



Quantum-assisted interferometers for astrometry

Andrei Nomerotski

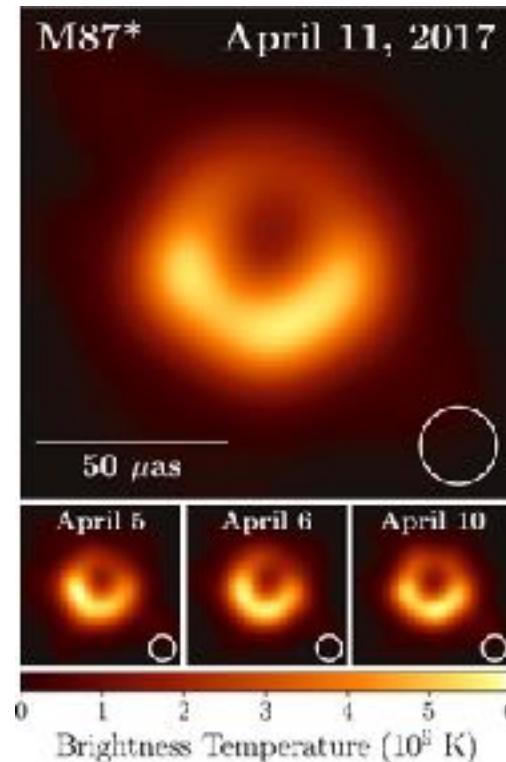
Florida International U &
Czech Technical U



CAPADS group in Czech TU



EHT black hole image



sensitive to features
on angular scale

$$\Delta \theta \sim \frac{\lambda}{b}$$

2019 ApJL 875

Black hole in the center of M87 imaged at 1.3mm

Achieved by radio interferometry with ~ 10000 km baselines

Radio

$$\bar{n} \gg 1$$

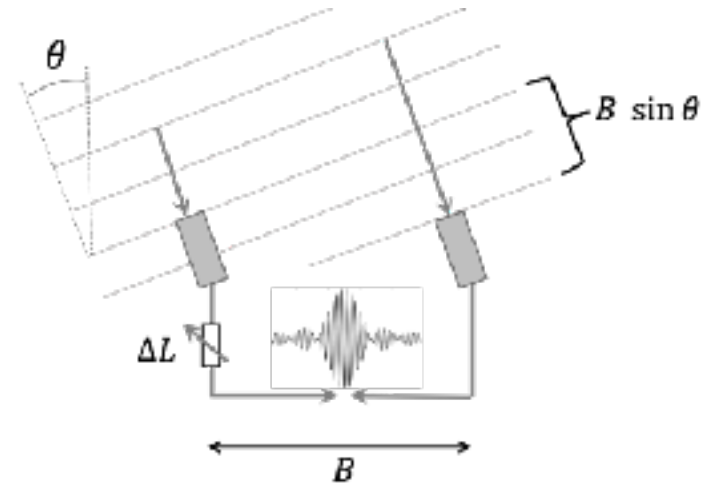
mode population

Can literally record entire waveform
interfere later offline

Optical

$$\bar{n} \ll 1$$

mode population



One photon at a time!

Need to bring paths to common point in real time

**Need path length delayed and
stabilized to better than λ**

Max baselines ~ 100 m

Two-photon techniques

Second photon for quantum assist

PRL 106, 070503 (2012)

PHYSICAL REVIEW LETTERS

week ending
17 AUGUST 2012

Longer-Baseline Telescopes Using Quantum Repeaters

Daniel Gottesman¹

Perimeter Institute for Theoretical Physics, Waterloo, Ontario, Canada

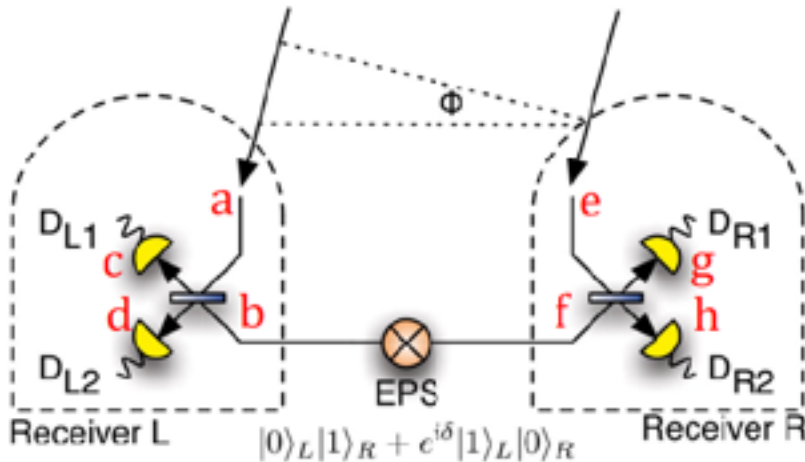
Thomas Jennewein²

Institute for Quantum Computing, University of Waterloo, Waterloo, Ontario, Canada

Sarah Croke¹

Perimeter Institute for Theoretical Physics, Waterloo, Ontario, Canada

(Received 25 October 2011; revised manuscript received 22 May 2012; published 16 August 2012)



$$\Psi^{\text{Initial}} = \psi_1 \psi_2 = \frac{1}{2} \underbrace{(\hat{a}^\dagger + e^{i\delta_1} \hat{e}^\dagger)}_{\text{Sky photon}} \underbrace{(\hat{b}^\dagger + e^{i\delta_2} \hat{f}^\dagger)}_{\text{Ground photon}}$$

Beam
Splitters

$$\begin{aligned} \hat{a}^\dagger &\rightarrow (\hat{c}^\dagger + \hat{d}^\dagger)/\sqrt{2} & \hat{b}^\dagger &\rightarrow (\hat{c}^\dagger - \hat{d}^\dagger)/\sqrt{2} \\ \hat{e}^\dagger &\rightarrow (\hat{g}^\dagger + \hat{h}^\dagger)/\sqrt{2} & \hat{f}^\dagger &\rightarrow (\hat{g}^\dagger - \hat{h}^\dagger)/\sqrt{2} \end{aligned}$$

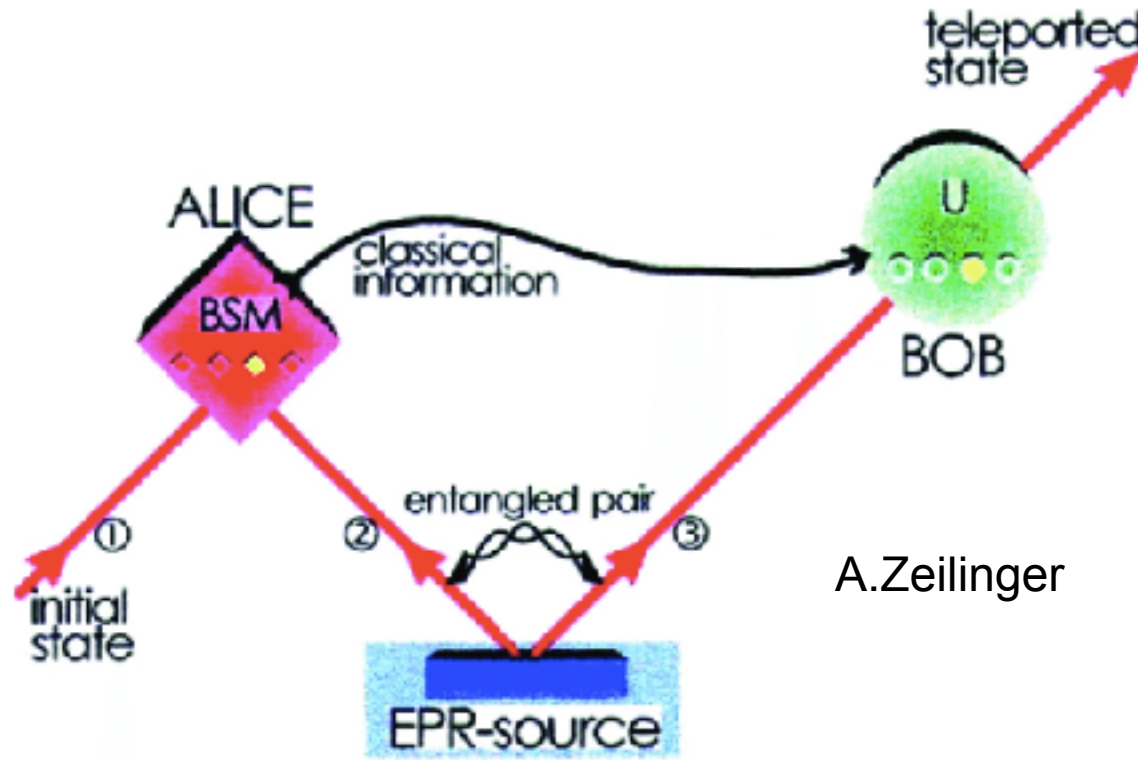
$$\Psi^{\text{Output}} = (1/4)(\hat{c}^\dagger \hat{c}^\dagger - \hat{d}^\dagger \hat{d}^\dagger + e^{i(\delta_1 + \delta_2)}(\hat{g}^\dagger \hat{g}^\dagger - \hat{h}^\dagger \hat{h}^\dagger) + (e^{i\delta_1} + e^{i\delta_2})(\hat{c}^\dagger \hat{g}^\dagger - \hat{d}^\dagger \hat{h}^\dagger) + (e^{i\delta_1} - e^{i\delta_2})(\hat{c}^\dagger \hat{h}^\dagger + \hat{d}^\dagger \hat{g}^\dagger))$$

$$\begin{aligned} P(c^2) = P(d^2) = P(g^2) = P(h^2) &= 1/8 \\ P(cg) = P(dh) &= (1/8)(1 + \cos(\delta_1 - \delta_2)) \\ P(ch) = P(dg) &= (1/8)(1 - \cos(\delta_1 - \delta_2)) \end{aligned}$$

