

# ELKO SPINOR AS DARK MATTER

## specific signatures and constraints for the interacting model

(Recently accepted in Nucl. Phys. B)

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# Table of Contents

- 1 General remarks
- 2 Definition of ELKO spinor
- 3 Amplitudes

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**Dark matter** has strong evidence due to gravitational effects such as

- Galaxy rotation curve
- gravitational lensing effects

**However,**

- Low interaction with standard matter particles

Therefore, some **NON-Luminous** matter must be throughout the universe.

**According to cosmology 85% of the total mass of the universe is associated to dark matter.**

# General remarks

- Standard gravitational theory implies in decaying velocity outside the region with the major quantity of matter (Luminous) at  $R \approx 15$  (kpc) light years.
- The observed curve must be explained by some non-luminous matter in the region  $> R$

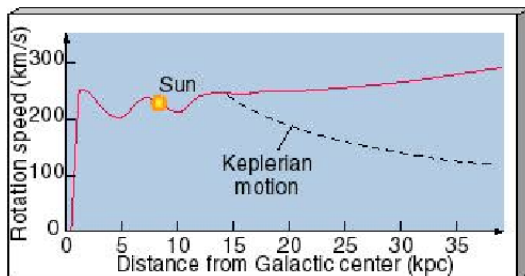


Figure: credit-astro.physics.uiowa.edu. (Milk way)

**ELKO** spinor is defined as eigenspinors of charge conjugation operator. It transforms differently than Dirac ones under discrete symmetries  $C; P; T$ . They are candidates for a non-standard Wigner class.

**It has some important characteristics:**

- Are the expansion coefficients of Mass dimension one fermion fields
- Each species Just obey KG equation
- Then, **cannot** enter the standard model doublets in a **renormalizable** manner
- It is also important to highlight the fact that ELKO spinors enter as CLASS 5 spinors in LOUNESTO classification, DIFFERENTLY FROM MAJORANA fermions. Another characteristic feature is the fact that they DO NOT HAVE A DEFINITE HELICITY, each quiral component has a different one.

# XENON-nT (noble gas)-Gran Sasso-Italy

- On the ground, low contamination, and  $v = 232\text{Km/s}$  with relation to DM approximately in rest in the center of the galaxy.
- Recognize scintillation due to nucleon recoil.
- Well known background. Searching for excess with statistical significance in terms of standard deviation.



# Experimental searches

However, there are recent searches for Dark matter in nucleon recoil LAB experiments

- Xenon-nT, Lux, zeplin, CDMS, etc (scintillation)
- **phonon mediated detections-2022**
- Cross section  $10^{-39} \text{cm}^2$
- LAB **WIMP Mass expected to be of order  $GeV$**  event rate peak 2206.06772
- Tritium contamination in XENON (2020): fake  $m \approx 2; 3 \text{Kev}$  WIMP
- Previous estimate from ELKO and Galaxy rotation curve  $m \approx 0; 1 \text{KeV}$

To be detected in LAB (ordinary matter devices.) there must be coupling involving an **INTERMEDIATE boson** and the ELKO and also one involving this boson and the nucleons



- Considering an approximately flat universe.
- **A coupling with photon would be fine** (It is possible to find an exotic RENORMALIZABLE one.)
- But how to keep the darkness property with such coupling?
- **We show that it leads to ZERO non-polarized Compton-like and pair annihilation squared amplitudes!**

# Table of Contents

- 1 General remarks
- 2 Definition of ELKO spinor
- 3 Amplitudes

# Defining ELKO

Expressing the CHARGE CONJUGATION involutive operator as

$$C = \begin{pmatrix} 0 & i \\ i & 0 \end{pmatrix} K; \quad (1)$$

One can define the ELKO, its **EIGENSPINORS** given in a  $\frac{1}{2};0$   $0;\frac{1}{2}$  representation (**INTRINSICALLY NEUTRAL**)

$$C^S(p) = +^S(p); \quad C^A(p) = -^A(p); \quad (2)$$

$$^S = \lim_{p \neq 0} +i \begin{bmatrix} (p) \\ (p) \end{bmatrix}; \quad ^A = \lim_{p \neq 0} -i \begin{bmatrix} (p) \\ (p) \end{bmatrix}; \quad (3)$$

with the right/left handed structures transforming as  $\frac{1}{2};0 = 0;\frac{1}{2}$

$$(\sigma \cdot \mathbf{p})(k) = (k); \quad (4)$$

$$\sigma \cdot \mathbf{p} (p) = (p_{\square}); \quad (5)$$

# Defining ELKO

- Defining parity as  $P = m^{-1} \not{p}$  and the time-reversal operator  $T = i \gamma_5 C$  yields [Phys. Reports 2205.04754v1 (2022)]  
 $(CPT)^2 = 1$ ; and  $fC;Pg = 0$  (6)

