

FUNDAÇÃO DE AMPARO À PESQUISA DO ESTADO DE SÃO PAULO



Magnetic Reconnection, and particle acceleration and propagation around Black Holes

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Municipal theater

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- 6. Conclusions & take-home message



- Solar flares;





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$$V_{\rm rec} = V_A S^{-1/2}$$

where the Lundquist number is

$$S = \frac{LV_A}{\eta}$$

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Turbulent Magnetic Reconnection



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Lazarian & Vishniac, 1999, ApJ



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 λ_{\parallel}

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"Big power law in the sky" (ISM)

(Armstrong, Rickett & Spangler, 1995, ApJ)



Neutrino VS gamma-ray flux from NGC1068



Electromagnetic observations (26)

0.1 to 100 GeV gamma-rays (40,41)







hotopion (π^{\pm} component





component





hotopion (π^0 component



otopion $(\pi^0 \text{ component})$



hotopion $(\pi^0 \text{ component})$





air production



 $\overline{\mathbf{IC}}$

ducti

(IceCube Collaboration, 2022, Science)



part of the AGN!

ir production

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Electromagnetic observations (26) 0.1 to 100 GeV gamma-rays (40,41) > 200 GeV gamma-rays (42)



What may accelerate these protons in the surroundings of the SMBH?



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Possible configuration of the magnetic field lines for an accretion flow into a black hole









Particles can be accelerated in the magnetic discontinuity according to a first-order Fermi process:

 $V_{\rm rec}$





Implies an exponential growth of the energy with time!

 ΔE

F

 $V_{\rm rec}$







(Rodríguez-Ramírez, de Gouveia Dal Pino & Alves Batista, 2019, ApJ)



Search for reconnection sites in 2D & 3D GRMHD simulations of accretion flows



(Kadowaki et al. 2018; de Gouveia Dal Pino et al. 2018)



In classical regime, we use the AMUN code (Kowal, 2009) to solve the isothermal non-ideal MHD equations:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \boldsymbol{v}) = 0,$$

$$\frac{\partial \rho \boldsymbol{v}}{\partial t} + \nabla \cdot \left[\rho \boldsymbol{v} \boldsymbol{v} + \left(a^2 \rho + \frac{B^2}{8\pi} \right) I - \frac{1}{4\pi} \boldsymbol{B} \boldsymbol{B} \right] = \boldsymbol{f},$$

$$\frac{\partial A}{\partial t} + E = 0,$$



and the initial magnetic field is given by

$$\vec{B} = B_z \hat{z} + \hat{z} \times \nabla \psi,$$

where

$$\psi = \frac{1}{2\pi} \tanh\left(\frac{y}{h}\right) \cos(\pi x) \sin(2\pi y) \overset{\scriptscriptstyle n}{\longrightarrow}$$





(Vicentin et al., in prep.)



(Vicentin et al., *in prep.*)



We can recover the Sweet-Parker regime for low-Lundquist numbers!

The reconnection rate starts to deviate from SP for high-Lundquist numbers (plasmoid instability)

(Vicentin et al., *in prep.*)

Time = 2.40



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For the case where we inject forced turbulence initially in the domain (up to $0.1 t_A$), we can reach high values of

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even after the turbulence injection is stopped.





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- Turbulent magnetic reconnection around BHs can explain VHE emission from these compact sources
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- Turbulent magnetic reconnection around BHs can explain VHE emission from these compact sources
 - from theory and global GRMHD sims.
- Classical MHD simulations have showed that the system remains turbulent even after injection is stopped





Any question?

Giovani H. Vicentin IAG USP

