# Galactic cosmic rays (II)

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## Outline

- Motivation
- Ø Fundamental observations
  - Spectrum
  - Anisotropy
  - Composition
- 8 Key insights
  - Grammage
  - Cosmic ray clocks
  - Rigidity-dependence
  - Source candidates
- The transport equation
- 6 Exercises
- **6** Shock acceleration
  - Shocks
  - Macroscopic approach
  - Microscopic approach
  - Additional effects
- Galactic transport
  - Leaky box model
  - 1D model
  - Green's function
  - Numerical codes

- **③** Open question 1: The positron excess
  - Sources of positrons
  - Acceleration of secondaries
- Open question 2: Self-confinement
  - Gamma-ray haloes
  - Near-source transport
- Open question 3: A swiss-cheese galaxy
- Open question 4: The ionisation puzzle
- Open question 5: Small-scale anisotropies
  - Data
  - Test particle simulations
- Open question 6: Diffuse emission
  - Modelling
  - Results
  - 3D gas maps
- Summary & Conclusions

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## The transport equation

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#### Definitions

• We will consider the differential (in momentum p) particle density  $\psi$ ,

$$\psi(\mathbf{r},\boldsymbol{p},t) = 4\pi \boldsymbol{p}^2 f_0(\mathbf{r},\boldsymbol{p},t) = \frac{\mathrm{d}\boldsymbol{n}}{\mathrm{d}\boldsymbol{p}}$$

• Can also express this as differential in energy E,  $\tilde{\psi},$ 

$$\tilde{\psi}(\mathbf{r},\boldsymbol{\rho},t) = \frac{\mathrm{d}\boldsymbol{n}}{\mathrm{d}\boldsymbol{E}} = \frac{\mathrm{d}\boldsymbol{\rho}}{\mathrm{d}\boldsymbol{E}}\psi(\mathbf{r},\boldsymbol{\rho},t) = \frac{\boldsymbol{E}}{\boldsymbol{\rho}\boldsymbol{c}^2}\psi(\mathbf{r},\boldsymbol{\rho},t)$$

• What is measured is the intensity

$$J(\mathbf{r}, p, t) = \frac{\mathrm{d}^4 N}{\mathrm{d}A \,\mathrm{d}t \,\mathrm{d}\Omega \,\mathrm{d}E} = \frac{v}{4\pi} \tilde{\psi}(\mathbf{r}, p, t) = p^2 f_0(\mathbf{r}, p, t)$$

#### The transport equation

$$\begin{split} \frac{\partial \psi_j}{\partial t} = & \nabla \cdot \left( \kappa \cdot \nabla \psi_j - \mathbf{U} \psi_j \right) \\ &+ \frac{\partial}{\partial p} \left( p^2 D_{pp} \frac{\partial}{\partial p} \frac{1}{p^2} \psi_j \right) \\ &+ \frac{\partial}{\partial p} \left( -\frac{\mathrm{d}p}{\mathrm{d}t} \psi_j + \frac{p}{3} \left( \nabla \cdot \mathbf{U} \right) \psi_j \right) \\ &- v n_{gas} \sigma_j \psi_j - \frac{\psi_j}{\tau_j} \\ &+ v n_{gas} \sum_{k > j} \sigma_{k \to j} \psi_k + \sum_{k > j} \frac{\psi_k}{\tau_{k \to j}} \\ &+ S_j \end{split}$$

spatial diffusion and advection

momentum diffusion

momentum change incl. adiabatic

spallation and decay

spallation and decay

primary sources

## Mathematical description

Ginzburg & Syrovatskii (1964)



Blandford, Ostriker (1978)

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### CRs are a collisionless plasma

- From local measurement  $arepsilon_{CR} \sim 1\,eV\,cm^{-3} \Rightarrow \textit{n}_{CR} \sim 10^{-9}cm^{-3}$
- Inelastic collisions, typically  $\sigma \sim 100\,{\rm mb} = 10^{-25}\,{\rm cm}^2$

$$egin{aligned} &\Rightarrow\lambda=(\mathit{n_{\mathsf{CR}}}\sigma)^{-1}\ &=10^{34}\,\mathsf{cm}\simeq3 imes10^{12}\,\mathsf{kpc} \end{aligned}$$

$$\rightarrow$$
 (mean free path)  $\gg$  (size of the system)

 $\rightarrow\,$  CRs described by the collisionless Boltzmann equation:

$$\frac{\mathrm{d}f}{\mathrm{d}t} = \frac{\partial f}{\partial t} + \dot{\mathbf{x}} \cdot \nabla f + \dot{\mathbf{p}} \cdot \nabla_{\mathbf{p}} f = 0$$