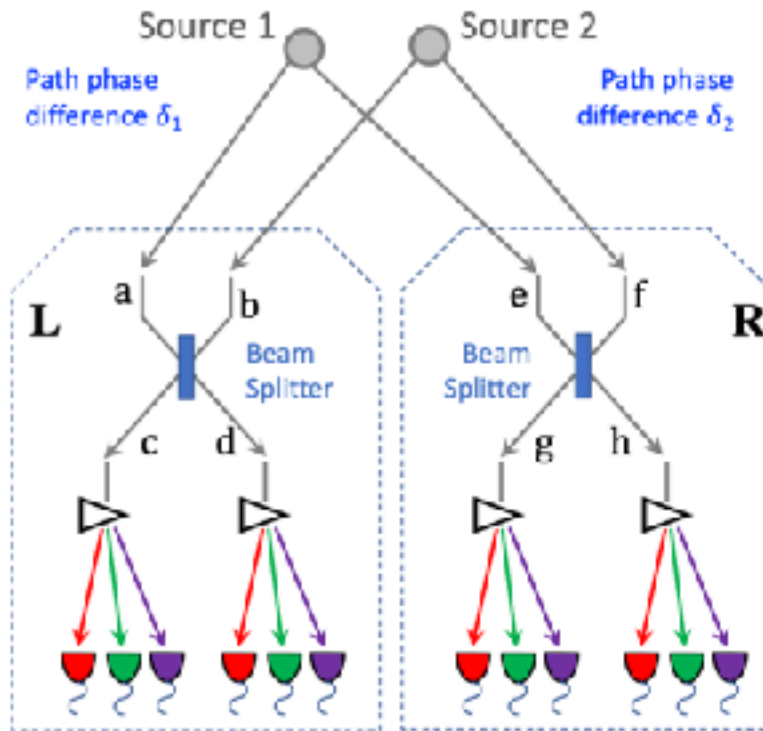
- Correlation of counters are sensitive to sky photon phase → direction
- Need to transfer the photon quantum state → can use quantum networks, this will allow long distances
- Verified experimentally in quantum optics settings: M Brown et al, PRL 131, 210801 (2023)

Quantum Network

- Attenuation in fibers → need quantum repeater to reproduce qubits
 - Simple amplification will not conserve the quantum state
- Qubit teleportation: produce entangled photons and send them to two locations
- Bell State Measurement (BSM) on one photon will collapse the wave function of the other one (or swap entanglement, or teleport photon)



Idea: use another star as source of coherent states for the interference



$$\begin{aligned}
 P(c^2) &= P(d^2) = P(g^2) = P(h^2) = 1/8 \\
 P(cg) &= P(dh) = (1/8)(1 + \cos(\delta_1 - \delta_2)) \\
 P(ch) &= P(dg) = (1/8)(1 - \cos(\delta_1 - \delta_2))
 \end{aligned}$$

Full QFT calculation

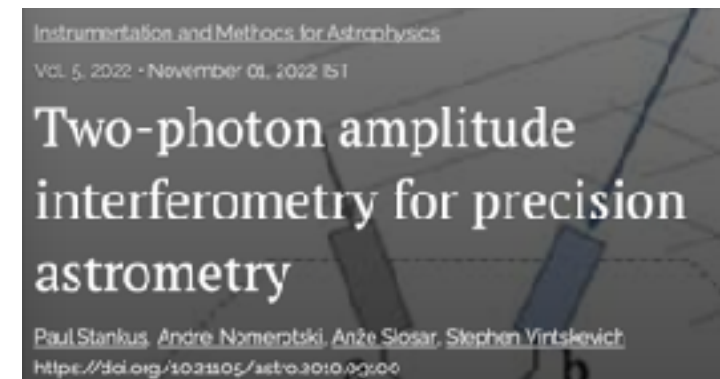
$$\begin{aligned}
 N_c(xy) &= \eta_1 \eta_2 A^2 \int_0^{T_r} P_{L,R,\tau}^{\text{two photons}} d\tau = \\
 &A^2 \eta_1 \eta_2 T_r \left[\underbrace{(I_1 + I_2)^2}_{\text{Rates}} + \underbrace{I_1^2 \frac{\tau_c g_{11}}{T_r} + I_2^2 \frac{\tau_c g_{22}}{T_r}}_{\text{HBT}} \pm \right. \\
 &\left. 2I_1 I_2 \frac{\tau_c g_{12}}{T_r} \cos \left(\frac{\omega_0 B (\sin \theta_1 - \sin \theta_2)}{c} + \frac{\omega_0 \Delta L}{c} \right) \right] \quad (30)
 \end{aligned}$$

New oscillatory term!

Relative path phase difference $\delta_1 - \delta_2$ can be extracted from the coincidence rates of four single photon counters: c, d, g and h

Perfect to start exploring this approach

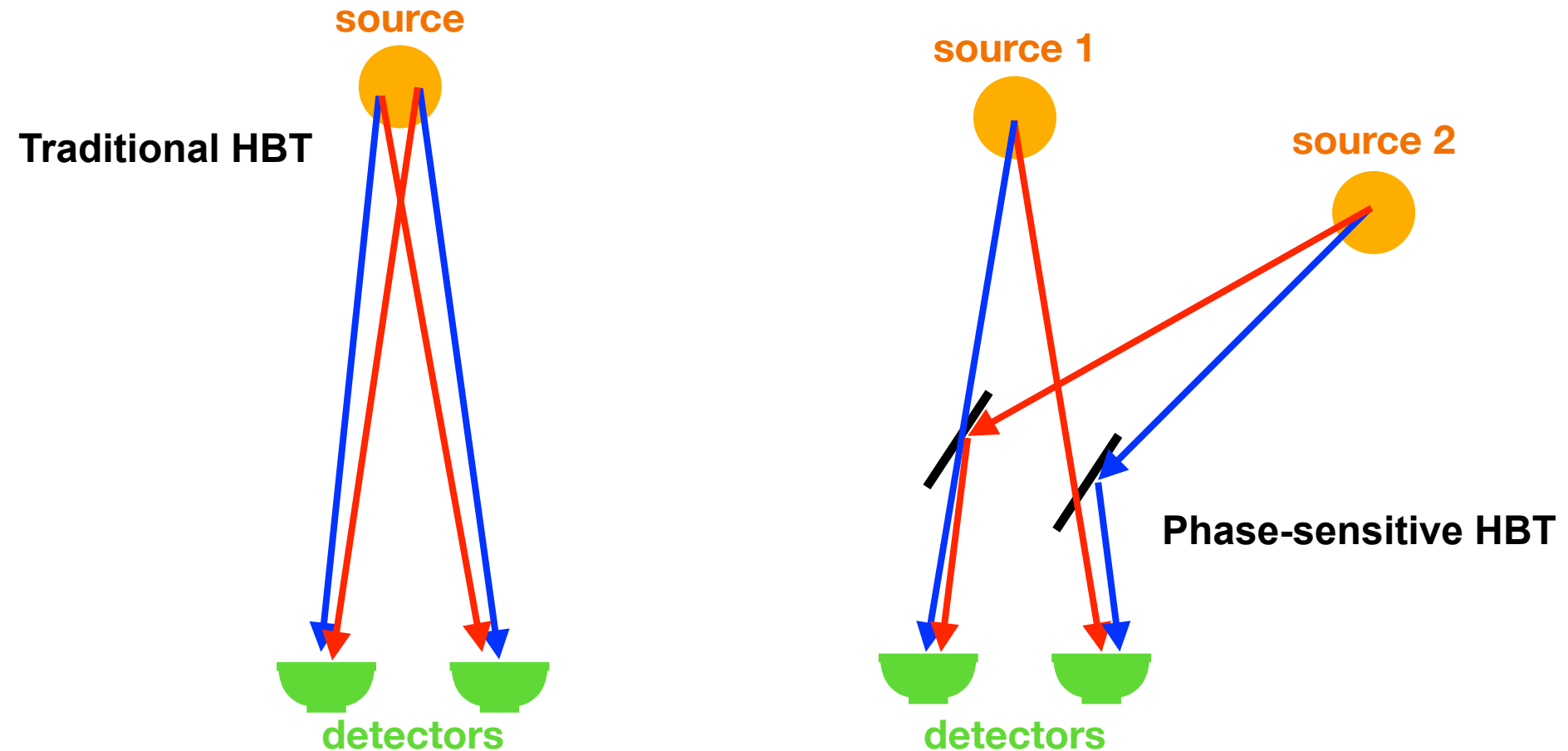
<https://arxiv.org/abs/2010.09100>



Hanbury Brown – Twiss Interferometry

If points are close enough two options of photons paths are coherent = photon phases not so different and they interfere

Interference produces photon bunching or HBT effect



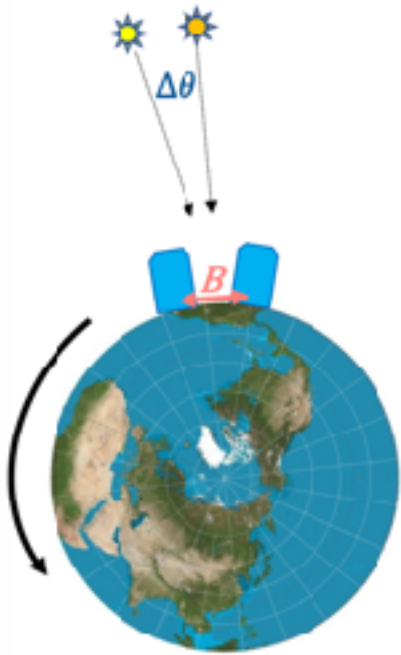
Earth rotation fringe scan

$$\langle N(xy) \rangle = \frac{k(S_1 + S_2)^2}{8} \left[1 \pm V_{2PS} \cos \left[\frac{2\pi B}{\lambda} (\sin \theta_1 - \sin \theta_2) + \frac{2\pi \Delta L}{\lambda} \right] \right]$$

This will evolve as the Earth rotates

$$\langle N_{xy} \rangle(t) = \bar{N}_{xy} [1 \pm V \cos(\omega_f t + \Phi)]$$

Coincidence rates oscillate

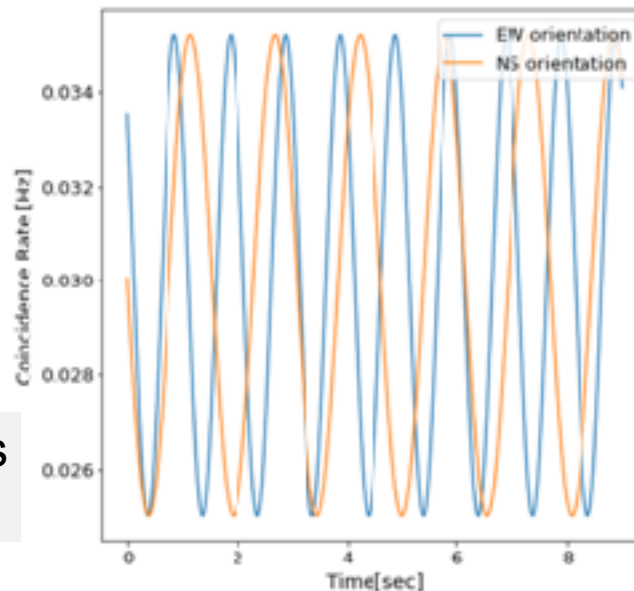


$$\omega_f = \frac{2\pi B \Omega_{\oplus} \sin \theta_0}{\lambda} \Delta \theta$$



Fringe oscillation rate is a direct measure of sources' opening angle!

