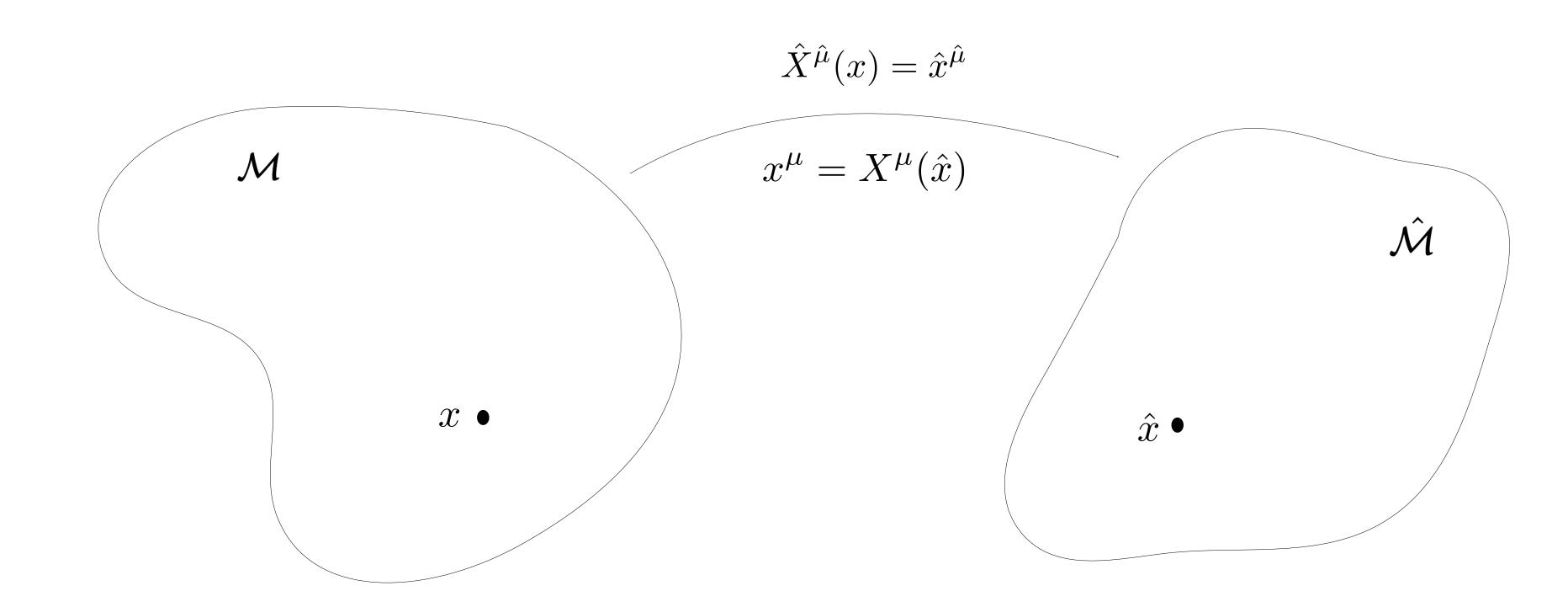
$$R(x) \mapsto \varphi * R(x)$$

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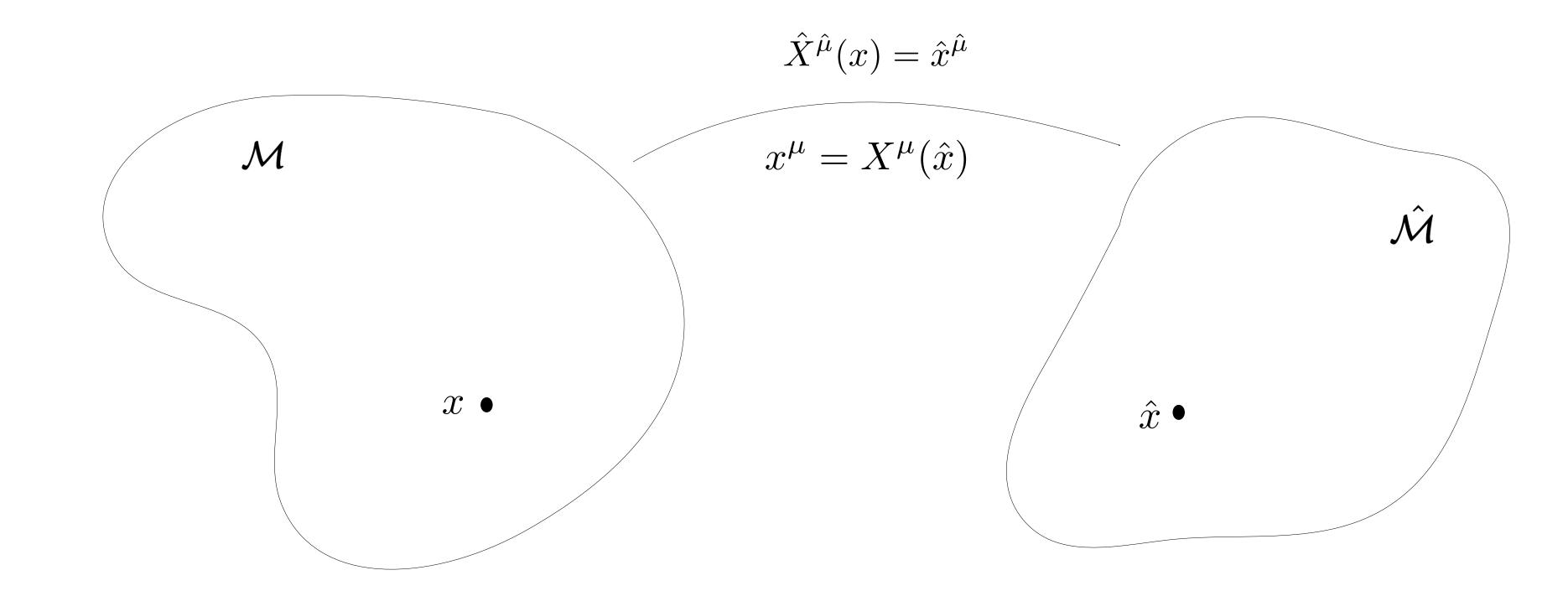
composed transformation leaves the tensor invariant