- ELKO transform differently than Dirac: (Possible non-Standard Wigner Class)

$$P \psi^{S=A}(p) = i m \psi^{S=A}(p) \quad (7)$$

$$T \psi^{S=A}(p) = i m \psi^{S=A}(p) \quad (8)$$

- ELKO: general non-vanishing momenta by boost  
 $\psi^{S=A}(p) = D(L(p)) \lim_{p \rightarrow 0} \psi^{S=A}(p)$ :

Field expansion prescription:

$$f(x) = \int \frac{d^3 p}{(2\pi)^3} \frac{1}{(2mE(\mathbf{p}))^{1/2}} \times \left[ a(\mathbf{p}) e^{-ip^\mu x_\mu} + b^\dagger(\mathbf{p}) e^{ip^\mu x_\mu} \right]$$

# Table of Contents

- 1 General remarks
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# Rotationally covariant basis, Dual, spin sums (Nuclear Physics B, Vol. 987 (2023) 116092)

This basis overcomes a recent criticism, successfully implementing FULL ROTATION COVARIANCE (in a Wigner anti-linear setting.)

$$u_1(p) = S_+(p) ; \quad u_2(p) = S(p) \quad (10)$$

$$u_3(p) = i A_+(p) ; \quad u_4(p) = i A(p); \quad (11)$$

and the anti-self-conjugate one has

$$\begin{aligned} u_1(p) &= A_+(p) ; & u_2(p) &= A(p) \\ u_3(p) &= i S_+(p) ; & u_4(p) &= i S(p); \end{aligned} \quad (12)$$

with both bases being defined in terms of the ELKO spinors.

The dual is defined as

$$\bar{u}_i(p) = P^{\dagger} u_i(p) ; \quad \bar{u}_i(p) = P^{\dagger} u_i(p) ; \quad (13)$$

with  $P = \frac{\not{p}}{m}$ .

# The Dual, spin sums

- Successfully implementing Well-defined expressions for products, spin sums, etc

$$u^{\dagger}(p) u(p) = m \quad v^{\dagger}(p) v(p); \quad (14)$$

with spin sums

$$\sum_{\lambda} u^{\dagger}(p) u(p) = 2m \quad \sum_{\lambda} v^{\dagger}(p) v(p); \quad (15)$$

with completeness relation

$$\sum_{\lambda} u^{\dagger}(p) u(p) + \sum_{\lambda} v^{\dagger}(p) v(p) = 4m;$$

- ELKO do not obey Dirac, just Klein-Gordon equation

$$p \cdot S(p) = im \cdot S(p); \quad p \cdot A(p) = im \cdot A(p); \quad (16)$$

# Lagrangian and $Z$ conjugation

- The model has the following Lagrangian ensuring KG equation

$$L_{\text{free}}(x) = \partial_\mu \dot{f}(x) \partial^\mu f(x) - m^2 \dot{f}(x) f(x) \quad (17)$$

It is invariant under  $Z$  conjugation (**generalized Hermitian conjugation**)

$L_{\text{free}} = L_{\text{free}}^Z$  **defined as**

$$i \int d^3p \, a_{S=A}(\mathbf{p}) e^{i p \cdot x} = i \int d^3p \, a_{S=A}(\mathbf{p}) e^{-i p \cdot x} \quad (18)$$

- The interaction lagrangian couples directly with physical fields being immediately gauge invariant (**photon interaction with neutral fermions!**)

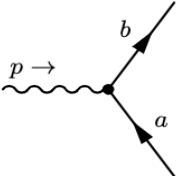
$$L = \frac{g}{2} i \bar{f} \not{F} f \quad (19)$$

With  $\frac{1}{2} [ \dots ]$  and  $F = \partial_\mu A_\nu - \partial_\nu A_\mu$