Can measure with high precision



example of oscillations for pair of stars

$$\sigma[\Delta \theta] \sim 10 \mu\text{as} (\sim 10^{-11} \text{ rad})$$

Related Ideas

Astrometry with Extended-Path Intensity Correlation

Ken Van Tilburg,^{1,2,*} Masha Baryakhtar,^{3,†} Marios Galanis,^{4,‡} and Neal Weiner^{1,§}

¹*Center for Cosmology and Particle Physics, Department of Physics,
New York University, New York, NY 10003, USA*

²*Center for Computational Astrophysics, Flatiron Institute, New York, NY 10010, USA*

³*Department of Physics, University of Washington, Seattle WA 98195, USA*

⁴*Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada*

(Dated: July 10, 2023)

arxiv.org/abs/2307.03221

Tang et al. *Light: Science & Applications* (2025)14:46
<https://doi.org/10.1038/s41377-024-01684-y>

Official journal of the CIOMP 2047-7538
www.nature.com/lsa

ARTICLE

Open Access

Phase-dependent Hanbury-Brown and Twiss effect for the complete measurement of the complex coherence function

Xuan Tang¹, Yunxiao Zhang², Xueshi Guo², Liang Cui², Xiaoying Li²  and Z. Y. Ou¹ 

Possible impact on astrophysics and cosmology

<https://arxiv.org/abs/2010.09100>

offers orders of magnitude better astrometry with major impact

- Parallax: improved distance ladder
- Proper star motions - sensitive to Dark Matter
- Black hole accretion disk imaging
- **Gravitational waves through coherent motions of stars - microHz range**
- Exoplanets, microlensing etc

Requirements for detectors

Photons must be indistinguishable so close enough in frequency and time to interfere → temporal & spectral binning :

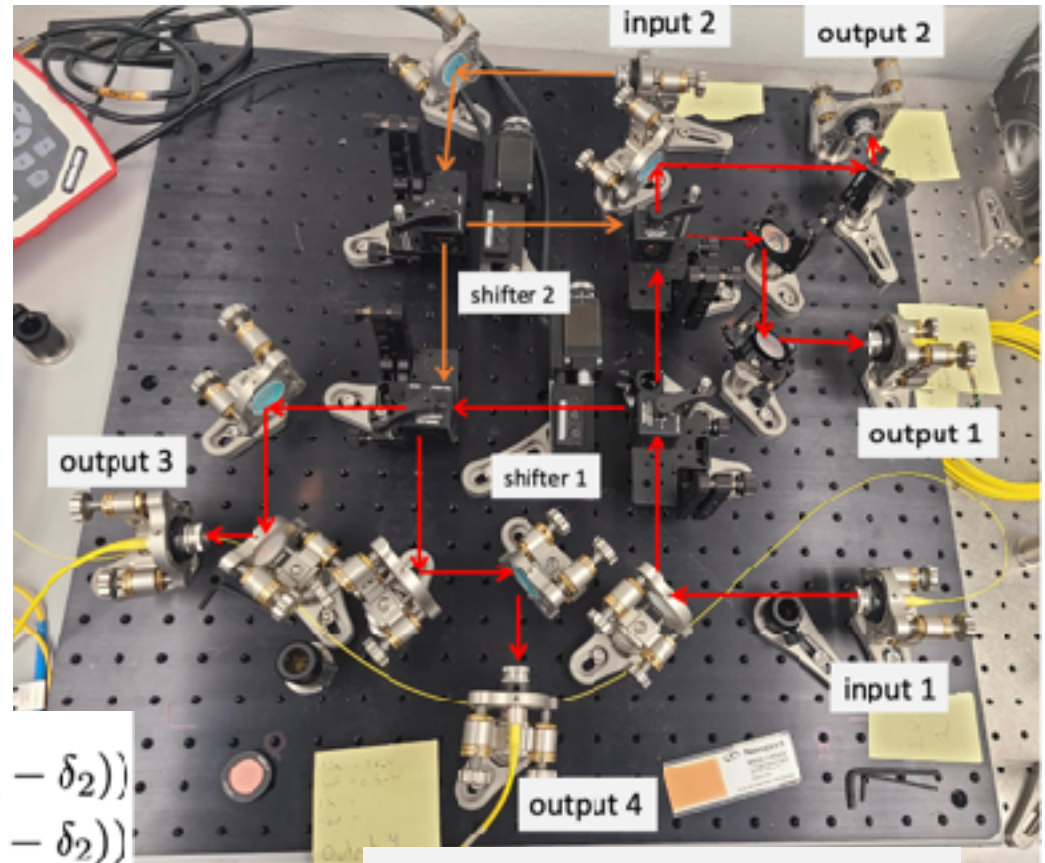
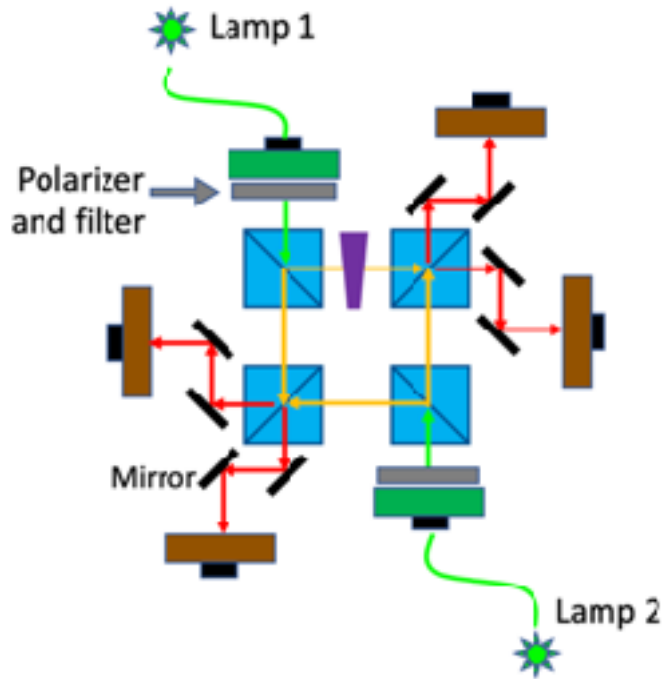
need $\sim 10 \text{ pm} * 10 \text{ ps}$ $\Delta t * \Delta E \geq \hbar/2$

**Need fast spectrometers
at Heisenberg Uncertainty Principle limit
(or Fourier limit, or Heisenberg-Gabor limit)**

then achieve maximum contrast for HBT peaks, no resolution smearing

- Fast imaging techniques are the key
 - Promising technologies: **SPADs**, SNSPDs
 - Target 10 - 100 ps resolution
- Spectral binning: diffraction gratings, echelle spectrometers
- High photon detection efficiency

Benchtop Verification

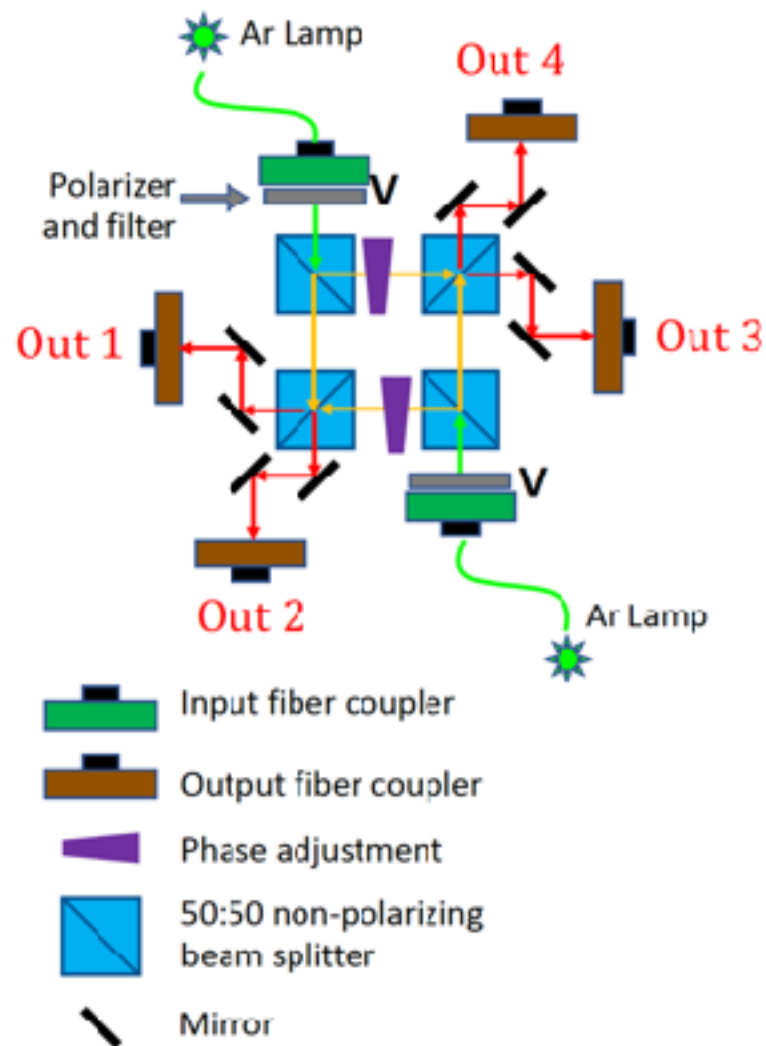
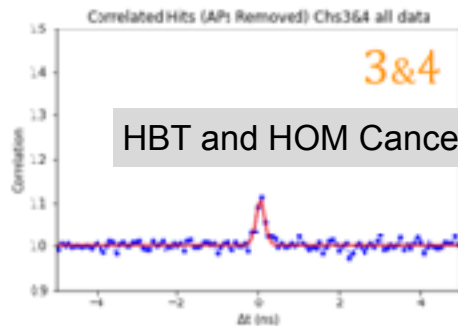
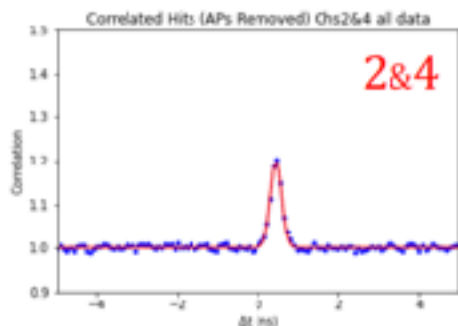
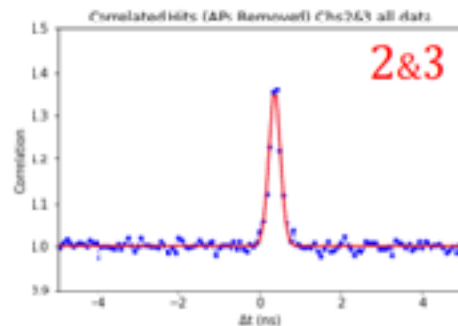
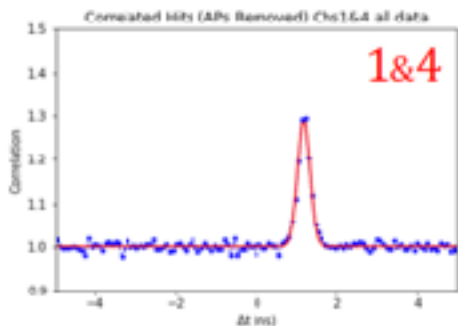
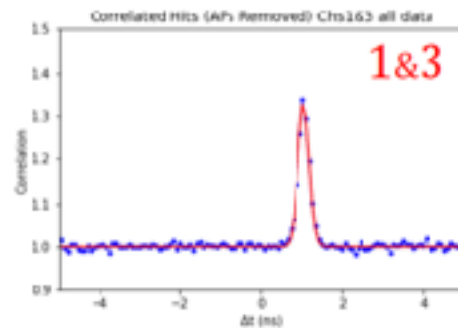
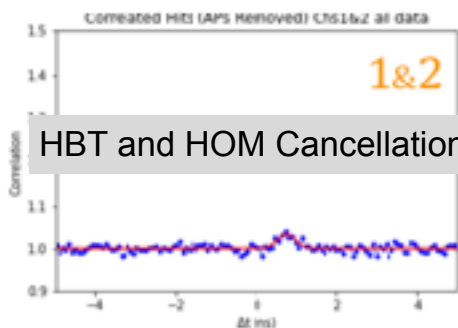