$$R(X) \mapsto \varphi * R(\varphi^{-1}(X)) = R(X)$$



Introduce a dynamical fields

$$\phi^a(x) = \{g_{\mu\nu}, \text{matter fields}\}\$$

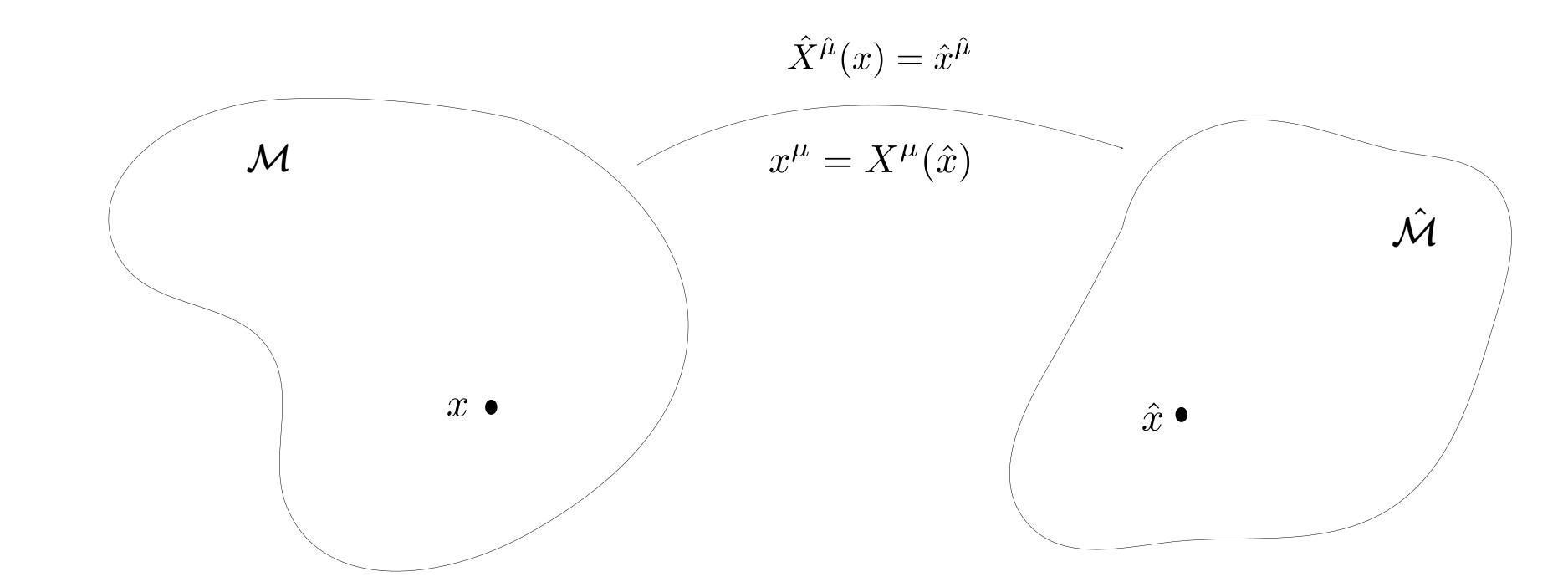


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