## Diffusion

- But CRs interact with the magnetic field in the background plasma
- Homogeneous field  $\boldsymbol{B}_0 \to gyration$
- Magneto-hydrodynamic (MHD) equations allow for small-amplitude solutions  $\delta B$ : plasma waves



• Can perturbatively derive a diffusion equation:

$$rac{\partial \psi}{\partial t} - 
abla \cdot \left(\kappa \cdot 
abla - \mathbf{U}
ight)\psi = 0$$









If  $kr_g = 1 \implies$  coherent force over many gyrations



If  $kr_g \neq 1 \implies \mathbf{NO}$  coherent force

#### Diffusion

• Charged particles interact resonantly with wave:

$$k_{
m res} = r_g^{-1} \propto p^{-1}$$

• Scattering rate  $\Gamma$  is gyrofrequency  $\Omega$  reduced by density of resonant plasma waves:

$$\Gamma(k) \simeq \Omega rac{\delta ilde{B}^2(k)}{B_0^2} \simeq \Omega rac{k P(k)}{B_0^2}$$

• The power spectrum of resonant plasma waves is:

$$P(k) \sim \int \mathrm{d}^3 r \, \mathrm{e}^{i\mathbf{r}\cdot\mathbf{k}} \langle \delta \mathbf{B}(\mathbf{r}_0) \delta \mathbf{B}(\mathbf{r}_0 + \mathbf{r}) 
angle \propto k^{-5/3}$$

• Interaction with ensemble of random waves  $\rightarrow$  diffusion:

$$\kappa \simeq rac{v^2}{3\Gamma} \propto rac{1}{\Omega} \left( rac{k P(k)}{B_0^2} 
ight)^{-1} \Big|_{k=k_{
m res}} \propto p^{1/3}$$

#### Momentum diffusion

- Waves carry a momentum
- When scattering particles, can transfer momentum to CRs
- If angular distribution of waves is not isotropic, CRs get accelerated
- Stochastic process, can be described as diffusion in momentum space:

$$\frac{\partial f}{\partial t} - \frac{1}{p^2} \frac{\partial}{\partial p} \left( p^2 D_{pp} \frac{\partial f}{\partial p} \right) = 0$$
$$\frac{\partial \psi}{\partial t} - \frac{\partial}{\partial p} \left( p^2 D_{pp} \frac{\partial}{\partial p} \frac{1}{p^2} \psi \right) = 0$$

#### Momentum change

• CRs lose energy by interactions with radiation fields:

- Synchrotron emission
- Inverse Compton scattering
- . . . or gas:
  - Bremsstrahlung
  - Ionisation losses
  - Coulomb scattering
- Adiabatic losses or gains:

$$\left(\frac{\mathrm{d}\boldsymbol{p}}{\mathrm{d}t}\right)_{\mathsf{ad}} = -\frac{\boldsymbol{p}}{3}\left(\nabla\cdot\mathbf{U}\right)$$

• In transport equation:

$$\frac{\partial \psi}{\partial t} - \frac{\partial}{\partial p} \left( -\frac{\mathrm{d}p}{\mathrm{d}t} \psi + \frac{p}{3} \left( \nabla \cdot \mathbf{U} \right) \psi \right) = 0$$

## Spallation and decay



• In leaky box model, considered a global rate, e.g.

$$\frac{\mathrm{d}N_2}{\mathrm{d}t} = -\frac{N_2}{t_{\rm esc}} - N_2 \Gamma_2 + \Gamma_{1 \to 2} N_1$$

• In a diffusion model

$$rac{\mathrm{d}\psi}{\mathrm{d}t} \supset - extsf{vn}_{ extsf{gas}}\sigma_{j}\psi_{j} + extsf{vn}_{ extsf{gas}}\sum_{k>j}\sigma_{k
ightarrow j}\psi_{k}$$

• Additionally, decay:

$$rac{\mathrm{d}\psi}{\mathrm{d}t} \supset -rac{\psi_j}{ au_j} + \sum_{k>j} rac{\psi_k}{ au_{k
ightarrow j}}$$

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Exercise 1: Boron-to-carbon ratio

## Boron-to-carbon ratio (I)



 Solve a 1D, simplified version of the transport equation with homogeneous diffusion, advection, spallation, but no energy losses, radioactive decays or second-order Fermi acceleration, assuming steady-state,

$$-\kappa \frac{\mathrm{d}^2 \psi}{\mathrm{d}z^2} + V \frac{\mathrm{d}\psi}{\mathrm{d}z} = 2h\delta(z)(q - \Gamma\psi) \,.$$

- $\psi dE$  is the number density of CRs with kinetic energy in [E, E + dE]
- *κ* = *κ*(*E*) is the diffusion coefficient
- V is the advection speed away from the disk
- *h* is the disk height
- q = q(E) is the source spectrum
- $\Gamma = c\sigma_{tot}n_{gas}$  is the total destruction rate by inelastic collisions

#### Boron-to-carbon ratio (II)

**1** Show that for  $z \neq 0$ , the solution is of the form

$$\psi = A + B \mathrm{e}^{\frac{Vz}{\kappa}} \,,$$

with A, B functions of E.