$$P(cg) = P(dh) = (1/8)(1 + \cos(\delta_1 - \delta_2))$$

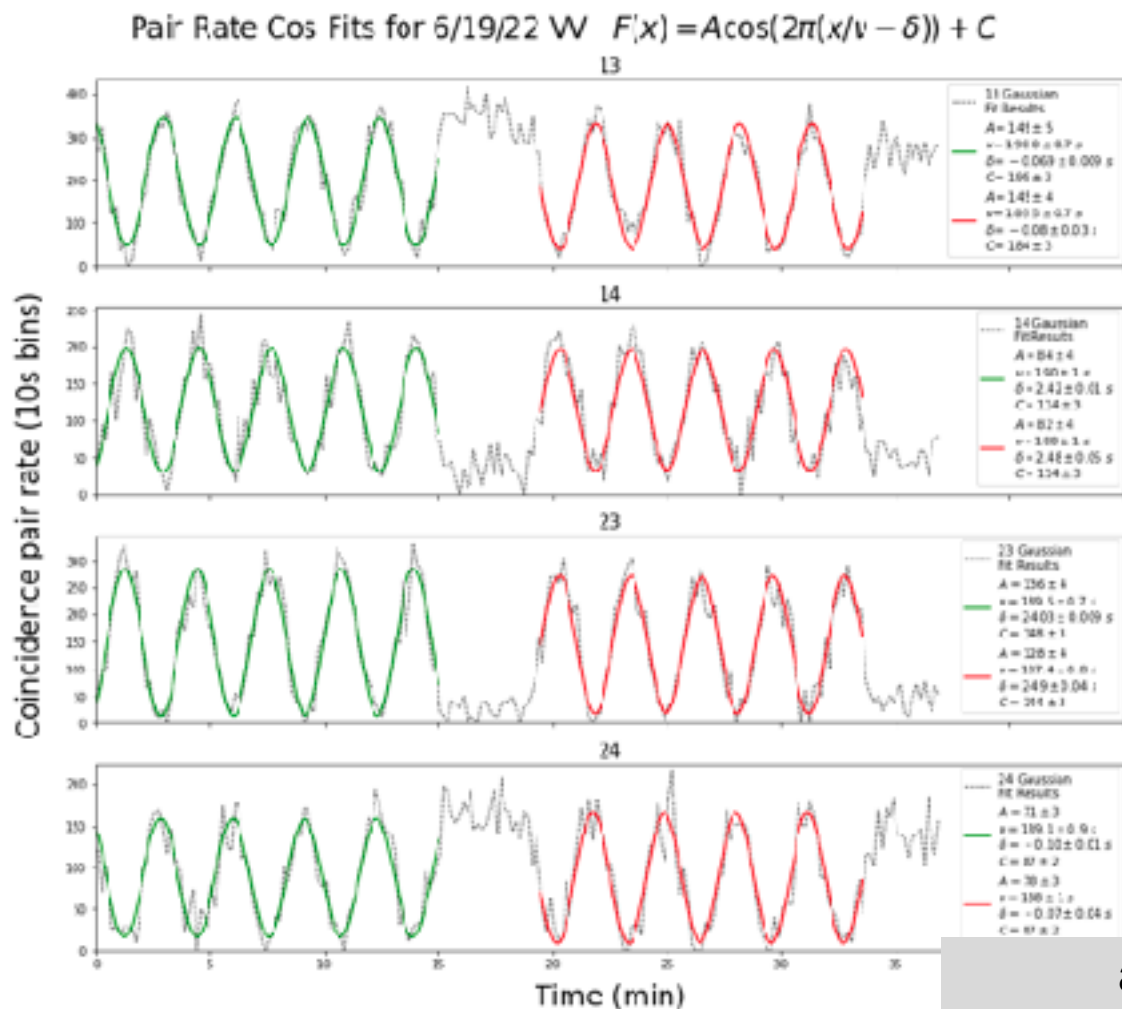
$$P(ch) = P(dg) = (1/8)(1 - \cos(\delta_1 - \delta_2))$$

SPAD and SNSPD readout

Polarized – V V



Phase dependence



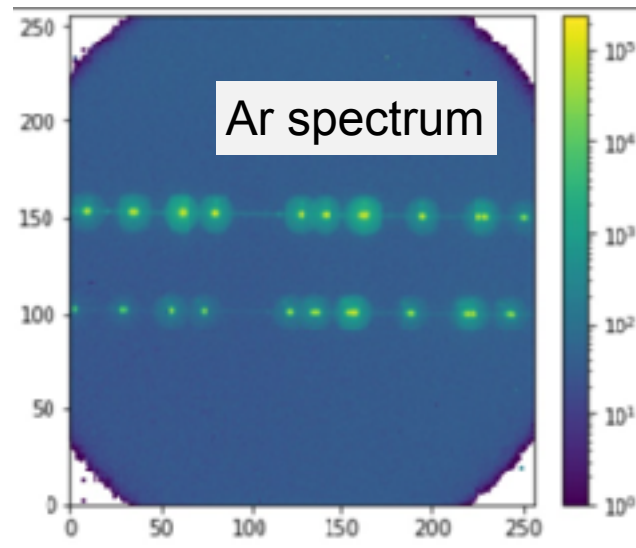
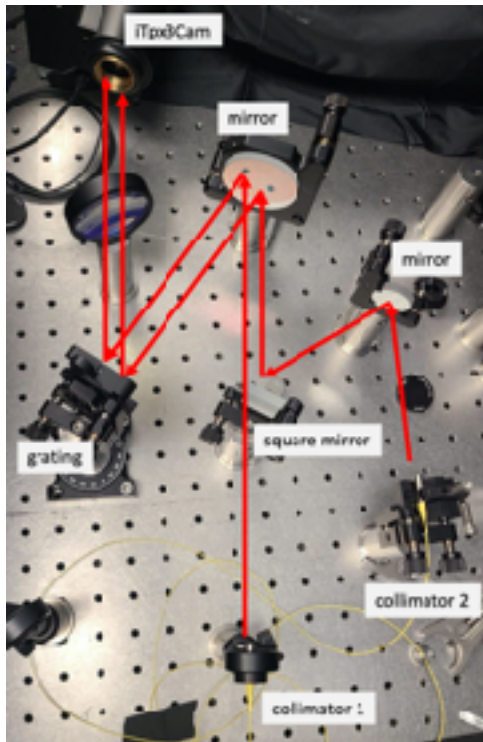
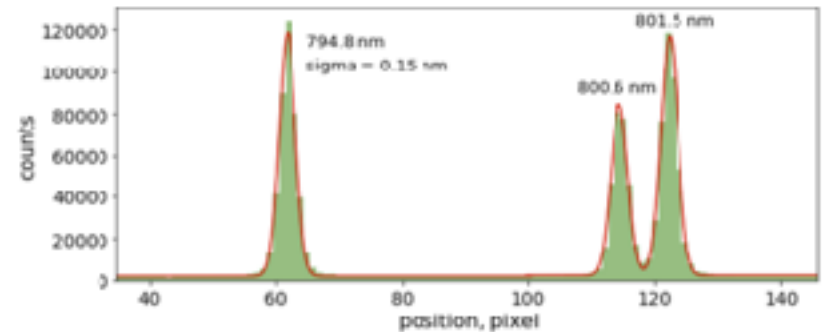
arxiv.org/abs/2301.07042
Optics Express 31, 44246-44258 (2023)