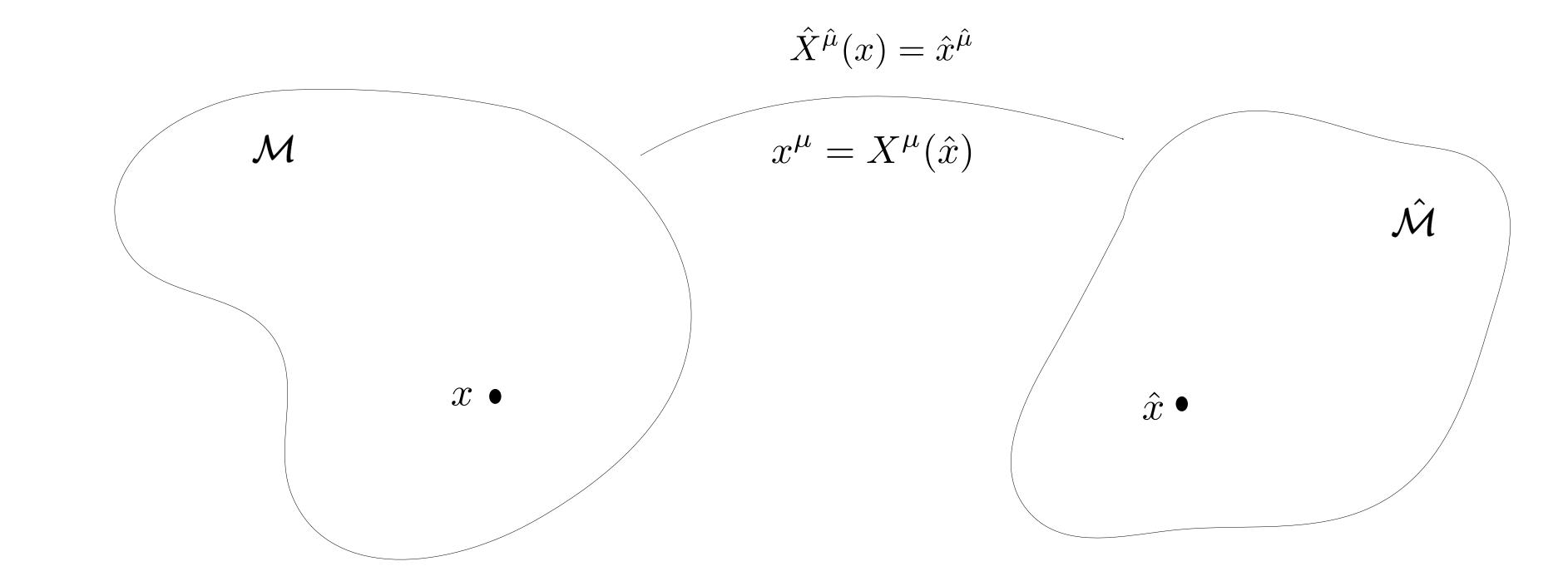
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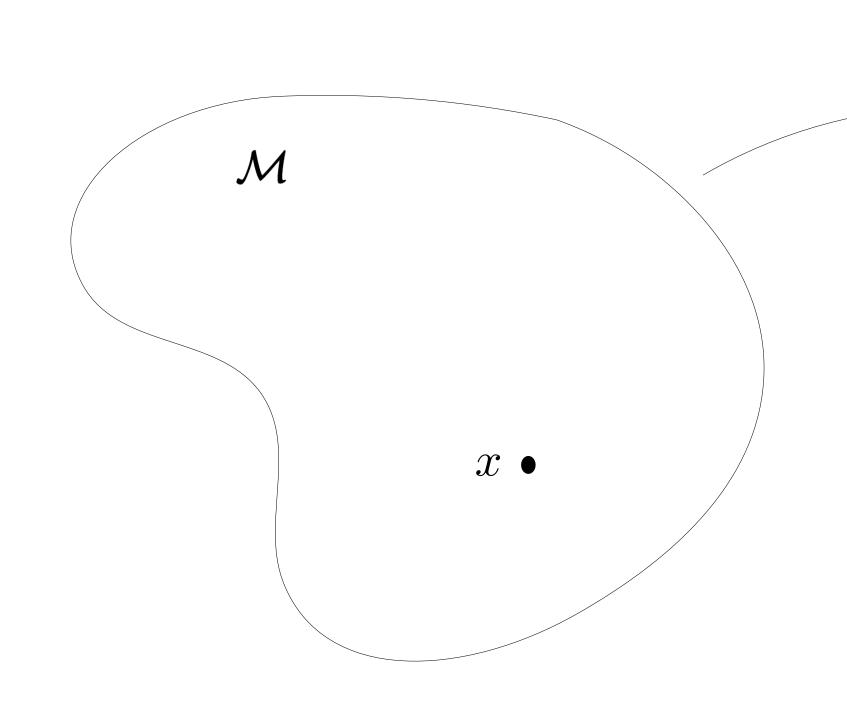
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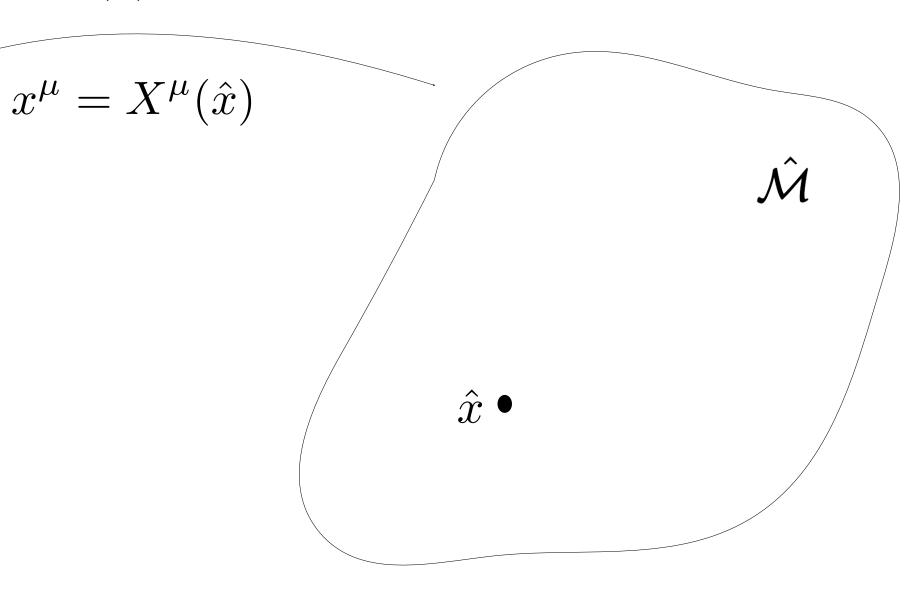


