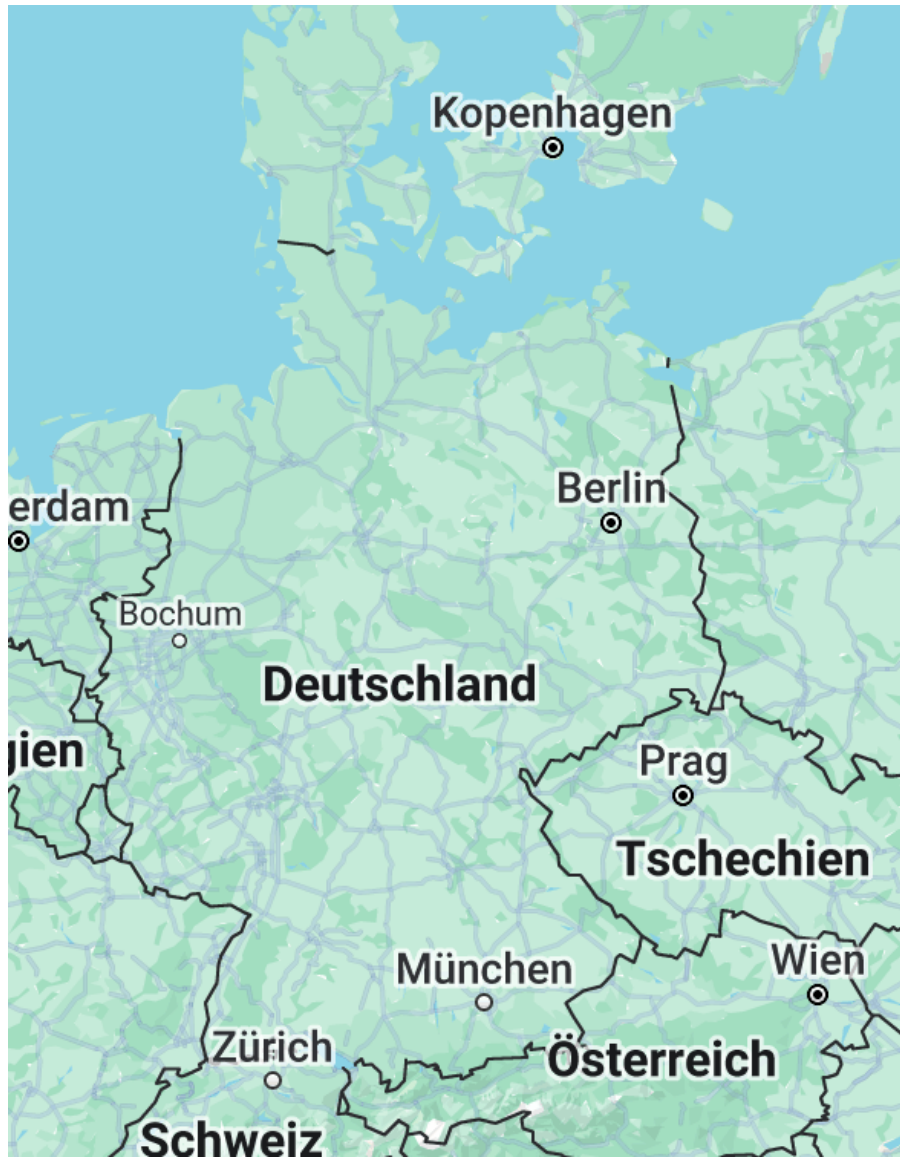
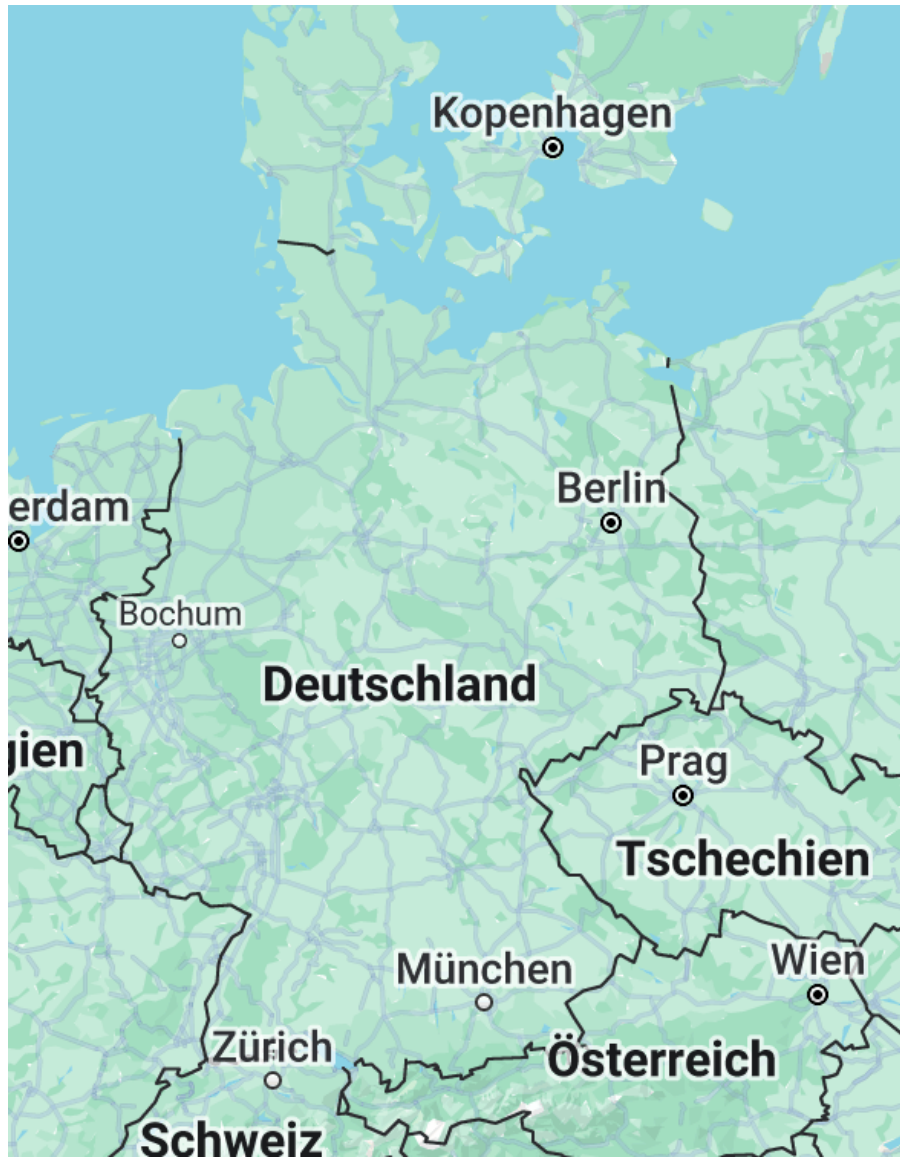