Population of HBT peaks as function of phase = phase oscillations

Next step: spectral binning

Spectral binning: spectrometer

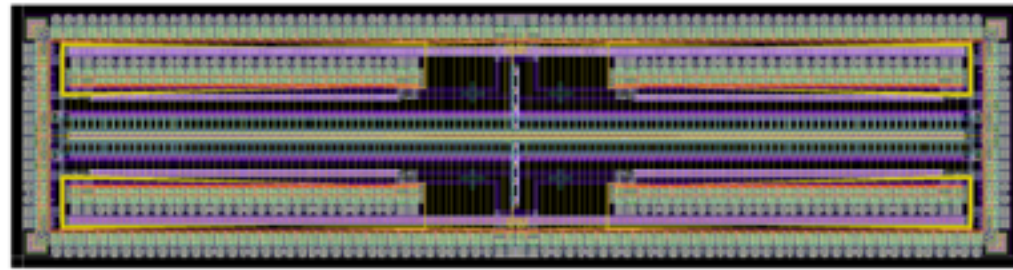
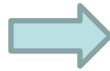
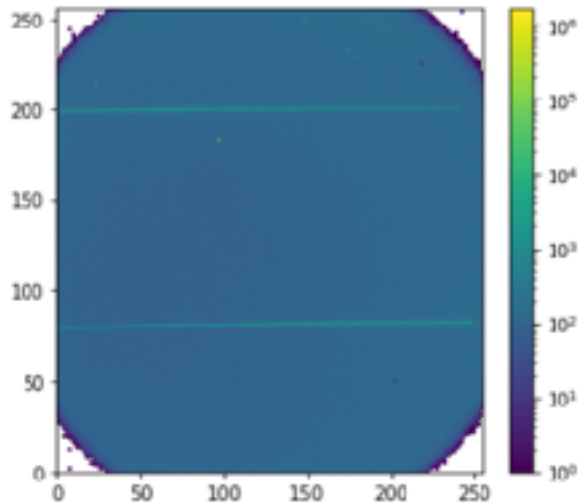
Two beams of thermal photons $\rightarrow$ diffraction grating
Based on intensified Tpx3Cam, ns time resolution



A.Nomerotski et al. Intensified Tpx3Cam, a fast data-driven optical camera with nanosecond timing resolution for single photon detection in quantum applications, arxiv.org/abs/2210.13713, published in JINST

Next steps: spectrometer based on LinoSPAD2

Diffracted photon stripe projected on to linear array



Spectrometer time resolution: 5 ns $\rightarrow$ 100 ps

Possible technologies: SPAD

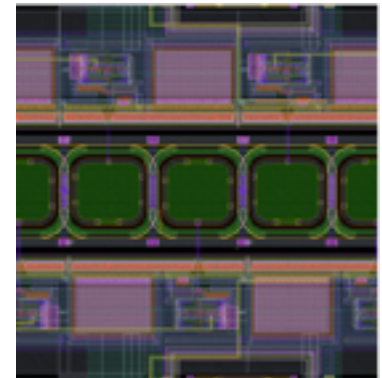
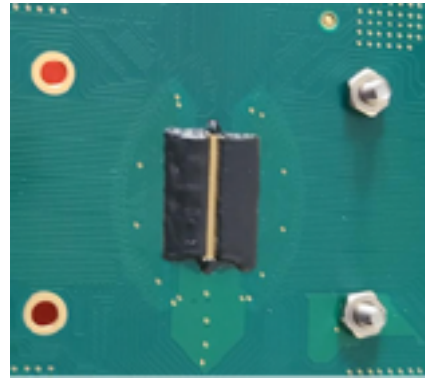
SPAD = single photon avalanche device

Semiconductor device: p-n junction with amplification

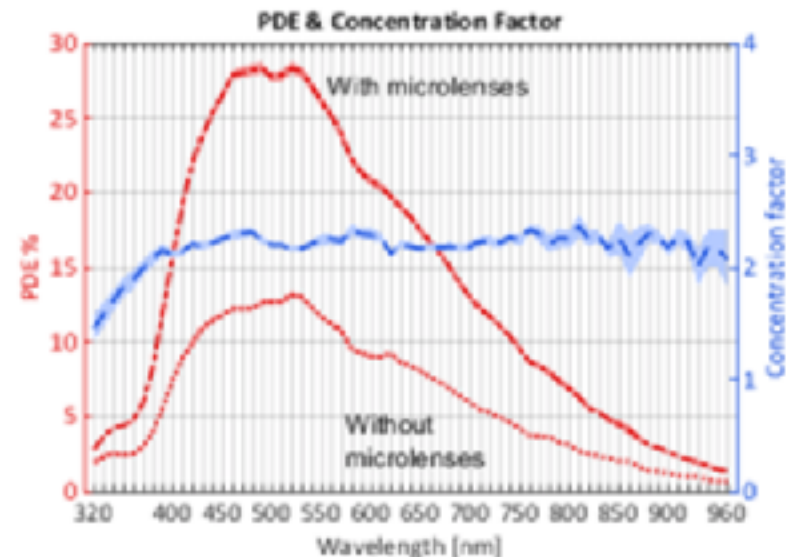
- 512 x 1 pixels
- 24 x 24 micron pixels
- Max PDE (with microlenses) ~ 30%
- 50 ps resolution

Developed in EPFL (Switzerland)
AQUA group (E.Charbon)

LinoSPAD2

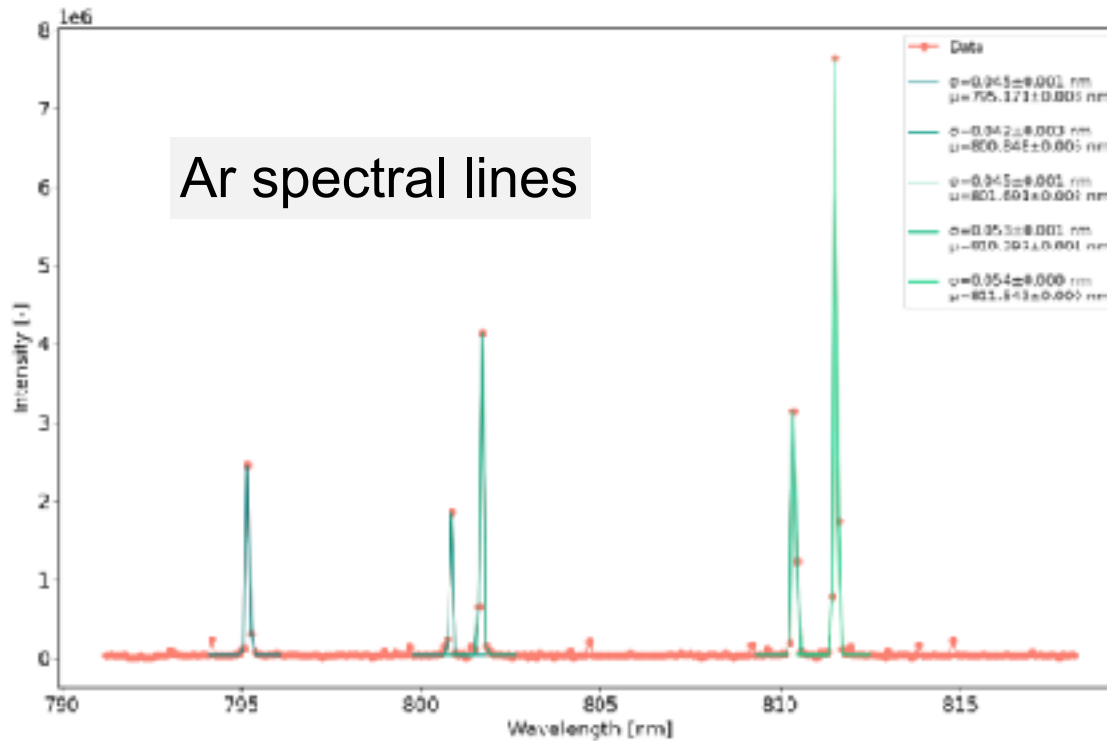


Close-up of SPADs



C. Bruschini et al, Linospad2: a 512x1 linear spad camera, in Quantum Sensing and Nano Electronics and Photonics XIX, vol. 12430, pp. 126–135, SPIE, 2023.