**2** Integrate the transport equation across the Galactic disk and use the symmetry of the solution with respect to z = 0 to find the following relation between A and B,

$$-VB = qh - \Gamma h(A + B)$$
.

B Let *H* be the halo half-height. Employ the boundary condition at |z| = H, that is  $\psi(\pm H) \equiv 0$ , in order to express *A* through B,

$$A = -Be^{\frac{VH}{\kappa}}$$
.

4 Putting everything together, show that

$$\psi = \frac{-qh\left(e^{\frac{Vz}{\kappa}} - e^{\frac{VH}{\kappa}}\right)}{V - \Gamma h\left(1 - e^{\frac{VH}{\kappa}}\right)} = \dots = \underbrace{\frac{2qh}{V\left(1 + \coth\left(\frac{VH}{2\kappa}\right)\right) + 2\Gamma h}}_{\psi(z=0)} e^{\frac{Vz}{2\kappa}} \frac{\sinh\left(\frac{V(H-z)}{2\kappa}\right)}{\sinh\left(\frac{VH}{2\kappa}\right)}.$$

### Boron-to-carbon ratio (III)

S Assume one primary species, C, and one stable secondary species, B, and find the analytic expression for the boron-to-carbon ratio (B/C) in the disk,

$$rac{\psi_2}{\psi_1} = rac{2hc\sigma_{1
ightarrow 2}n_{ ext{gas}}}{V\left(1+ ext{coth}\left(rac{VH}{2\kappa}
ight)
ight)+2\Gamma_2h}$$

6 Adopting the following fixed paramters,

Quantity	Symbol	Value
Half-height of CR halo	Н	5 kpc
Half-height of disk	h	0.1 kpc
Gas density in disk	n <sub>gas</sub>	$1{ m cm^{-3}}$
Boron inelastic cross-section	$\sigma_{\rm tot,B}$	230 mbarn
Production cross-section for boron	$\sigma_{C \to B}$	150 mbarn <sup>a</sup>

write a python script to plot the predicted B/C as a function of energy-per-nucleon  $E_n$ . For the diffusion coefficient, assume the following behaviour

$$\kappa(E) = \kappa_0 \left(\frac{E_n}{1\,\text{GeV}}\right)^{\delta}$$

Typical values are  $\kappa_0=0.1\,{\rm kpc}^2\,{\rm Myr}^{-1}$ ,  $\delta=1/3$  and  $\,V=10^{-2}\,{\rm kpc}\,{\rm Myr}^{-1}.$ 

<sup>a</sup>Really more like 30 mb, but here with fudge factor of 5 to take into account other channels.

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## Boron-to-carbon ratio (IV)

- In python, write a code to plot AMS-02 data on B/C. To this end, pip-install the CRDB package,
- 1 \$ pip install crdb

and use the following python fragment:

```
1 import crdb
2 from crdb.mpl import draw_table, draw_logo
3 import matplotlib.pyplot as plt
4 import numpy as np
5
6 tab = crdb.query("B/C", energy_type="EKN")
7 ## crdb query returns a structured Numpy array:
8 ## the energies can be accessed as 't.e',
9 ## the B/C values as 't.value'
10 ## and the statistical errors as 't.err sta'
11 plt.figure()
12 mask = crdb.experiment_masks(tab)['AMS02']
13 t = tab[mask]
14 draw_table(t, label="AMS02")
15 plt.xlabel(r"$E_n$ [GeV]")
16 plt.ylabel("B/C")
17 plt.loglog()
18 draw_logo(0.78, 1)
19 plt.show()
```

## Boron-to-carbon ratio (V)



Overplotting the AMS-02 data with your model prediction for the typical parameter values, you should find

Now find the best-fit  $\kappa_0$ ,  $\delta$  and V by minimising the  $\chi^2$ .

## Fitting the boron-to-carbon ratio (VI)

- A couple of caveats:
  - Solar modulation
  - Isotopes
  - Actual cross-sections
  - Carbon is not a pure primary
  - Height of CR halo is also a free parameter

Any questions?