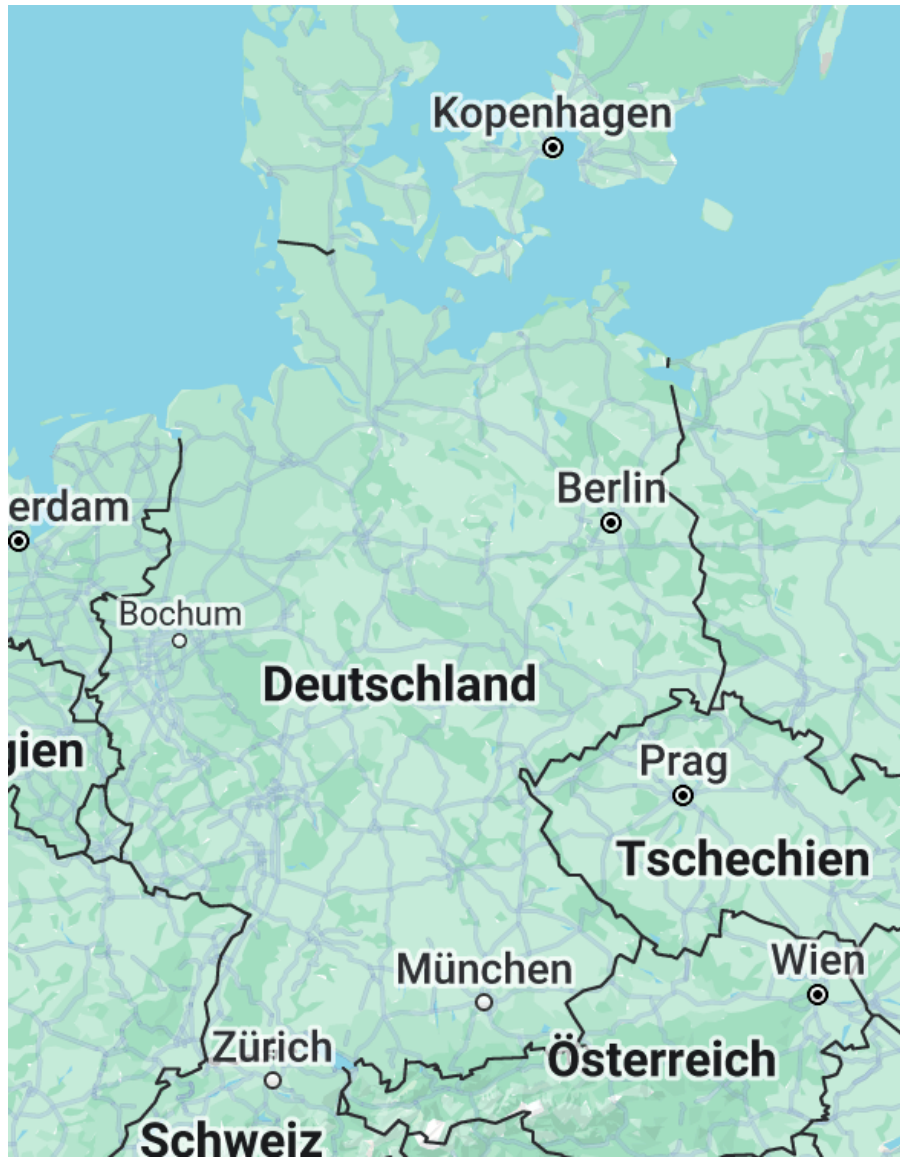
Neutrinos

Anna Franckowiak
Ruhr-University Bochum









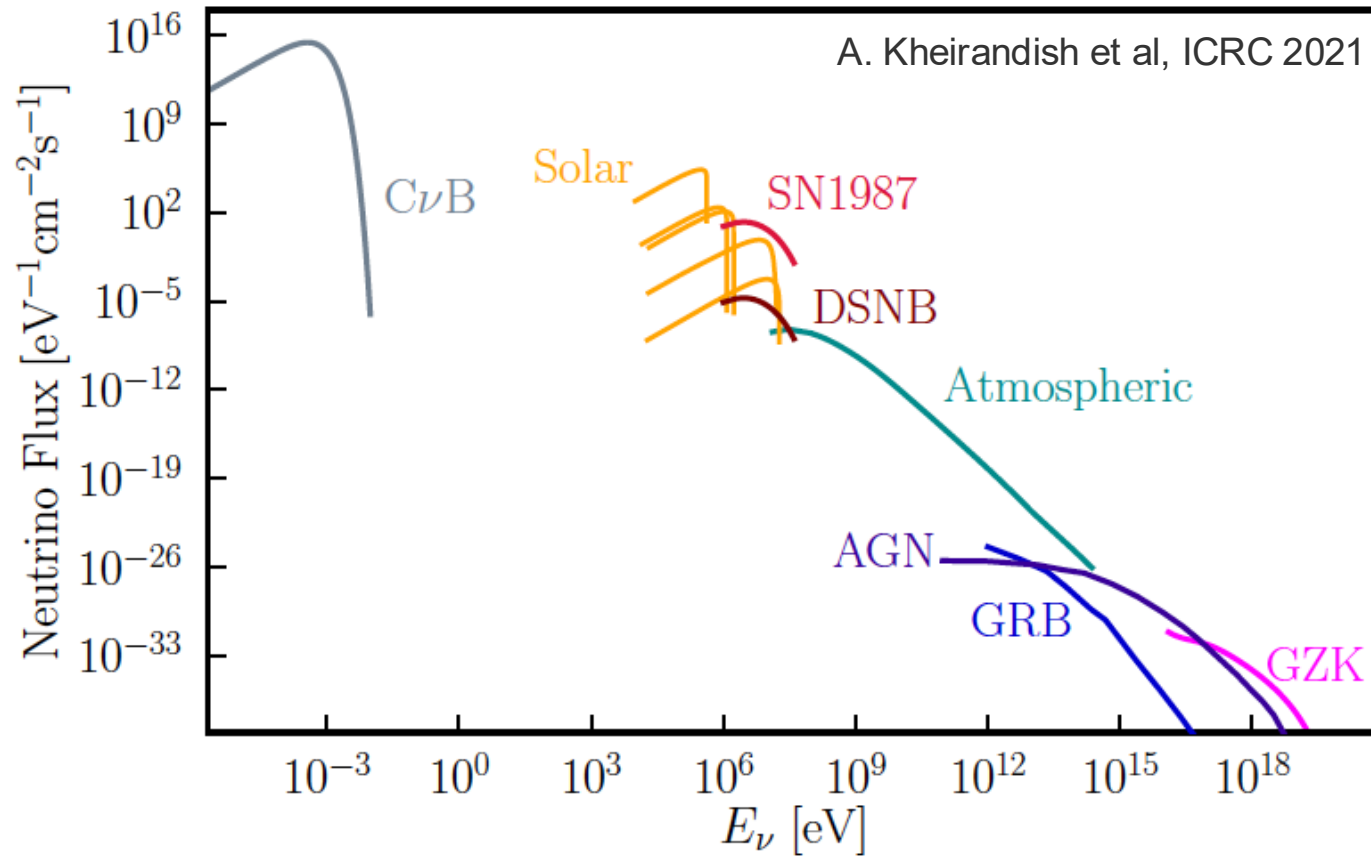
**Which neutrino sources do
you know?**

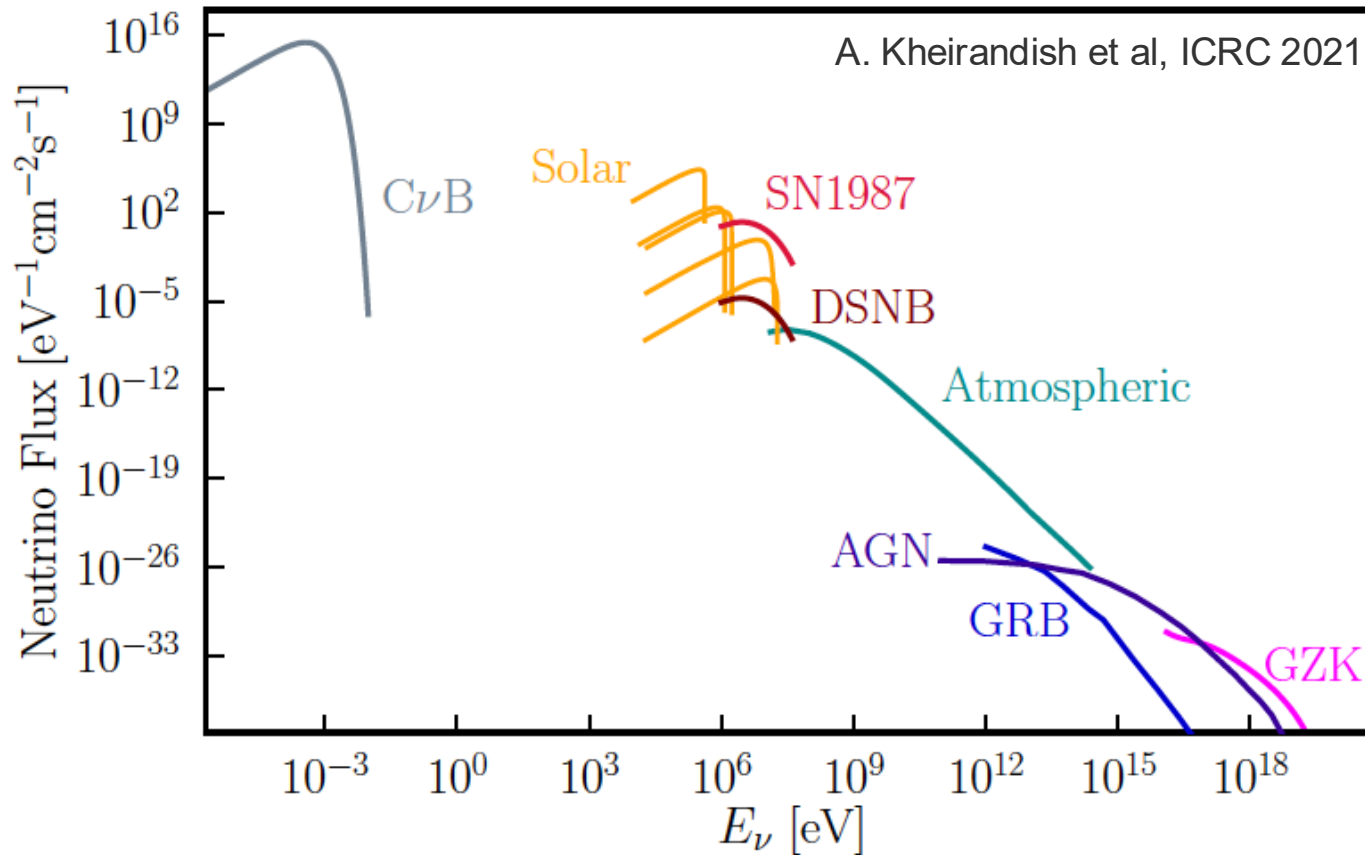
Content of the 3 lectures

- Low-energy (MeV) neutrinos
 - Sun
 - Supernovae
- High-energy (TeV-PeV) neutrinos
 - Astrophysics
 - Particle physics
- Even higher energies?

Learning Objectives

- What have we learned from MeV solar and supernova neutrinos?
- Are we ready for the next Galactic supernova?!
- How can we detect high-energy astrophysical neutrinos?
- What is the background in the search for high-energy cosmic neutrinos and how can we disentangle it from the signal?
- What have we learned so far from high-energy neutrino detections?
- How can we improve in the future?





↔

Ideas: neutrino capture
on unstable nucleus
(e.g. tritium)

↔

↔

Neutrino capture on
stable nucleus
(Chlorine, Gallium)

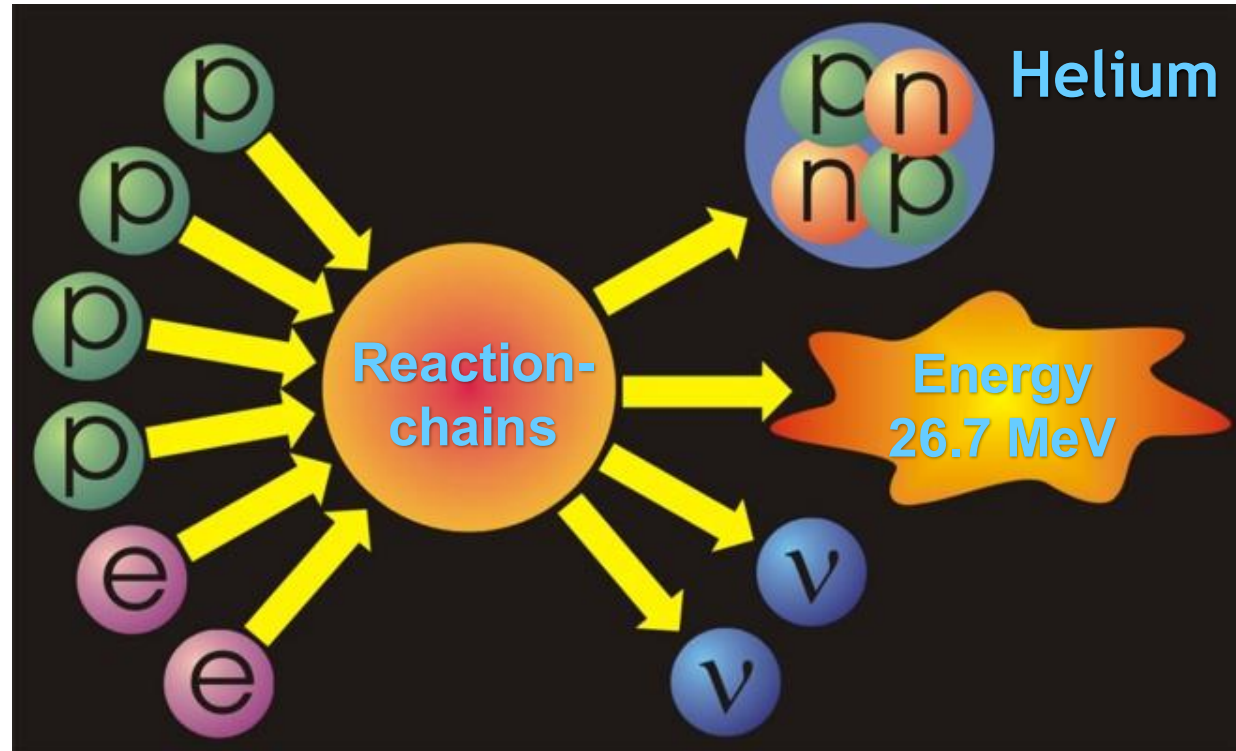
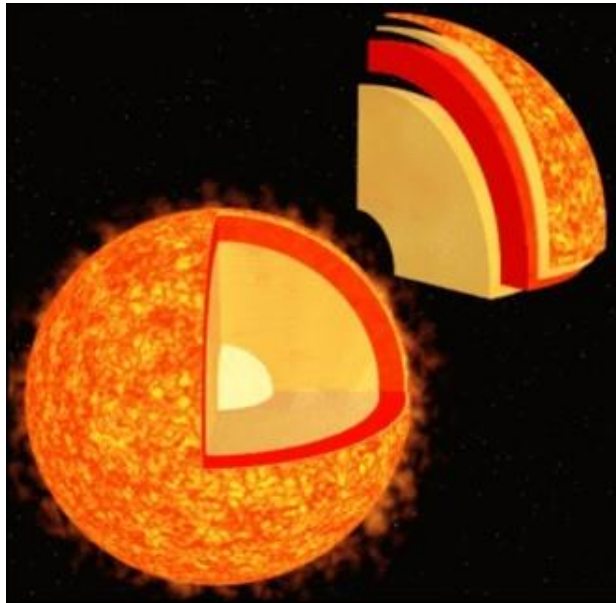
↔