Spectrometer with LinoSPAD2



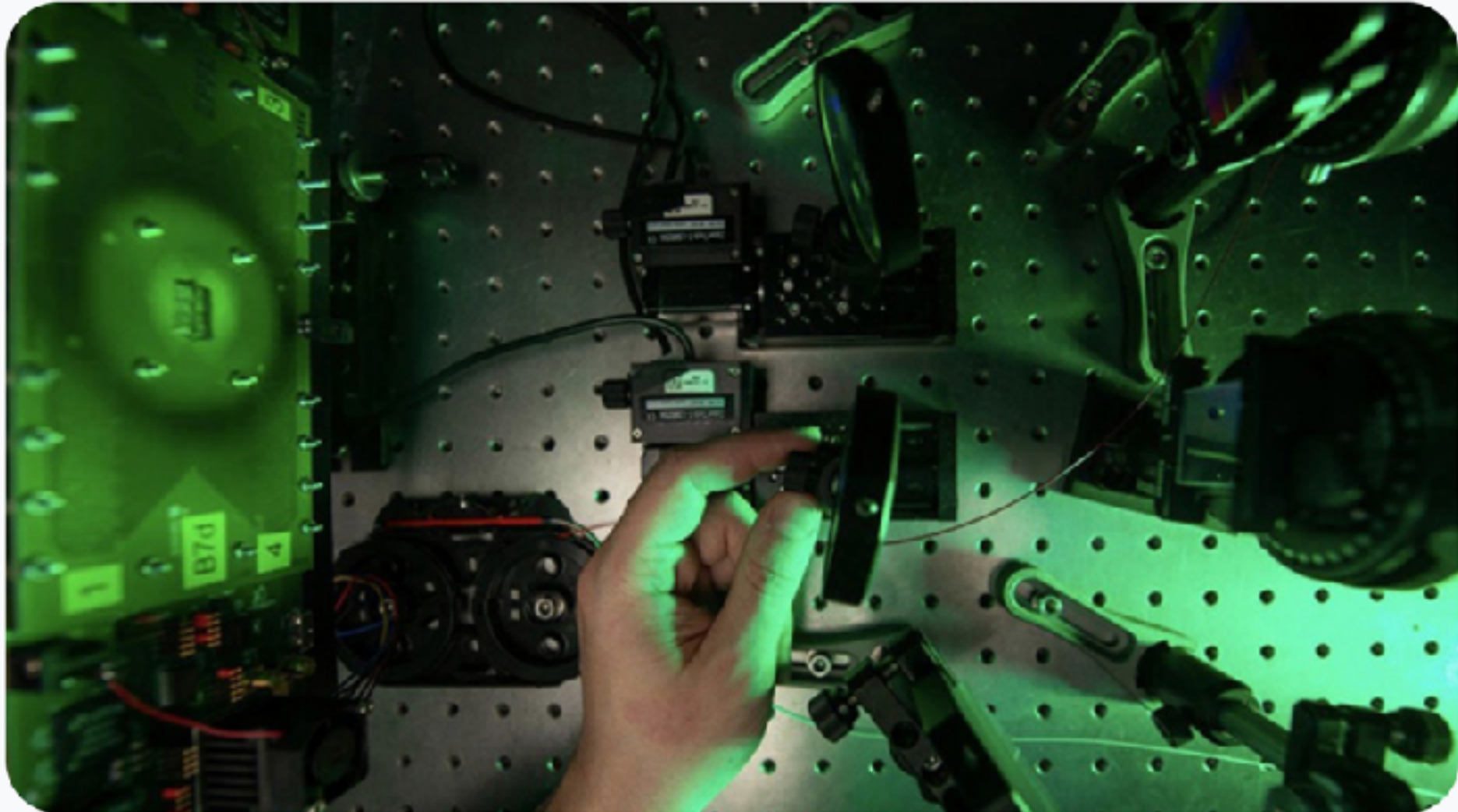
arxiv.org/abs/2304.11999

spectral resolution 0.04 nm

21

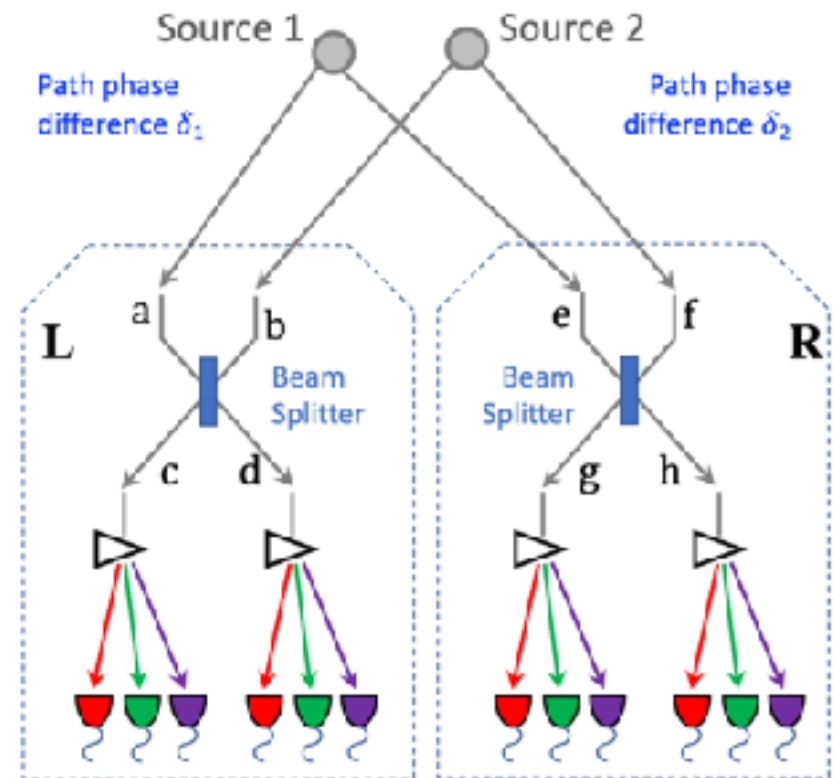
Achieved 0.04 nm spectral and 40 ps timing resolution
only x10 more than $\Delta t * \Delta E \geq \hbar/2$

Next talk



Important points for discussion

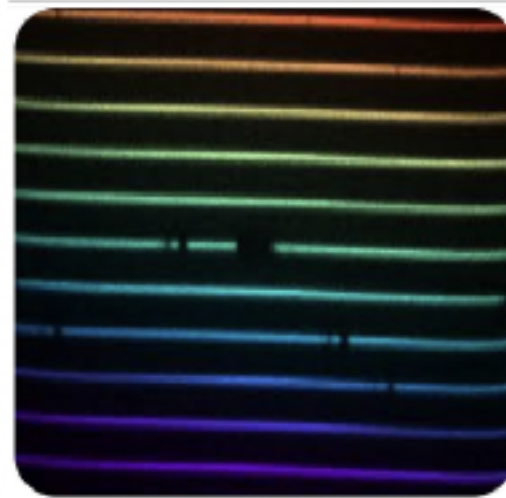
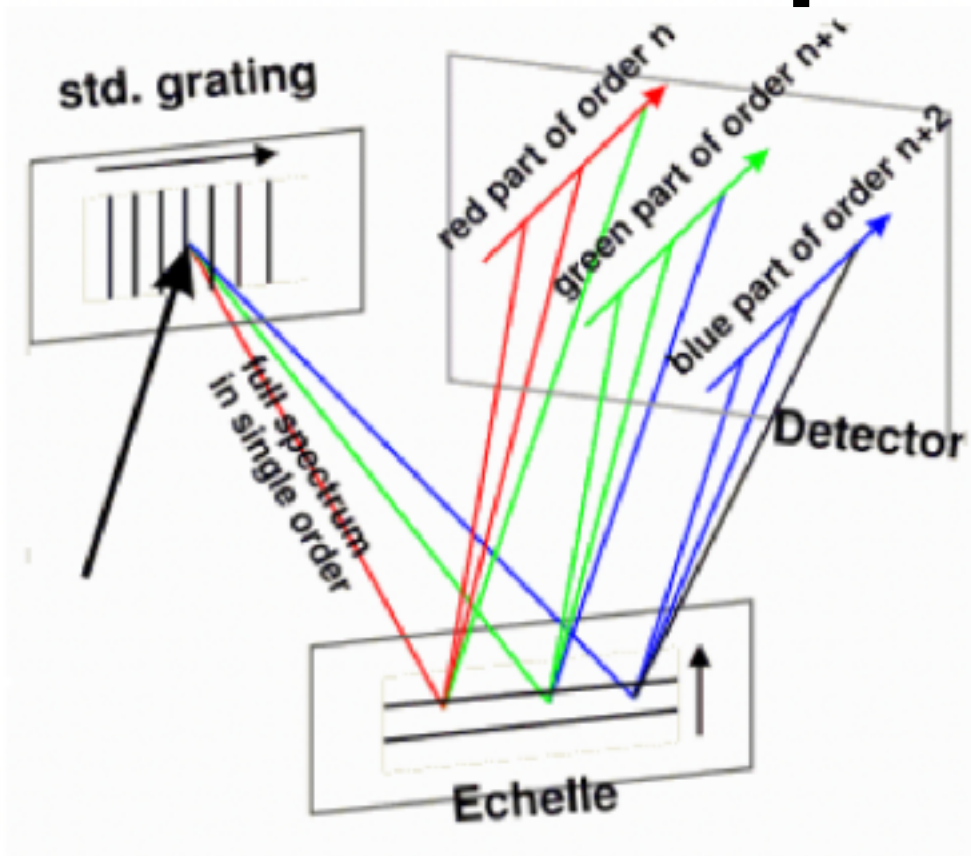
- For amplitude two-photon interferometry **light needs be coupled to SMF**
- **This is difficult!**
5-10 micron spot, non-trivial adaptive optics
- When achieved then interference and spectrometers are easy
- Work in progress to adapt spectrometer to MMF



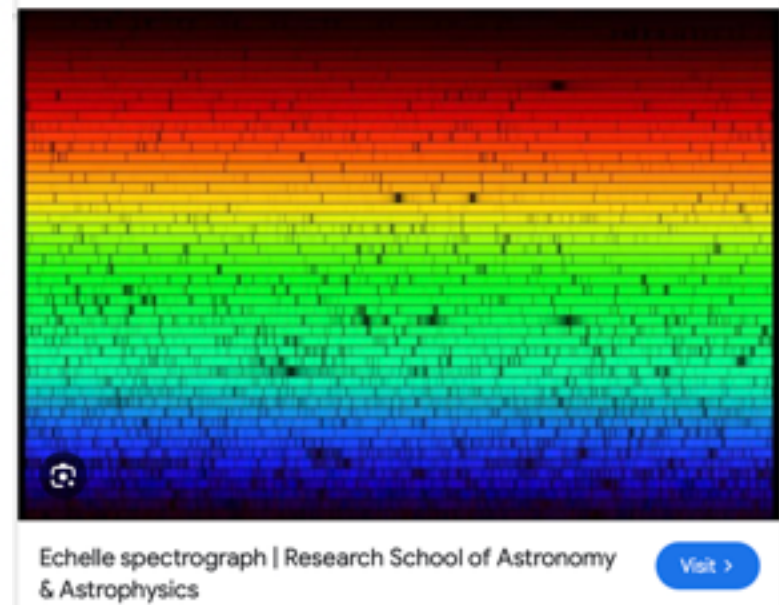
Sensor R&D

Ideas for 2d imaging sensors which can provide 20 ps resolution

Echelle spectrometer

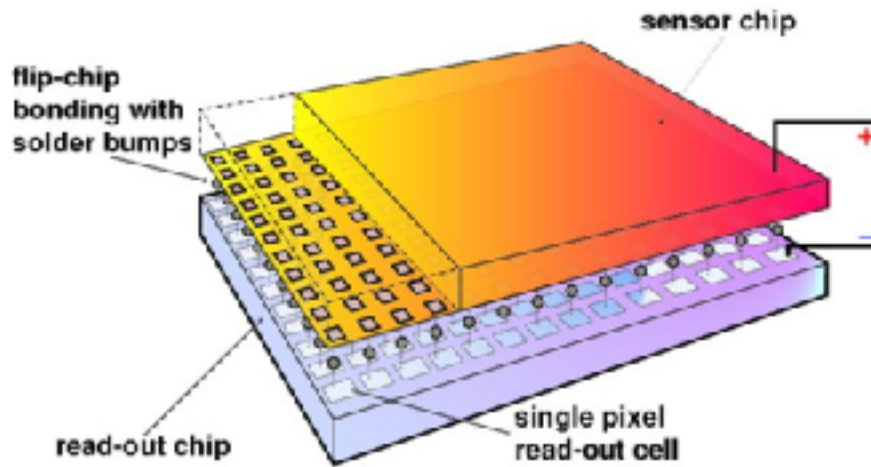


- Widely used in astronomy
- Achievable resolution $\sim$ few pm
- **Require 2D single photon detector**