# Interaction part

## The interaction vertex



A Feynman diagram showing a vertex where a wavy line (representing a boson) with momentum  $p$  and an arrow pointing right enters from the left. Two solid lines (representing fermions) exit from the vertex: one labeled  $b$  with an arrow pointing up and to the right, and another labeled  $a$  with an arrow pointing down and to the right.

$$= igp_{ba} \quad (20)$$

The BOSONIC sector is composed of the Podolsky Lagrangian (No mixing gauge)

$$L_{\text{pod}} = \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2M^2} @ F_{\mu\nu} @ F^{\mu\nu} + \frac{1}{2} \dot{a}^2 + \frac{1}{2} @ A_{\mu} @ A^{\mu} \quad (21)$$

Electron-positron scattering with 12 GeV .  $p_{\bar{s}}$  . 46.8 GeV yields the bound  $M < 370$  GeV.

- The amplitude for the Compton scattering reads

$$\Gamma(k) = \not{\epsilon}(k) \not{k}$$

$$\mathcal{M}_{\gamma\gamma \rightarrow \gamma\gamma} = \frac{g^2}{2m_s m^2} i \text{Tr}[\not{\epsilon}(p^\ell) \Gamma(k^\ell) \Gamma(k) \not{\epsilon}(k^\ell) \not{\epsilon}(k) \not{p}^\ell] + \frac{g^2}{2m_u m^2} i \text{Tr}[\not{\epsilon}(p^\ell) \Gamma(k) \Gamma(k^\ell) \not{\epsilon}(k^\ell) \not{\epsilon}(k) \not{p}^\ell]: (22)$$

Replacing the  $s$  channel by  $t$ ,  $\not{\epsilon}(p^\ell)$  by  $\not{\epsilon}(p^\ell)$  and with all external photons being conjugated, we obtain the amplitude for the DARK MATTER pair annihilation

# Compton and pair annihilation non-polarized process

- In both  $y$  and  $z$  prescription, the probability is proportional to

$$\sum_{\lambda, \lambda'} M_{\lambda, \lambda'} M_{\lambda, \lambda'}^* \quad (23)$$

is proportional to a linear combination of the following terms

$$\text{Tr}[\Gamma(k)\Gamma(k^0)\Gamma(k^0)\Gamma(k)] + \text{Tr}[\Gamma(k^0)\Gamma(k)\Gamma(k^0)\Gamma(k)] \quad (24)$$

with both being proportional to  $(k \cdot k^0)(k^0 \cdot k) = 0$  In interacting case, the massive excitation becomes UNSTABLE just contributing to internal lines. The asymptotic spectrum is free of massive photons.

- This is our definition of darkness

# Darkness survive full quantum corrections?

From the path integral  $U(1)$  invariance  $A(x) = \theta \Lambda$ ,  $f(x) = 0$   
(Neutral.)

$$\begin{aligned}
 Z &= N \int \mathcal{D}A \mathcal{D}\bar{f} \mathcal{D}f \\
 &= \int \exp i \int d^4x \left[ \frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2M^2} (\partial_\mu F^\mu)^2 \right. \\
 &\quad \left. + \partial_\mu \bar{f} \partial^\mu f - m^2 \bar{f} f + \frac{ig}{2} \bar{f} \gamma_5 f (F_{\mu\nu}) \right. \\
 &\quad \left. + \left( \frac{1}{M^2} + 1 \right) (\partial_\mu A^\mu)^2 + J A + \bar{T} f + \bar{f} T \right] ;
 \end{aligned} \tag{25}$$