Water / ice
Cherenkov
detectors

↔

Radio arrays,
cosmic-ray
detectors

Solar Neutrinos



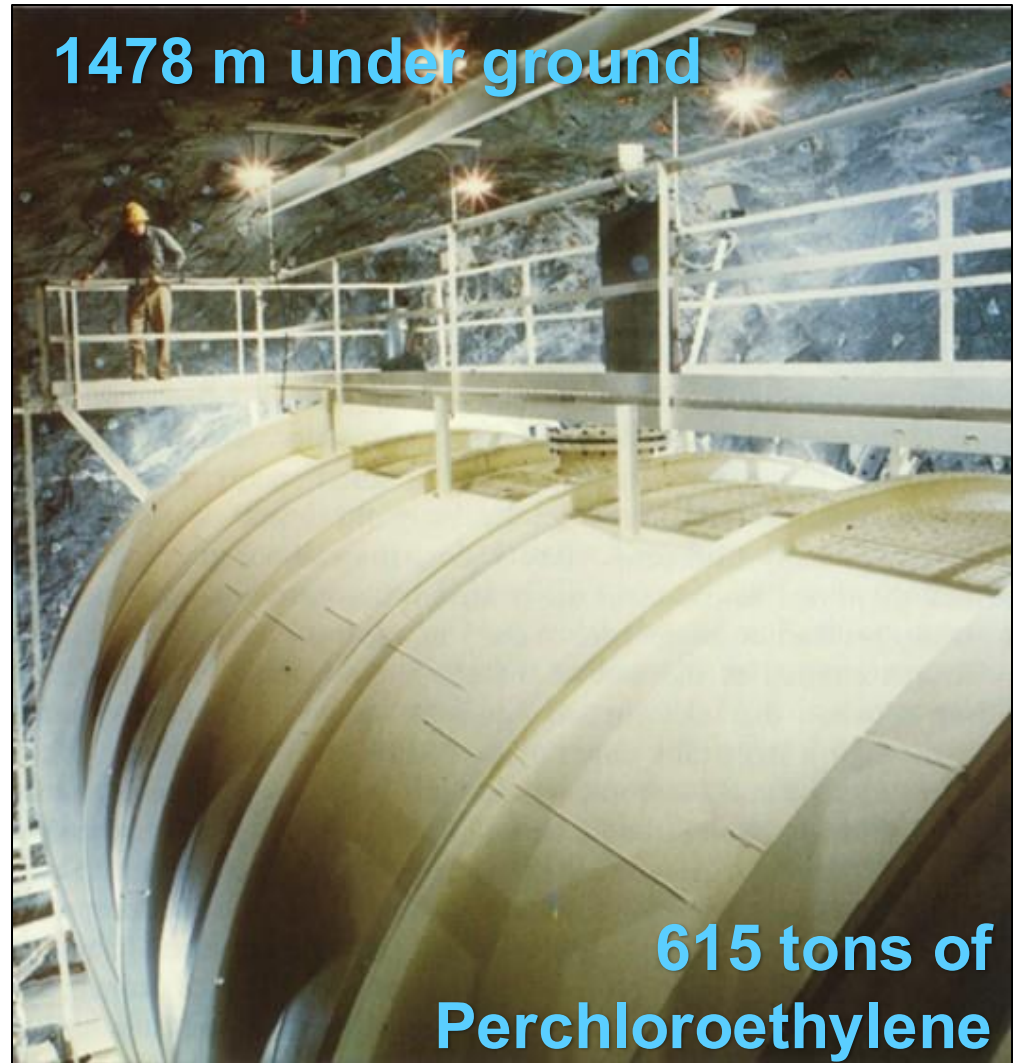
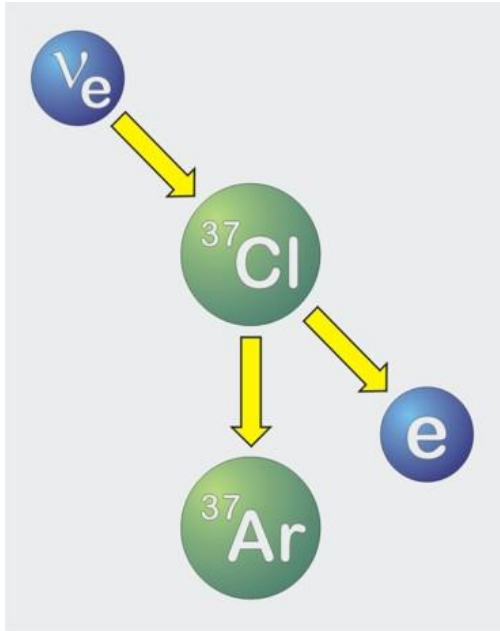
1967

Solar radiation:

- 98% light
- 2% **electron-neutrinos**

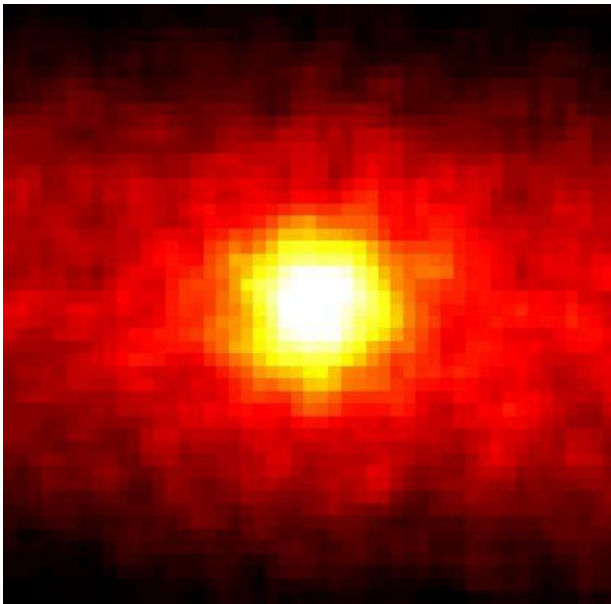
→ at Earth: 66 billion neutrinos / cm^2 / s

The Homestake Experiment



Solar Neutrinos

Astronomy Picture of the Day
June 5, 1998



The Sun in Neutrinos seen
by Super-Kamiokande

Combining neutrinos and electromagnetic information led to:

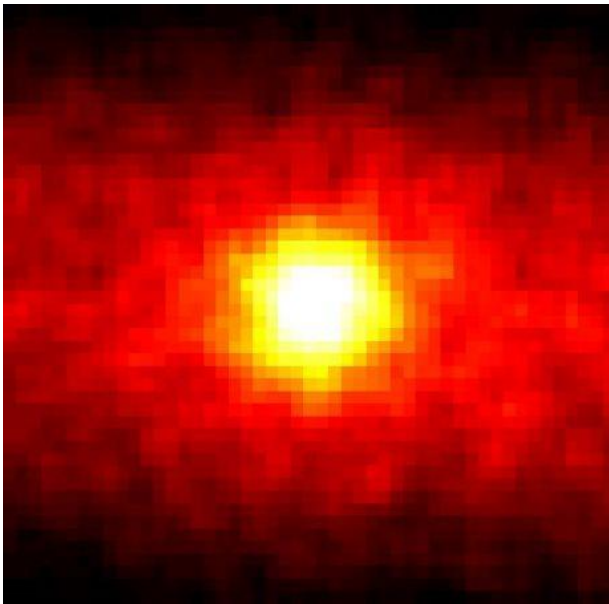
- **The solar neutrino problem**



Solar Neutrinos

Astronomy Picture of the Day

June 5, 1998



The Sun in Neutrinos seen
by Super-Kamiokande

Combining neutrinos and electromagnetic information led to:

- Confirmation of model of fusion
- New understanding of the standard model of particle physics
→ neutrinos oscillate → have mass!

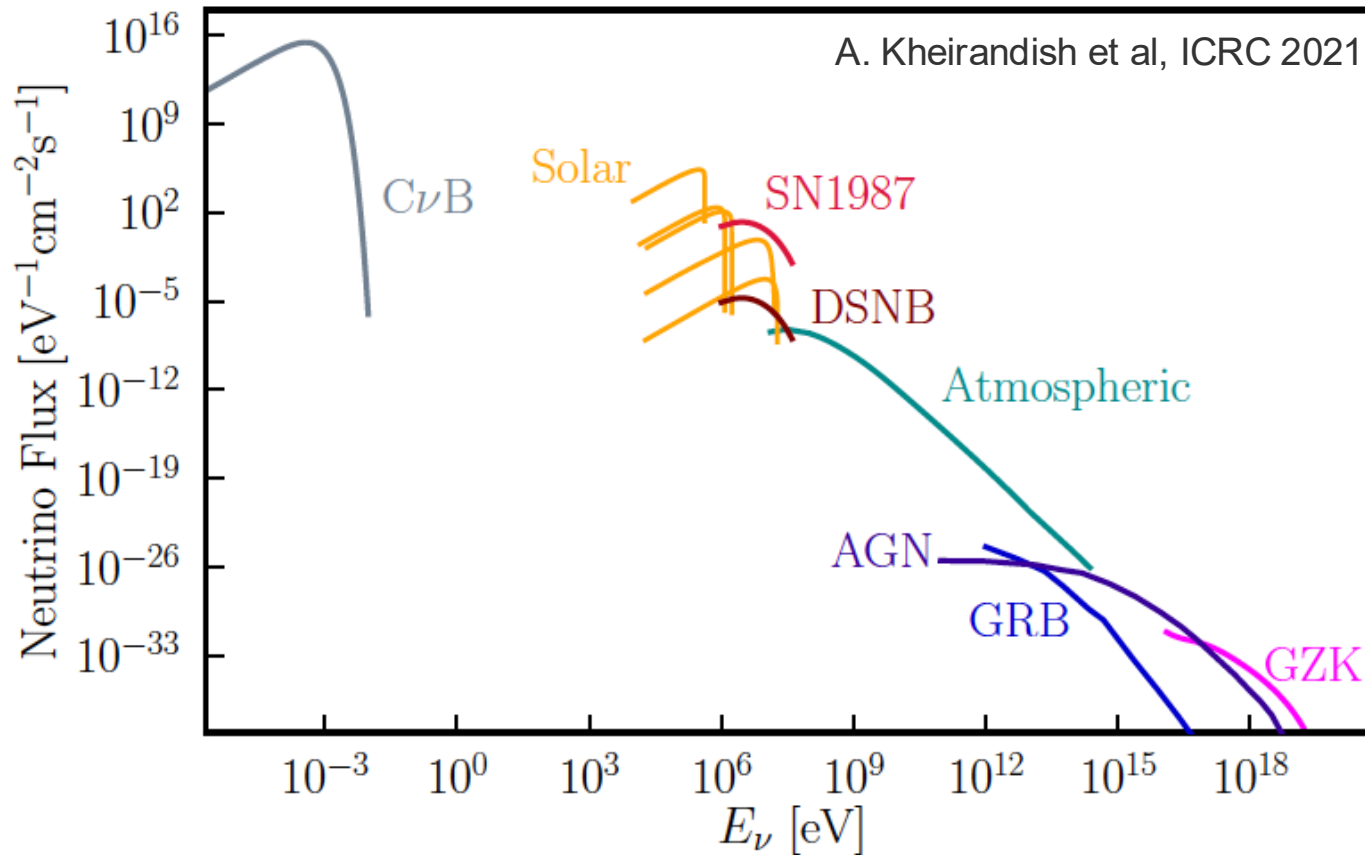




Take home message

Too few electron neutrinos have been detected from the sun, but the number matches expectations if all flavors are considered.

This is explained by neutrino oscillation, which only works when neutrinos have mass.



↔

Ideas: neutrino capture
on unstable nucleus
(e.g. tritium)

↔

↔

Neutrino capture on
stable nucleus
(Chlorine, Gallium)

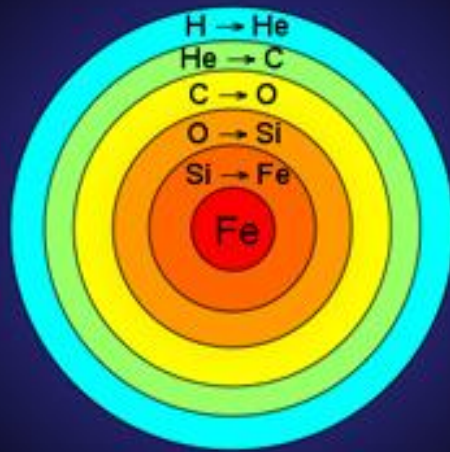
↔

Water / ice
Cherenkov
detectors

↔

Radio arrays,
cosmic-ray
detectors

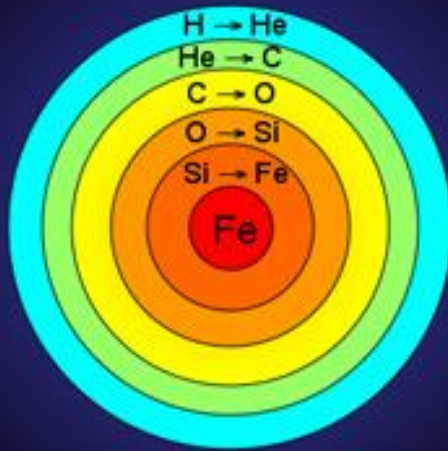
Supernova Neutrinos



For a 25 solar mass star:

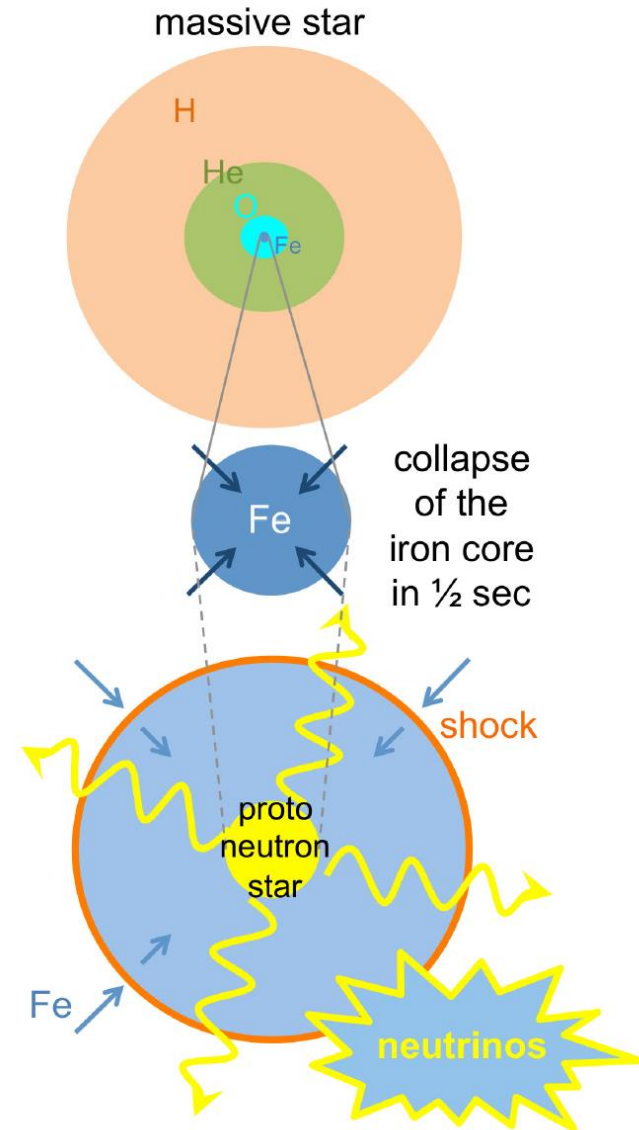
Stage	Duration
$H \rightarrow He$	7×10^6 years
$He \rightarrow C$	7×10^5 years
$C \rightarrow O$	600 years
$O \rightarrow Si$	6 months
$Si \rightarrow Fe$	1 day
Core Collapse	1/4 second