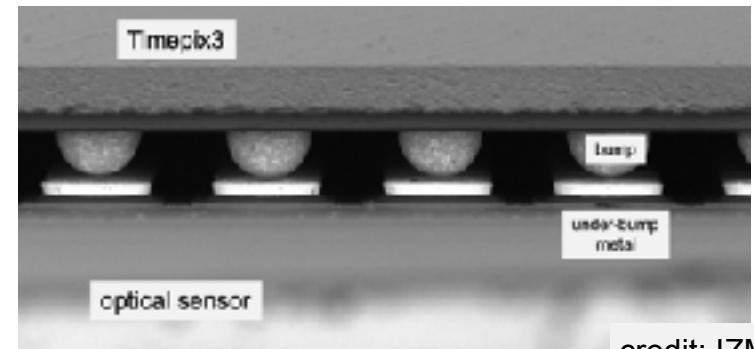


Fast hybrid pixel detectors

Have roots in R&D for LEP/LHC vertex detectors



Lukas Tlustos and Erik H. M. Heijne, Performance and limitations of high granularity single photon processing X-ray imaging detectors, in CERN proceedings (2005)



credit: IZM

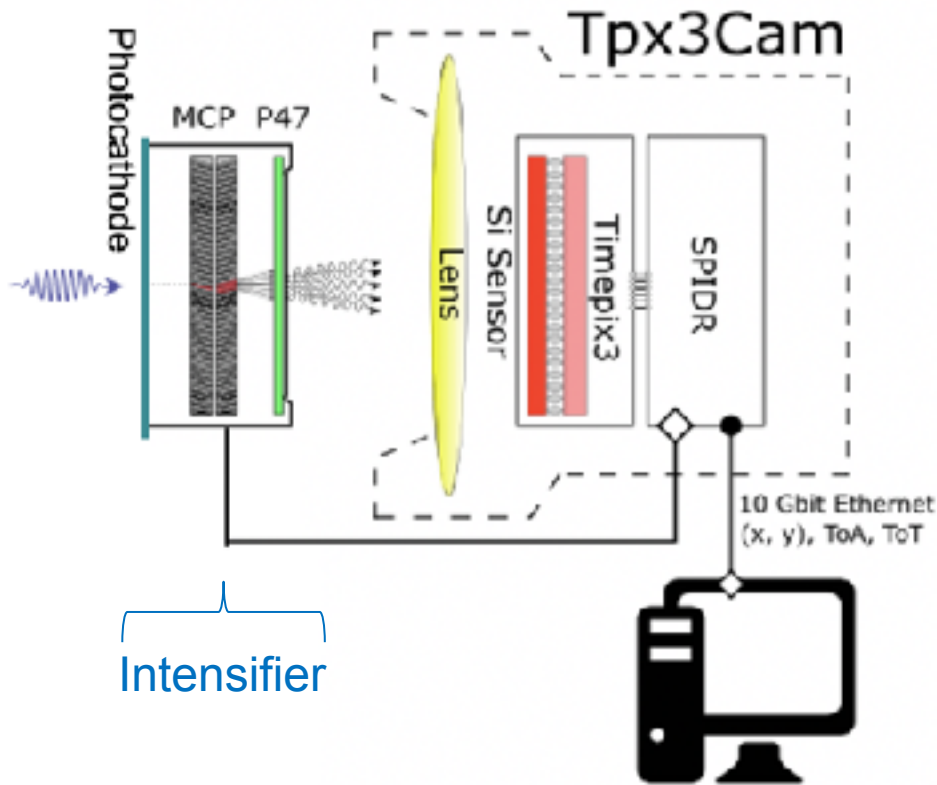
Decouple readout chip and sensor

- optimize technologies for chip and sensor separately

Use different sensors with same readout, versatile approach for x-rays (Si, CZT)

→ we will use **OPTICAL** sensors

Intensified fast camera: use off-the-shelf image intensifier



Timepix3 ASIC:

- 256 x 256 array, 55 x 55 micron pixel
 - 14 mm x 14 mm active area
- **1.56 ns timing resolution**
- **Data-driven readout**, 600 e min threshold

Optical Si sensor + Timepix3

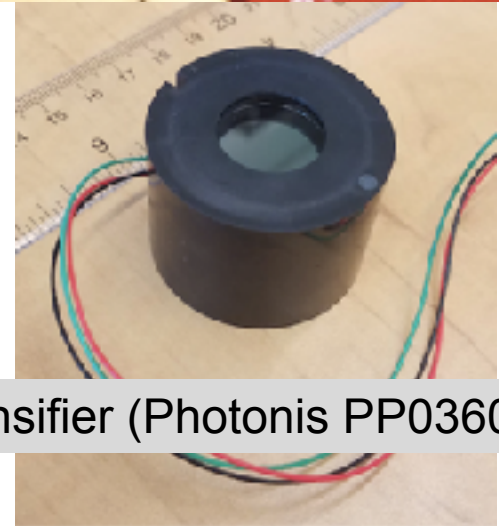
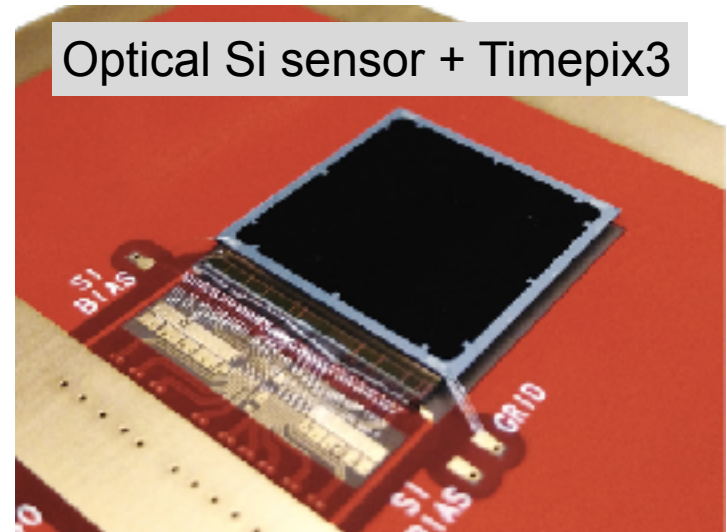


Image intensifier (Photonis PP0360EG)

Timepix3 → Timepix4

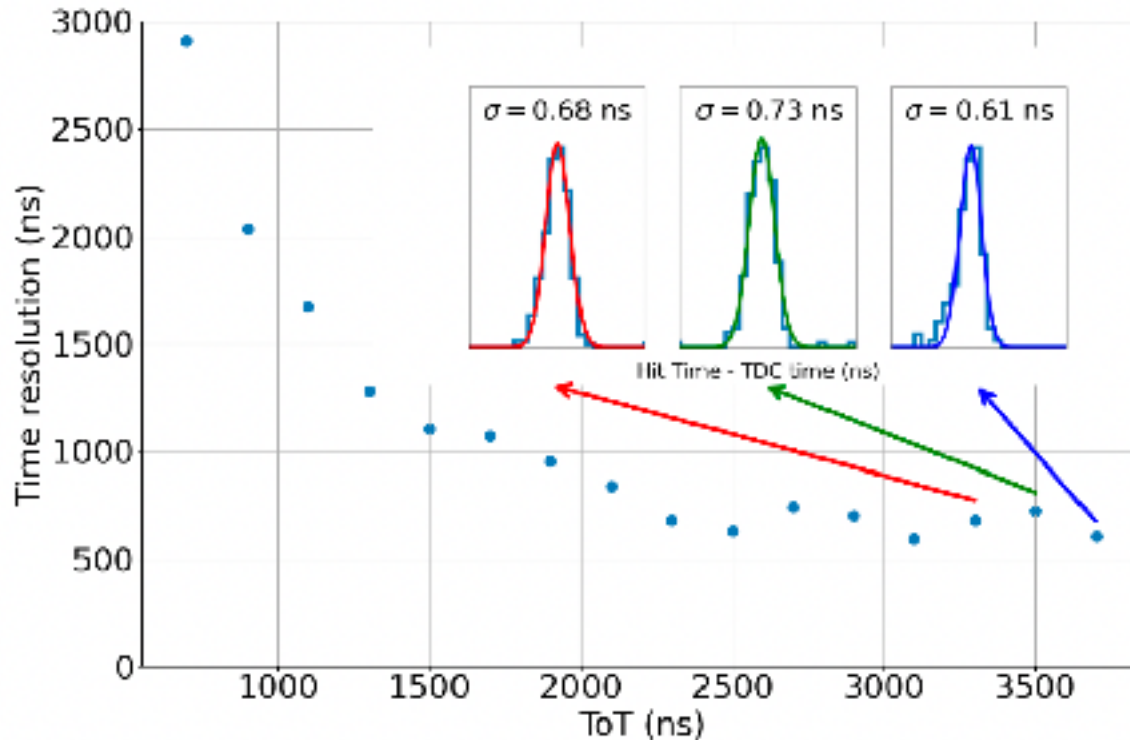
by Medipix4 collaboration

X. Llopart

		Timepix3	Timepix4
Technology		IBM 130nm	TSMC 65nm
Pixel Size		55 x 55 μm	$\leq 55 \times 55 \mu\text{m}$
Pixel arrangement		3-side buttable 256 x 256	4-side buttable 256 x 256 or bigger
Operating Modes	Data driven	PC (10-bit) and TOT (14-bit)	CRW: PC and iTOT (12...16-bit)
	Frame based	TOT and TOA	
Zero-Suppressed Readout	Data driven	< 80 MHits/s	< 500 MHits/s
	Frame based	YES	YES
TOT energy resolution		< 2KeV	< 1KeV
Time resolution		1.56ns	~200ps

X. Llopart, J. Alozy, R. Ballabriga, M. Campbell, R. Casanova, V. Gromov, E.H.M. Heijne, T. Poikela, E. Santin, V. Sriskaran, L. Tlustos, and A. Vitkovskiy. Timepix4, a large area pixel detector readout chip which can be tiled on 4 sides providing sub-200 ps timestamp binning. *Journal of Instrumentation*, 17(01):C01044, January 2022.