one can obtain the Ward-like identity

$$K \frac{\Gamma^3}{\bar{f}(p) f(\check{p}) A(k)} = 0 \tag{26}$$

# Darkness survive full quantum corrections?

The existence of the basis  $G_a = \{ \Gamma_4, \Gamma_5, \Gamma_6, \Gamma_7, \Gamma_8, \Gamma_9, \Gamma_{10}, \Gamma_{11}, \Gamma_{12}, \Gamma_{13}, \Gamma_{14}, \Gamma_{15}, \Gamma_{16}, \Gamma_{17}, \Gamma_{18}, \Gamma_{19}, \Gamma_{20}, \Gamma_{21}, \Gamma_{22}, \Gamma_{23}, \Gamma_{24}, \Gamma_{25}, \Gamma_{26}, \Gamma_{27}, \Gamma_{28}, \Gamma_{29}, \Gamma_{30}, \Gamma_{31}, \Gamma_{32}, \Gamma_{33}, \Gamma_{34}, \Gamma_{35}, \Gamma_{36}, \Gamma_{37}, \Gamma_{38}, \Gamma_{39}, \Gamma_{40}, \Gamma_{41}, \Gamma_{42}, \Gamma_{43}, \Gamma_{44}, \Gamma_{45}, \Gamma_{46}, \Gamma_{47}, \Gamma_{48}, \Gamma_{49}, \Gamma_{50}, \Gamma_{51}, \Gamma_{52}, \Gamma_{53}, \Gamma_{54}, \Gamma_{55}, \Gamma_{56}, \Gamma_{57}, \Gamma_{58}, \Gamma_{59}, \Gamma_{60}, \Gamma_{61}, \Gamma_{62}, \Gamma_{63}, \Gamma_{64}, \Gamma_{65}, \Gamma_{66}, \Gamma_{67}, \Gamma_{68}, \Gamma_{69}, \Gamma_{70}, \Gamma_{71}, \Gamma_{72}, \Gamma_{73}, \Gamma_{74}, \Gamma_{75}, \Gamma_{76}, \Gamma_{77}, \Gamma_{78}, \Gamma_{79}, \Gamma_{80}, \Gamma_{81}, \Gamma_{82}, \Gamma_{83}, \Gamma_{84}, \Gamma_{85}, \Gamma_{86}, \Gamma_{87}, \Gamma_{88}, \Gamma_{89}, \Gamma_{90}, \Gamma_{91}, \Gamma_{92}, \Gamma_{93}, \Gamma_{94}, \Gamma_{95}, \Gamma_{96}, \Gamma_{97}, \Gamma_{98}, \Gamma_{99}, \Gamma_{100} \}$  allows us to conclude (until 2-loops)

$$g\Gamma^{2 \text{ LOOP}}(p; \check{p}; k = p \quad \check{p}) = gF(p; \check{p}; k = p \quad \check{p}) \quad k : \quad (27)$$

This, together with Schwinger-Dyson equations ensuring the ELKO propagator structure, may prove **Darkness prevalence radiative corrections until  $g^6$  order in cross-section.**

- It is valid until  $g^6$  even when considering 4-point 1PI contributions and complete vertexes to Compton. However, box-box products violate the property. That is why it is valid just until  $g^6$  and not  $g^8$ .

# Darkness survives full quantum corrections?

- One loop contributions to Compton scattering. (The associated amplitude respects darkness for order  $g^6$ !)
- This is enough for phenomenological darkness (i.e. Standard QED measures 70 BOX events in hundred billion tentatives in LHC!... our coupling is  $10^5$  times smaller!)

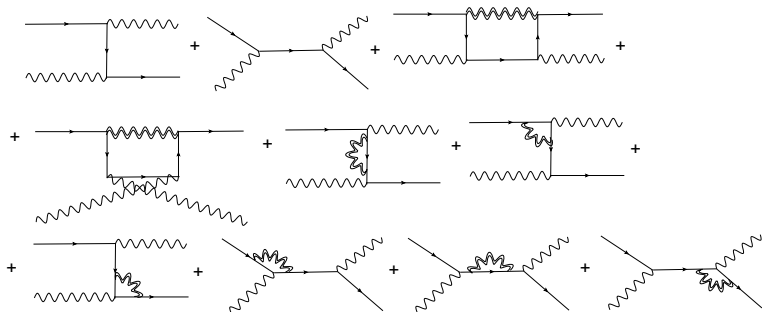


Figure: Compton radiative corrections.

# Nucleon Recoil Bounds

CONSIDERING...

- The effective **LOW ENERGY** vertex for photon-proton interaction is

$$V = \frac{1}{m^0} \quad p + \quad (28)$$

with  $m^0 = 1\text{GeV}$  denoting the proton mass.

- Using our machinery and the fact that the velocity  $232\text{Km/s}$  of the DM approaching the target made of protons
- The cross-section bound from LUX (for spin-dependent WIMP-proton interaction)  $10^{-39}\text{cm}^2$
- considering the very stringent limits, a WIMP with mass  $5\text{GeV}$  is our focus. Namely, (2206.06772) leads to an event rate peak at a possible WIMP with

$$M = 5\text{GeV} \quad (29)$$

WE CAN CALCULATE THE CROSS-SECTION  $f + p \rightarrow f + p$  (natural units.)

$$g = 10^{-5} \quad (30)$$