Supernova Neutrinos

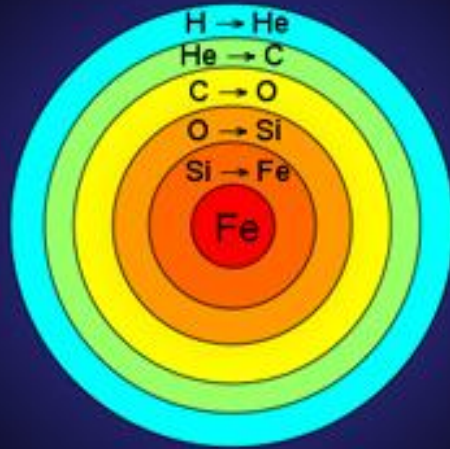


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Supernova Neutrinos

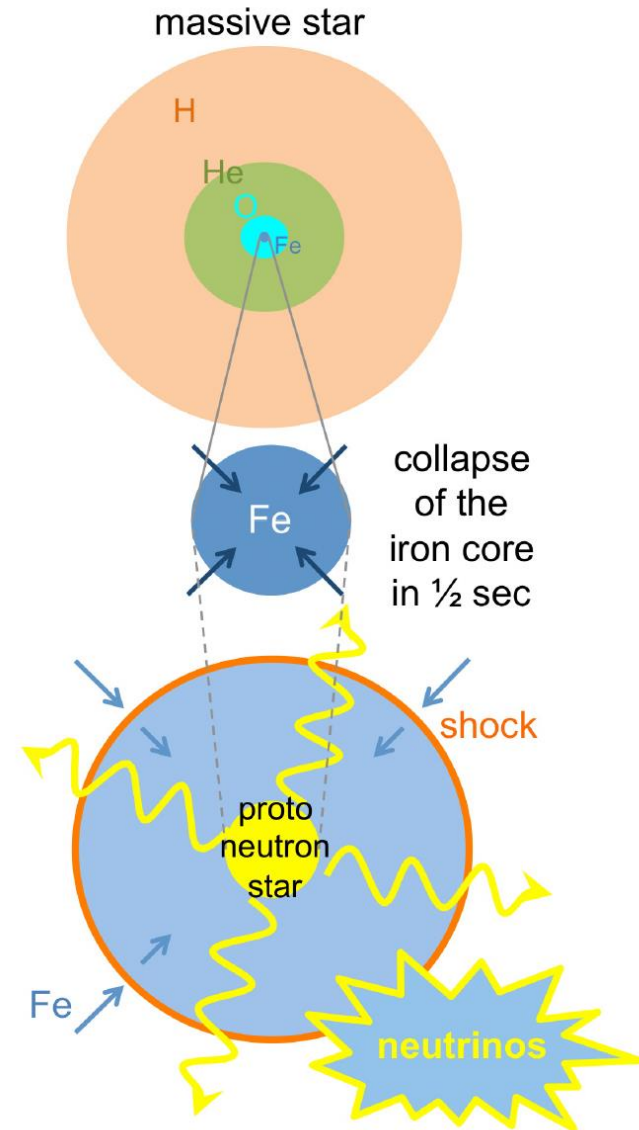


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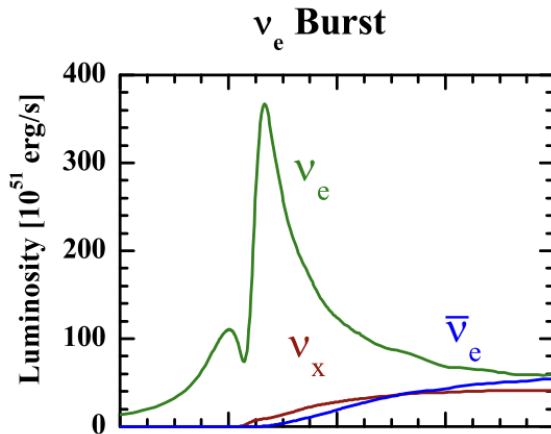
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Si → Fe	1 day
Core Collapse	1/4 second

Neutronization burst: $e^- + p \rightarrow n + \nu_e$
 Thermal neutrinos of all flavors

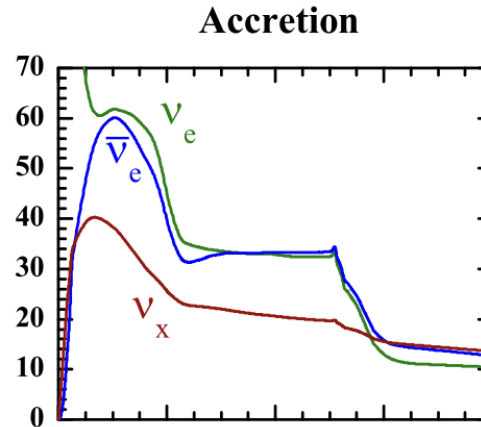
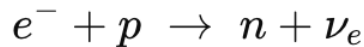
- $\sim 10^{58}$ neutrinos in ~ 10 s
- 99% of gravitational energy
- typical energy: 10-20 MeV



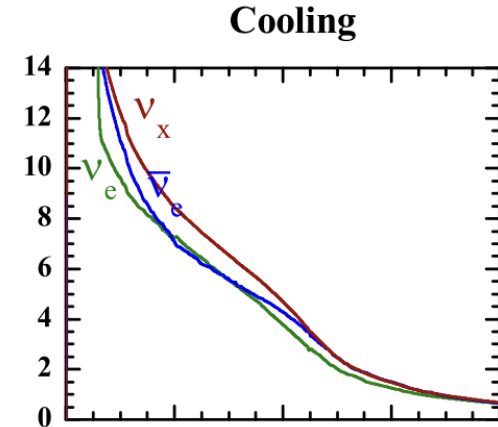
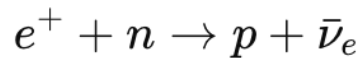
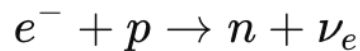
Expected neutrino signal



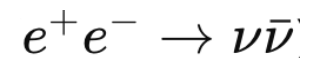
Initial neutronization burst



Material falls onto the proto-neutron star, e^- and e^+ in hot plasma



proto-neutron star loses trapped thermal energy mainly by neutrino emission of all flavors

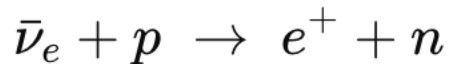


(Super) Kamiokande

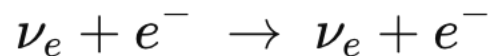


Detection channels:

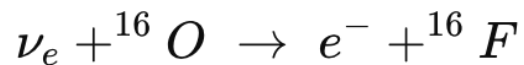
Inverse beta decay:



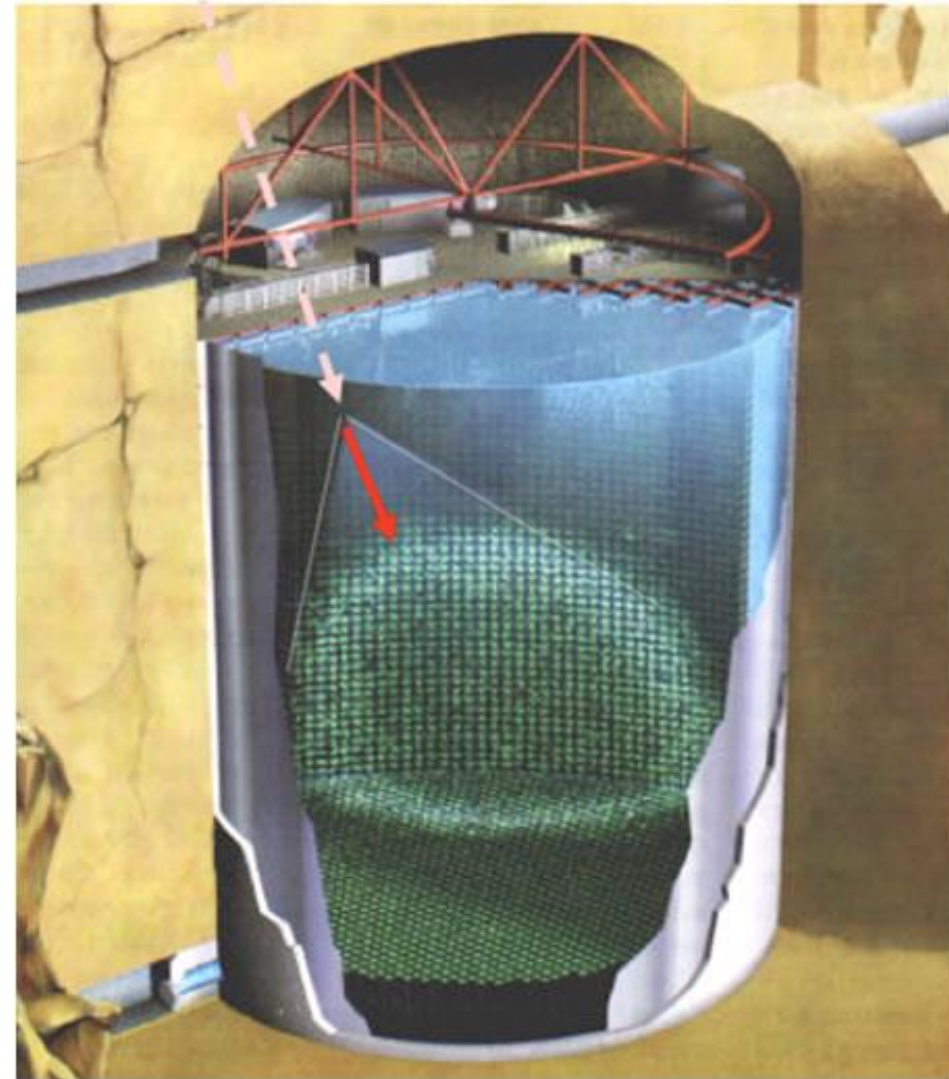
Elastic scattering:



Charged current absorption on oxygen:



Volume: 3kT (50kT)

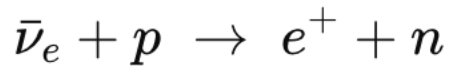


(Super) Kamiokande



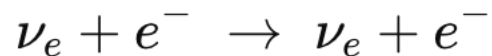
Detection channels:

Inverse beta decay:

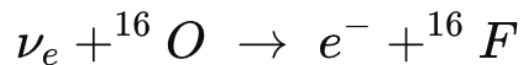


Most
efficient

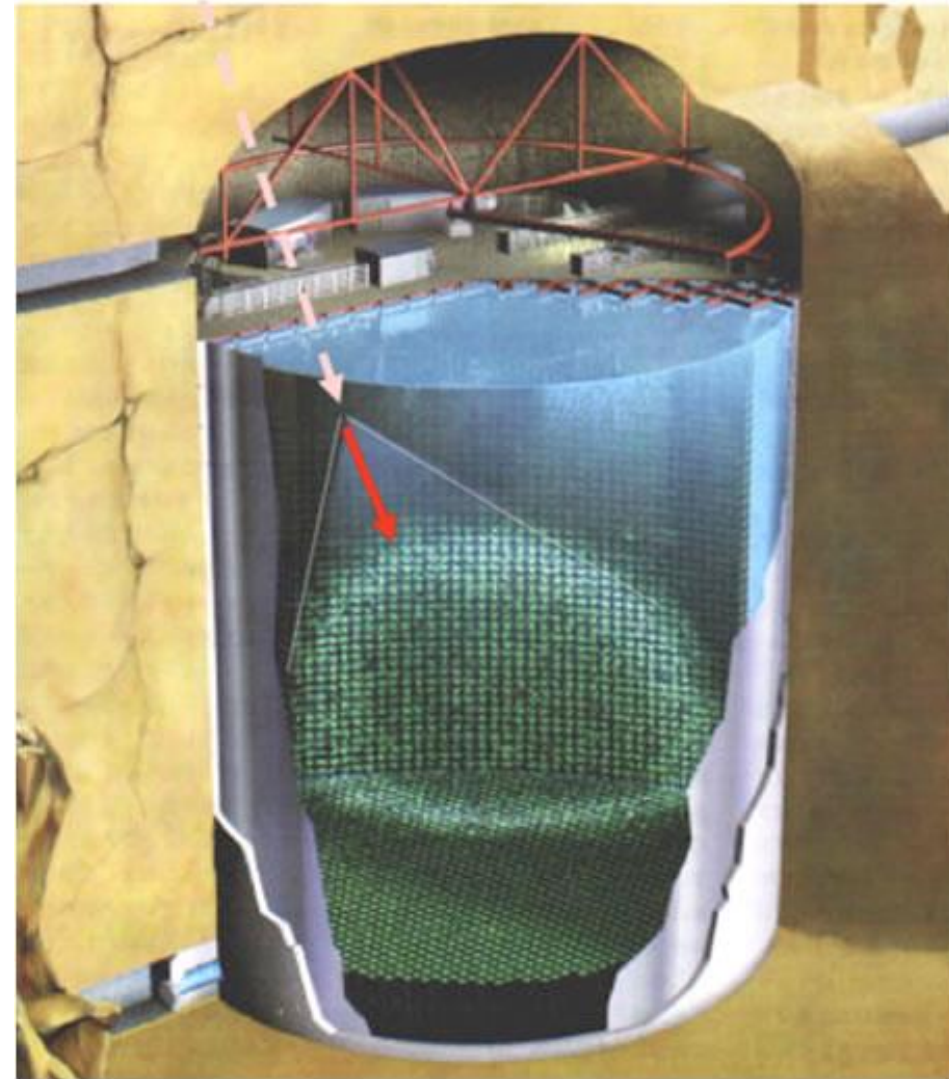
Elastic scattering:



Charged current absorption on
oxygen:

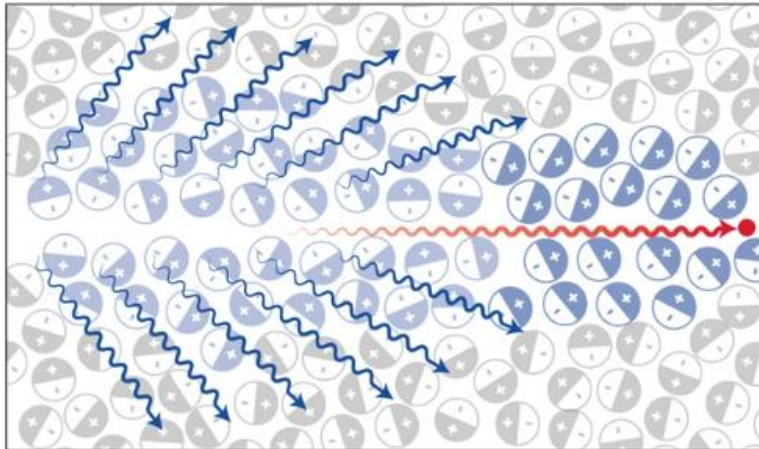
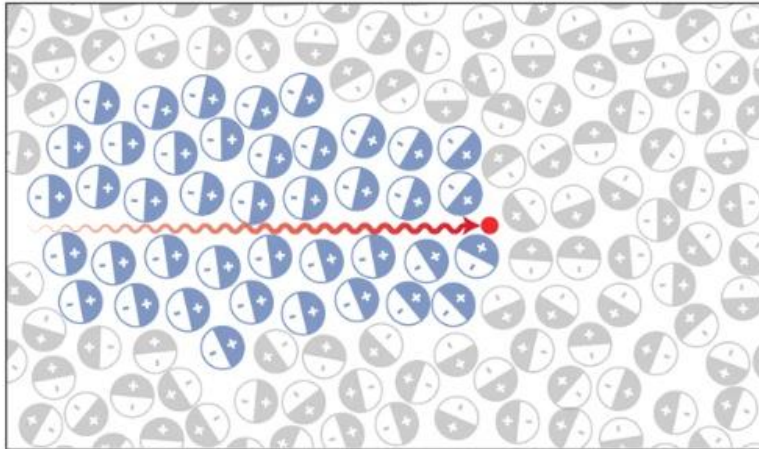


Volume: 3kT (50kT)

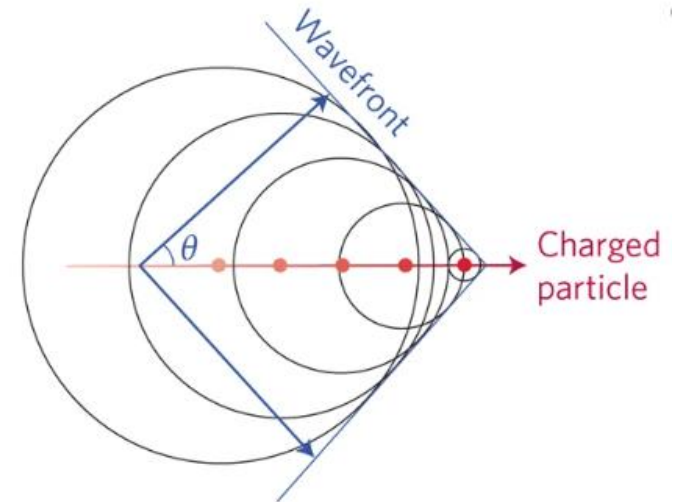
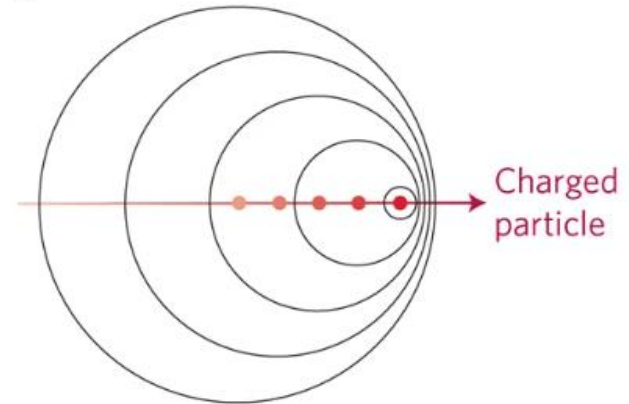


Different detection technique: Cherenkov Effect

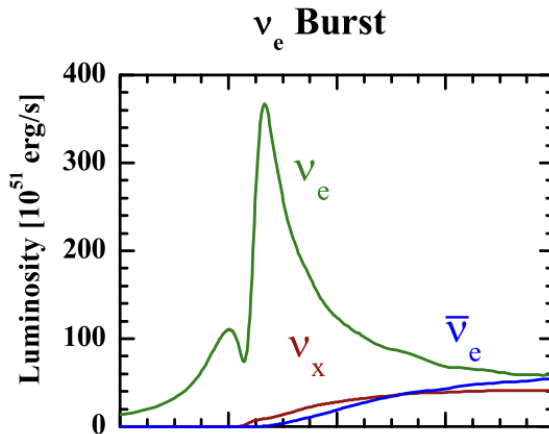
a



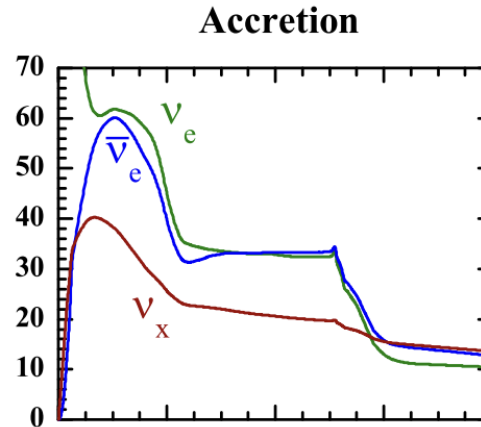
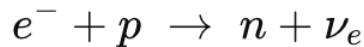
b



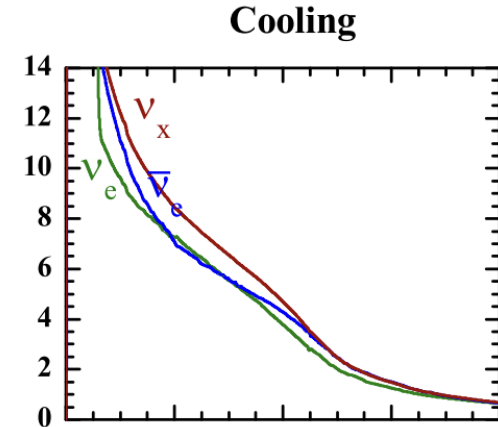
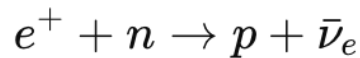
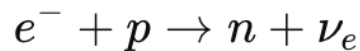
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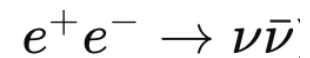
Initial neutronization burst



Material falls onto the proto-neutron star, e^- and e^+ in hot plasma



proto-neutron star loses trapped thermal energy mainly by neutrino emission of all flavors



Expected signal in Super-K dominated by **anti-electron neutrinos**

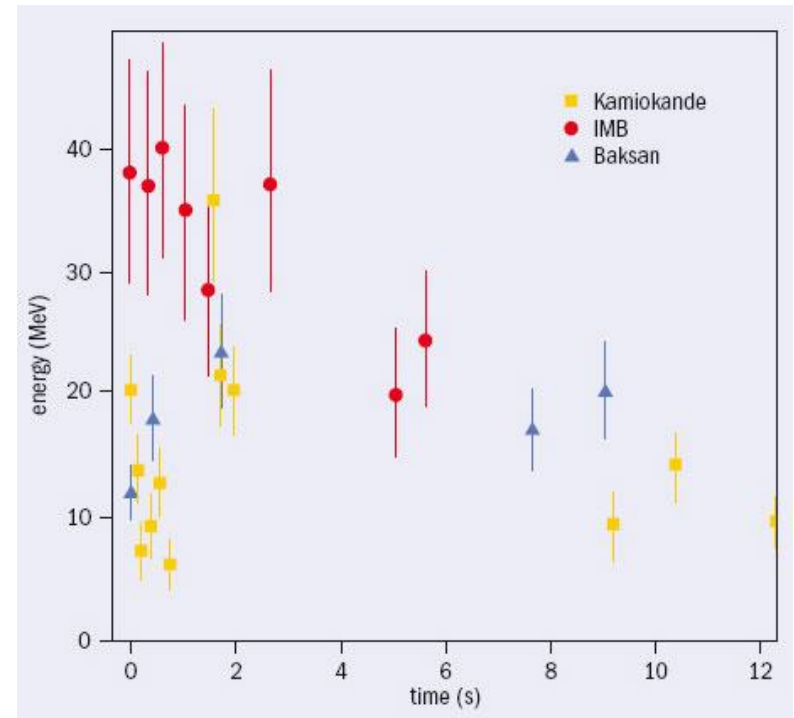
First and only Supernova neutrino detection

Optical detection of SN1987A in LMC



Neutrinos arrived 2-3h earlier

MeV neutrino burst



Location of Magellanic Clouds

Distance: 50kpc = 160000 Ly

NGC 2419

300 000 l.y.