First optical Tpx4 measurements



Single photon sensitive (intensified) optical Tpx4 camera

600 ps possible with high gain intensifier

Direct detection after MCP in Timepix



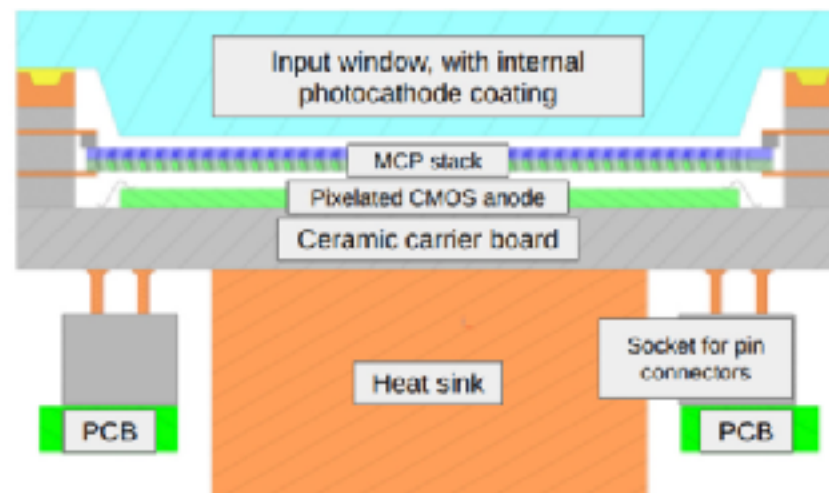
Nuclear Instruments and Methods in Physics
Research Section A: Accelerators,
Spectrometers, Detectors and Associated
Equipment

Volume 1082, Part 1, February 2026, 170965



Development and characterization of hybrid MCP-PMT with embedded Timepix4 ASIC used as pixelated anode

R. Bolzonella ^{a, b, *}, J. Alazay ^c, R. Ballabriga ^d, N.V. Biesuz ^a, M. Campbell ^e,
V. Cavallini ^{a, b}, A. Cotta Ramusino ^a, M. Fiorini ^{a, b}, E. Franzoso ^a, M. Guarise ^{a, b},
X. Llopart-Cudie ^c, G. Romolini ^{a, b}, A. Scaputi ^a



Has been implemented before with Timepix

Time resolution 60 ps (TBC)

Limitation: photocathode QE ~ 35%

Difficulties: HV and cooling in vacuum

Direct detection after MCP via resistive anode

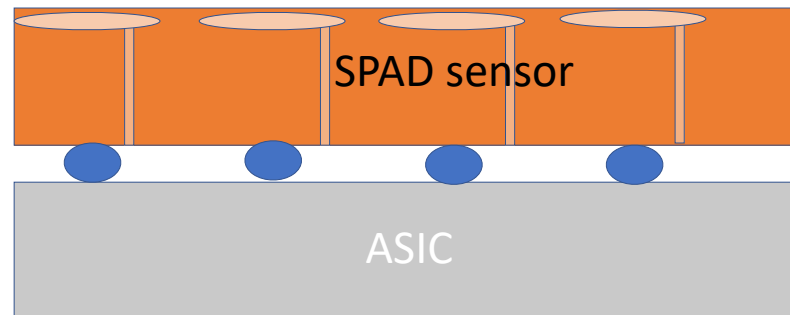


**PhotonPix™ Single
Photon Detector**

Verena Leopold's talk

20 ps timing in 2D

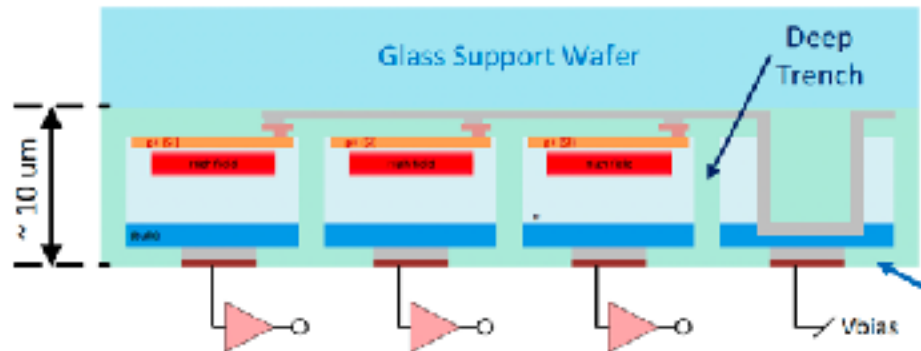
- 20 ps timing is needed for next round of CERN experiments in 10 years, there will be lots of investment in fast ASICs
- **Hybrid SPAD detector**
- Timepix4 chip: 200 ps bin (RMS 60 ps)
 - FE not designed for large SPAD signals, good time resolution but large dead time
 - Tests with SPADs planned in 2026
- **PicoPix chip: 40 ps bin**
 - Full size, will be used in one of CERN big experiments (LHCb)
 - Will have a SPAD compatible version
- Hybrid detector: SPAD + PicoPix



Difficulties: connectivity and cooling

20 ps timing in 2D

- 20 ps timing is needed for next round of CERN experiments in 10 years, there will be lots of investment in fast ASICs
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Possible solution

SPAD in CMOS: monolithic solution with 10 ps timing

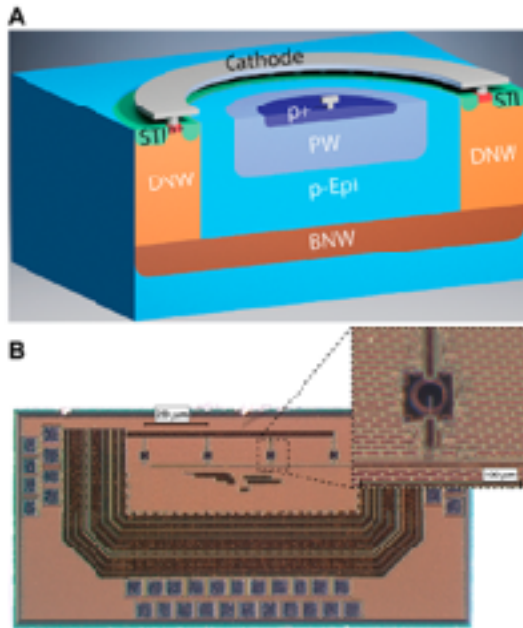
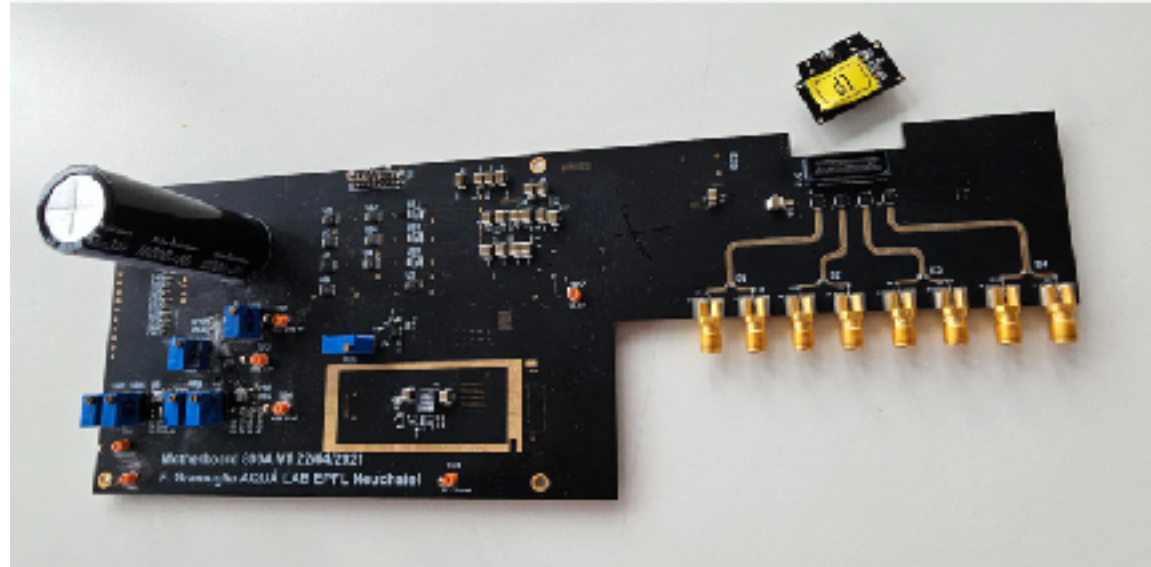


FIGURE 1 | (A): SPAD cross section. (B): Micrograph of the implemented chip embedding 25 μm diameter SPADs with integrated pixel circuit [21].



F. Gramaglia, M.-L. Wu, C. Bruschini, M.-J. Lee, and E. Charbon. A low-noise CMOS SPAD pixel with 12.1 ps SPTR and 3 ns dead time, *IEEE Journal of Selected Topics in Quantum Electronics* **28**, 1 (2022).

- **SuperSPAD sensor**
 - Developed in AQUA group in EPFL
 - Excellent results from QUASAR collaboration
 - 7.5 ps FWHM time resolution
- **QUASAR1 sensor** (Ivan Cardea's talk at this workshop)
 - Monolithic CMOS logic and TDCs (!)
 - Possible extension to full 2D version?

Main points to take home

- Two-photon interferometry allows independent stations over long baselines
- New ideas suggest quantum sensing technology can dramatically enhance astrometric precision, requires spectrometers and single photon cameras with 10 ps resolution
- Promising results with 40 ps & 40 pm spectrometers

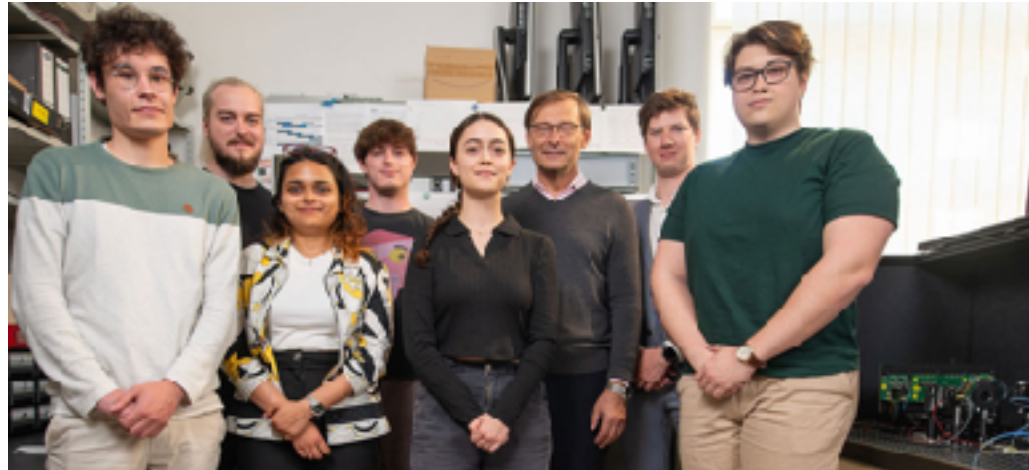
Broad program in quantum-assisted optical interferometry ahead, efforts underway to develop new timing technologies (also needed for other applications)

Synergy in techniques and instrumentation with SII community

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Thank you for your attention!



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