Frame fields: $e_{\mu}^{\hat{\mu}}(x) = \partial_{\mu} \hat{X}^{\hat{\mu}}(x)$

Invariant volume element: $ilde{e} = \det e_{\mu}^{\hat{\mu}}$

$$\delta(X(\hat{x}), x) = \tilde{e}(x) \,\delta(\hat{x}, \hat{X}(x))$$

$$\hat{X}^{\hat{\mu}}(x) = \hat{x}^{\hat{\mu}}$$



Relational Effective Average Action:

$$\Gamma_k^{\text{rel.}} \equiv \int d^4x \, \tilde{e}(x) \mathcal{L}_k^{\text{rel.}}(x) := \int d^4x \, \tilde{e}(x) \sum_i a_{\hat{I}_i}(k) E_{iI_i}^{\hat{I}_i}(x) O_i^{I_i}(x)$$

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$$= \int d^4 x \, \tilde{e} \left(a_0(k) + a_R(k) R + a_1(k) \delta_{\hat{\mu}\hat{\nu}} g^{\mu\nu} (\partial_{\mu} \hat{X}^{\hat{\mu}}) (\partial_{\nu} \hat{X}^{\hat{\nu}}) \right)$$

Find the fixed points

EAA

Identify the relational observables

PHYSICAL REFERENCE SYSTEM

Compute the flow of the observables

RELATIONAL EAA

Scaling dimension at the fixed point

Find the fixed points

EAA

$$\Gamma_k = \int d^4x \sqrt{-g} \left(\frac{1}{16\pi G_N(k)} (R - 2\Lambda(k)) + \frac{1}{2} \delta_{AB} g^{\mu\nu} \partial_{\mu} \varphi^A \partial_{\nu} \varphi^B \right) + \dots$$

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RESULT

Matter content

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SM + SF (type II)			
$SM + 3 \nu \text{ (type II)}$			

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θ_0	θ_R	θ_1
-4	-5.97643	-7.92358
-4	-5.97467	-7.8177
-4	-5.97505	-7.80603
-4	-5.98015	-7.78084

Small quantum

corrections.

-4 -6 -8

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- 6

Compare?

Small quantum corrections.

- 8

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at the level of the scaling exponents, e.g. scaling of the correlation functions

at the level of the RG computations