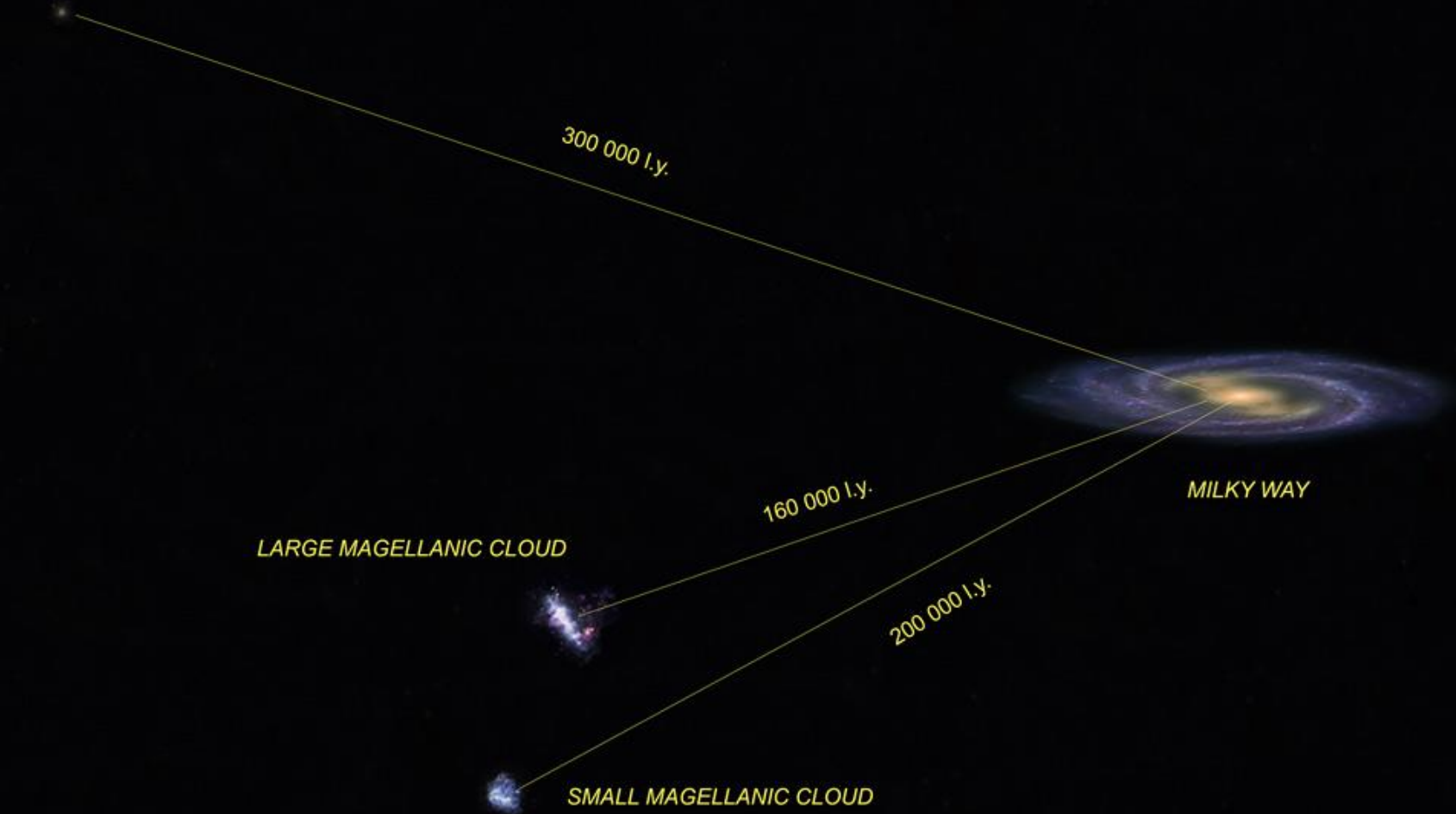
160 000 l.y.

200 000 l.y.

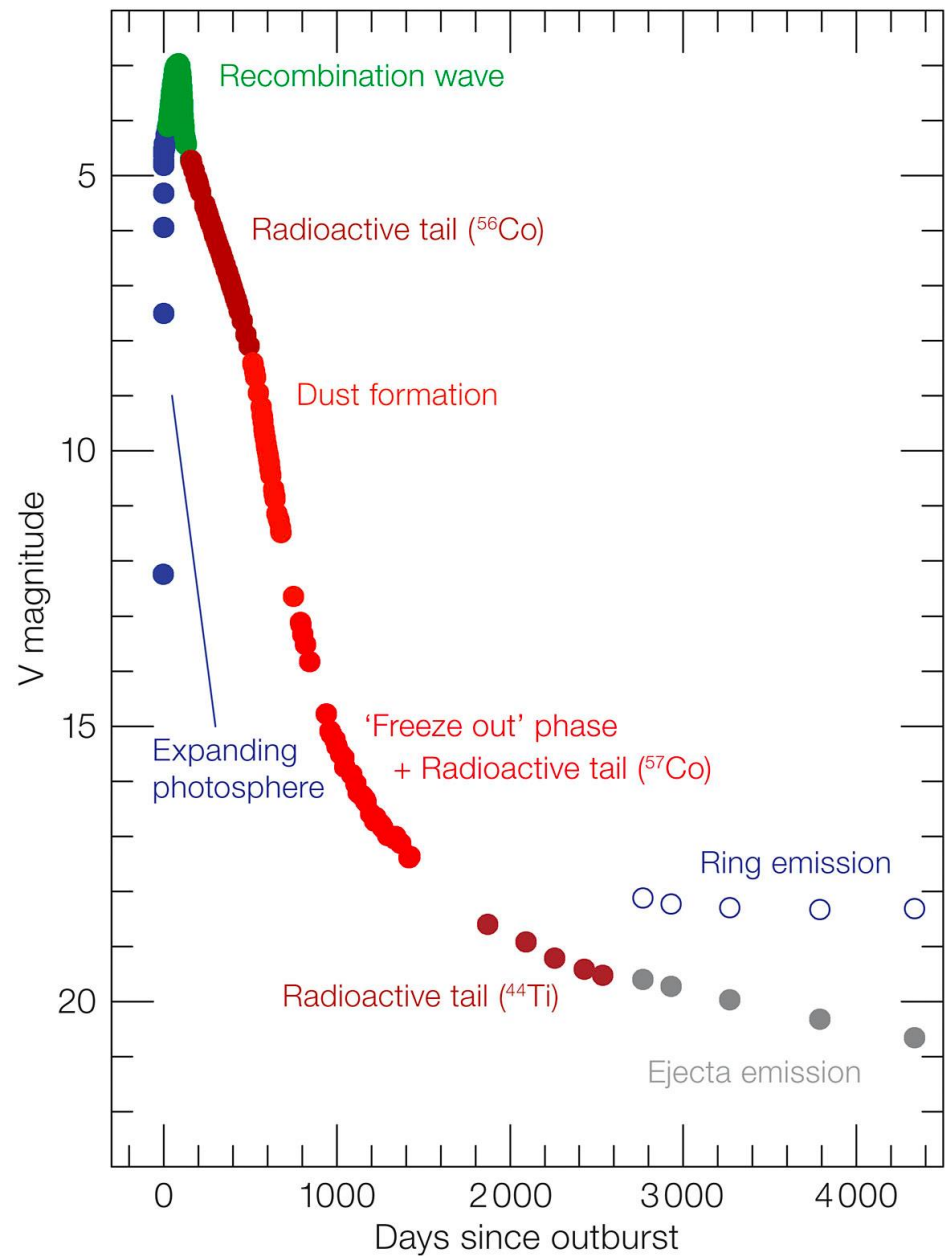
LARGE MAGELLANIC CLOUD

SMALL MAGELLANIC CLOUD

MILKY WAY

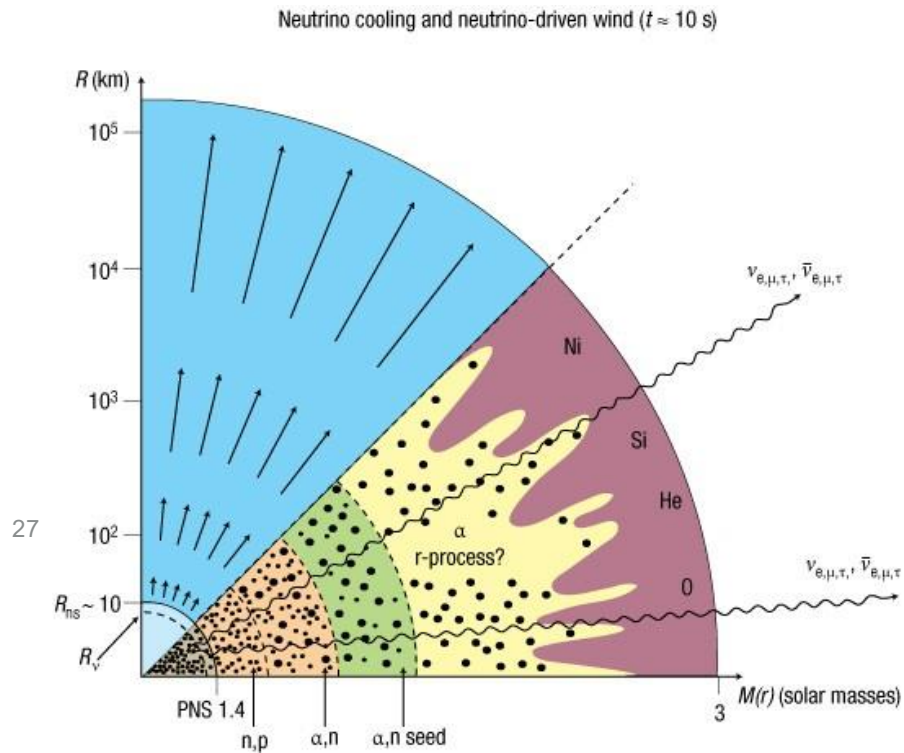


SN1987A light curve



First (and only) detection of a Supernova

First direct confirmation of our
basic picture of a stellar collapse

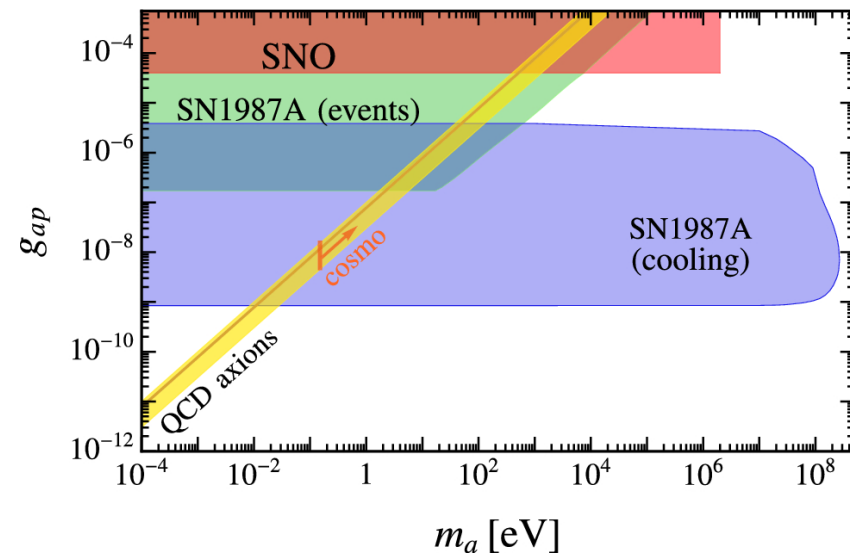
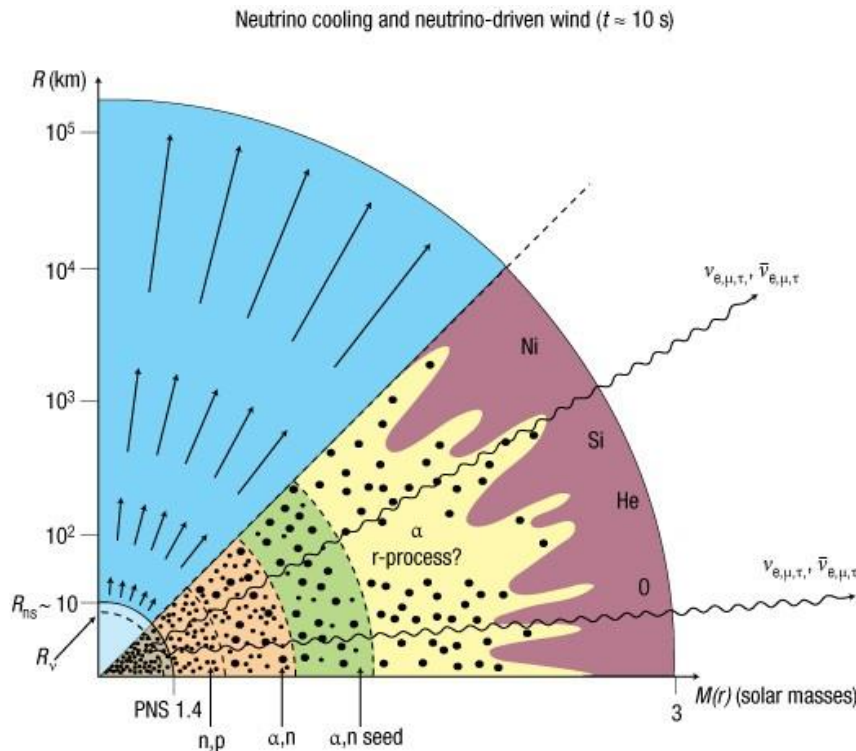


Woosely & Janka,
Nature Physics 2005

First (and only) detection of a Supernova

First direct confirmation of our basic picture of a stellar collapse

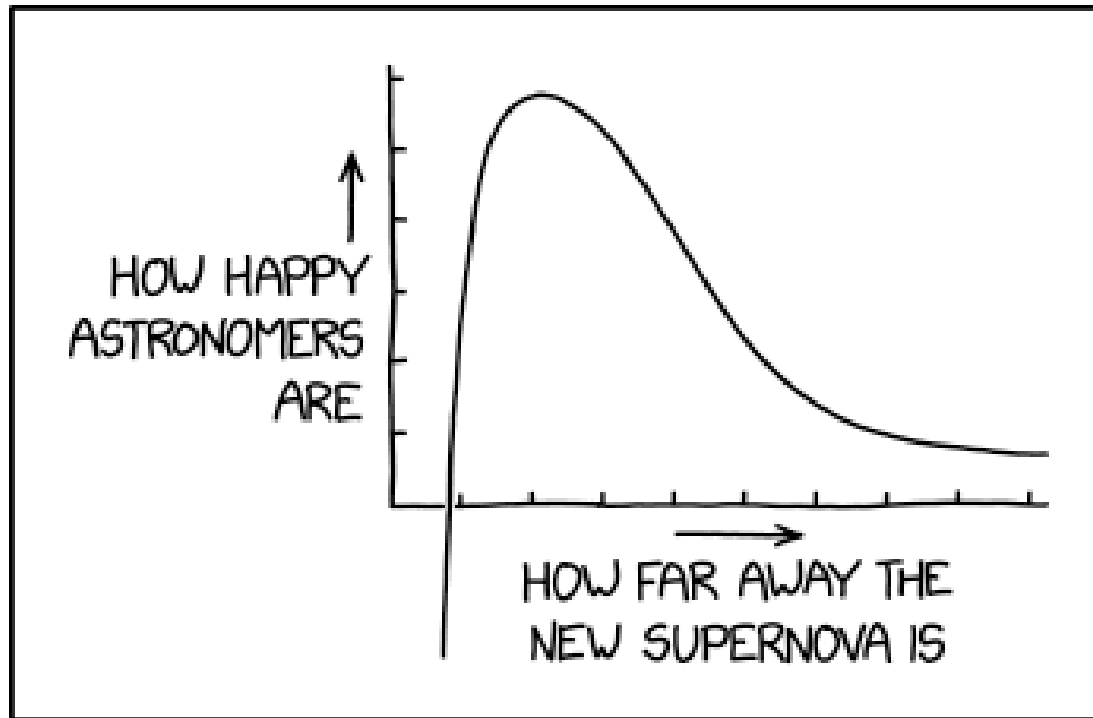
Constraints on exotic physics (e.g. axions)



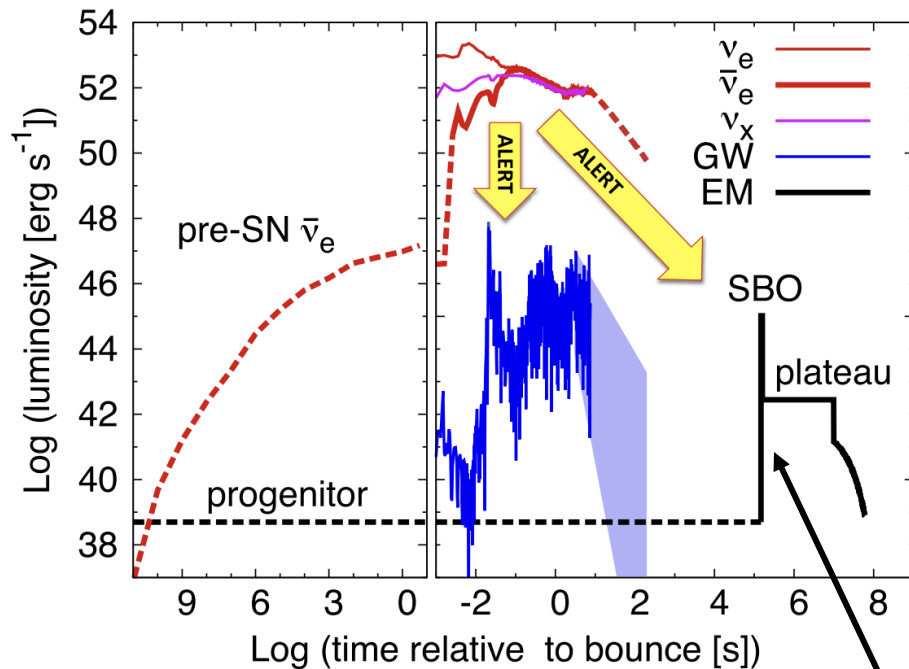
Lella et al. PRD 109 (2024)

Woosely & Janka,
Nature Physics 2005

Are we ready for the next Galactic Supernova?!

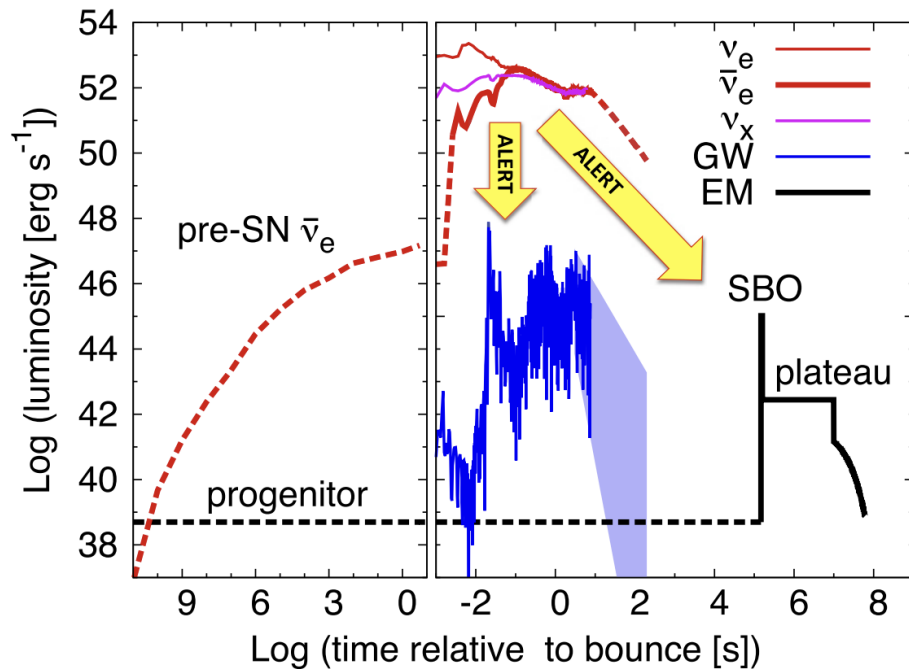


Supernova early warning system



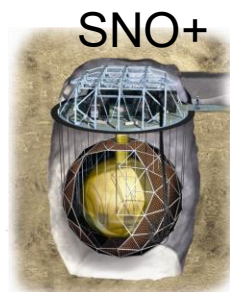
Shock breakout in X-ray
and UV emission

Supernova early warning system

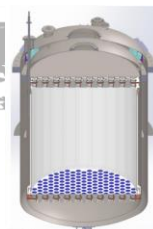


MeV neutrino burst as trigger for electromagnetic supernovae observations

Detectors participating in SNEWS



XENONnT



Baksan



Super-K



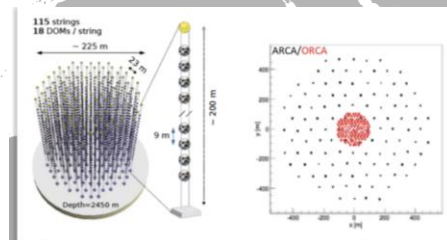
SBND



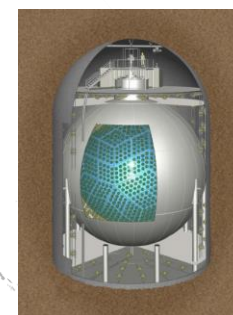
LVD



NOvA



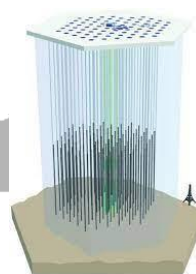
KM3NeT



KamLAND

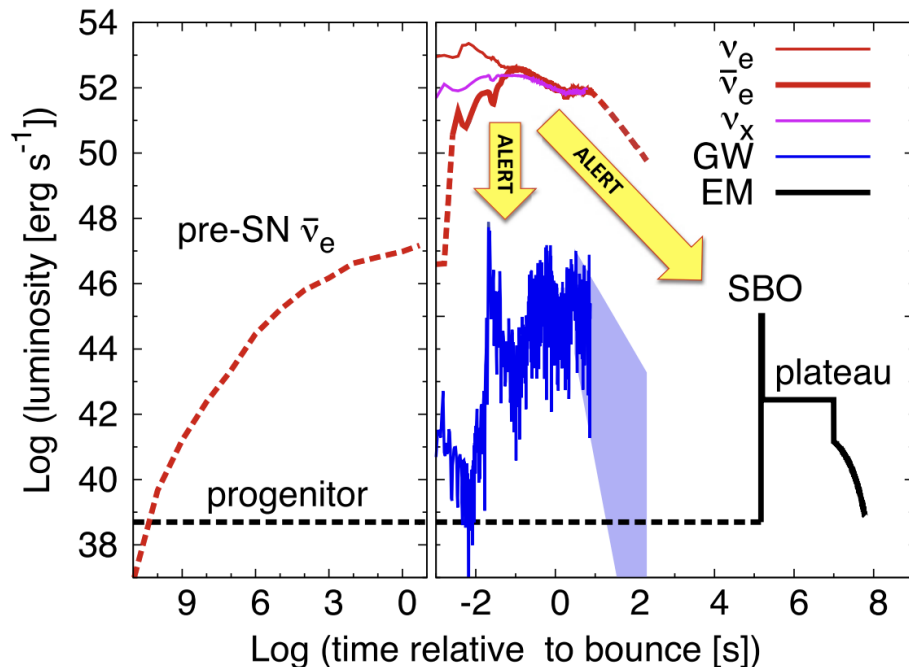


HALO



IceCube

Supernova early warning system

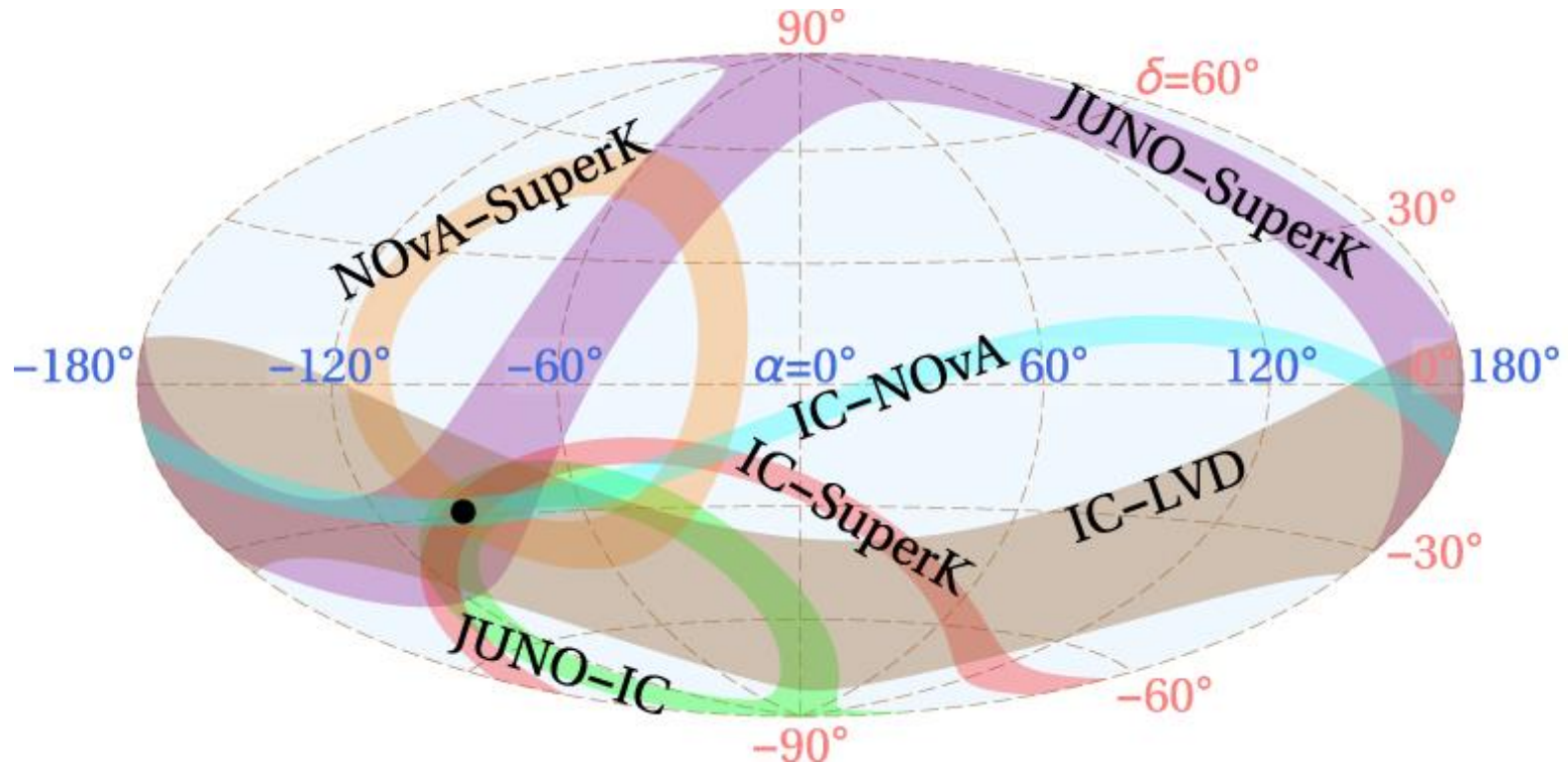


MeV neutrino burst as trigger for electromagnetic supernovae observations

SNEWS 2.0:

- new infrastructure
- public sub-threshold alerts
- pointing using inter-experiment triangulation
- searches for pre-supernova neutrinos

Supernova localization



Coordinated follow-up observations with wide-field-of-view instruments are necessary

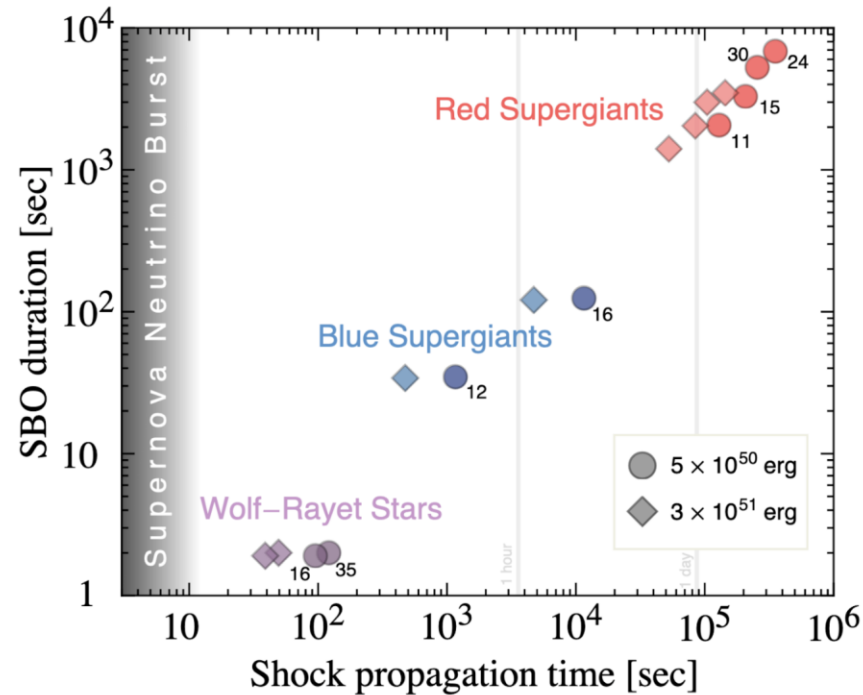
Catching the next Galactic Neutrino Supernova

- Unprecedented insights into the explosion mechanism
- Information about surrounding material
- Spatially resolved imaging of early phases of explosion

Catching the next Galactic Neutrino Supernova

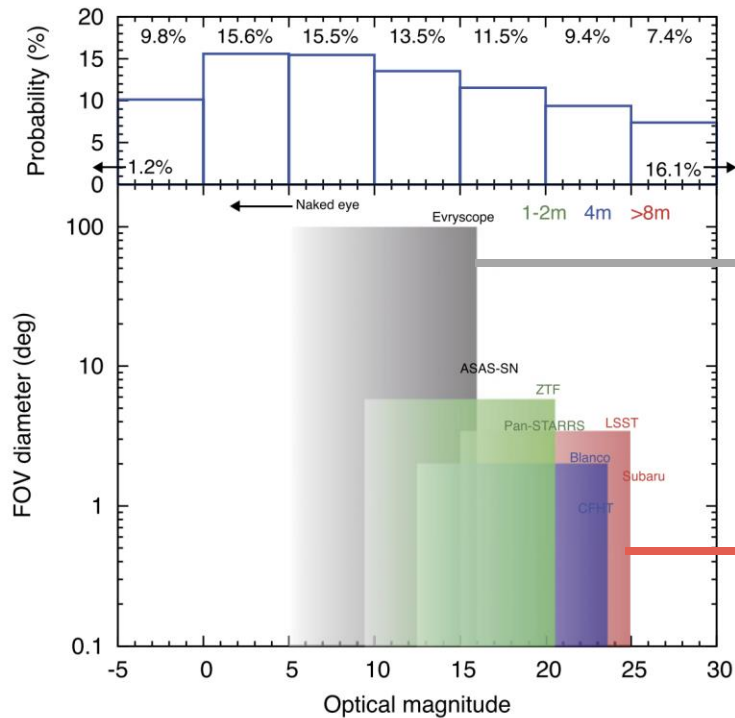
- Unprecedented insights into the explosion mechanism
- Information about surrounding material
- Spatially resolved imaging of early phases of explosion

Delay between neutrino burst and optical signal: 2 min to 2 days

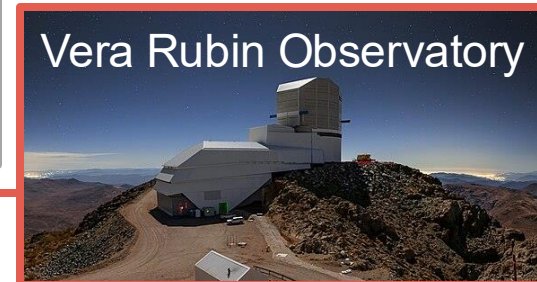
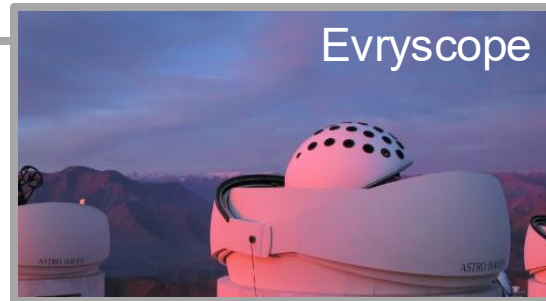


Catching the next Galactic Neutrino Supernova

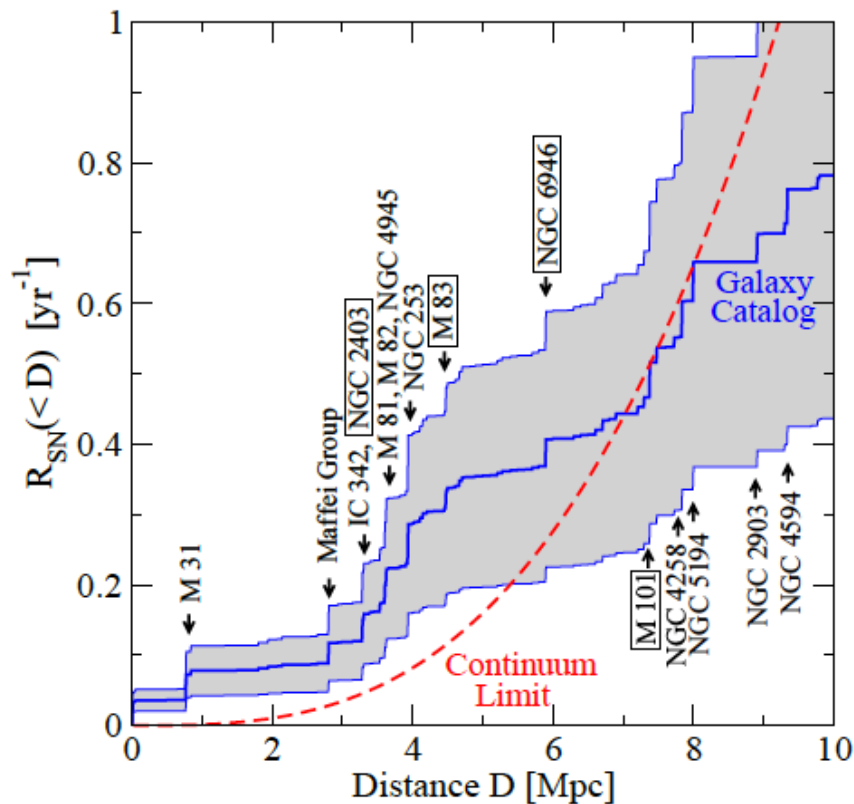
Detectability of shock breakout



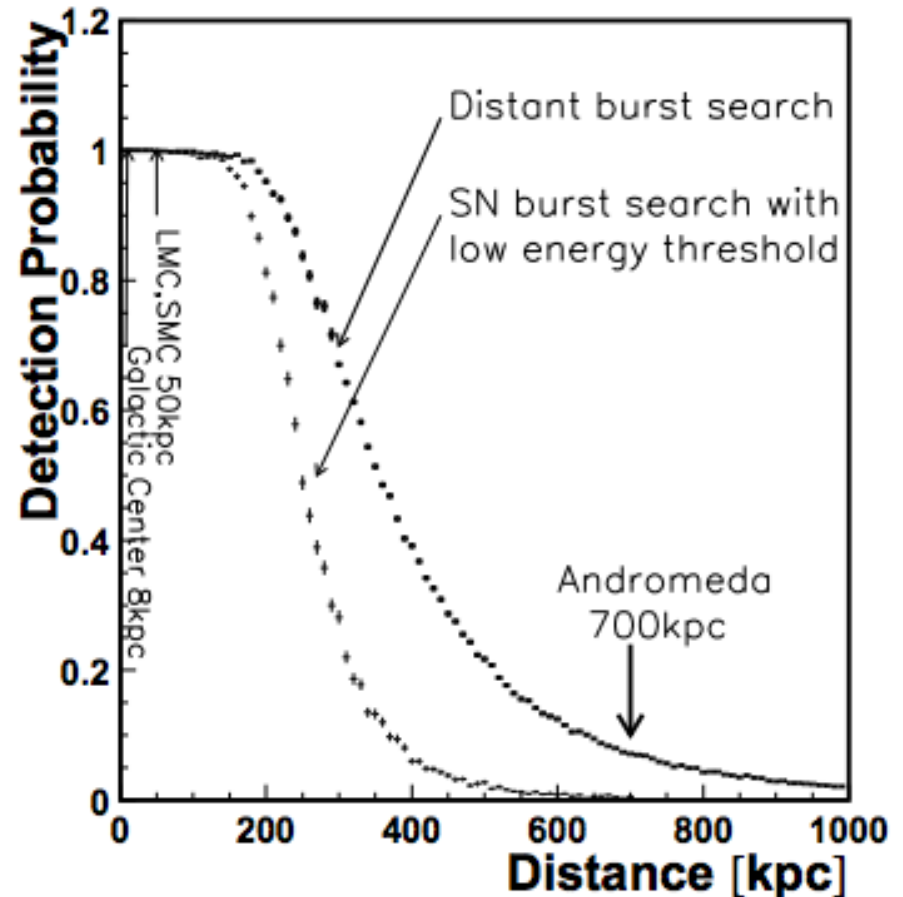
Optical counterpart can appear within minutes of neutrino alert → take full advantage of once-in-a-lifetime event



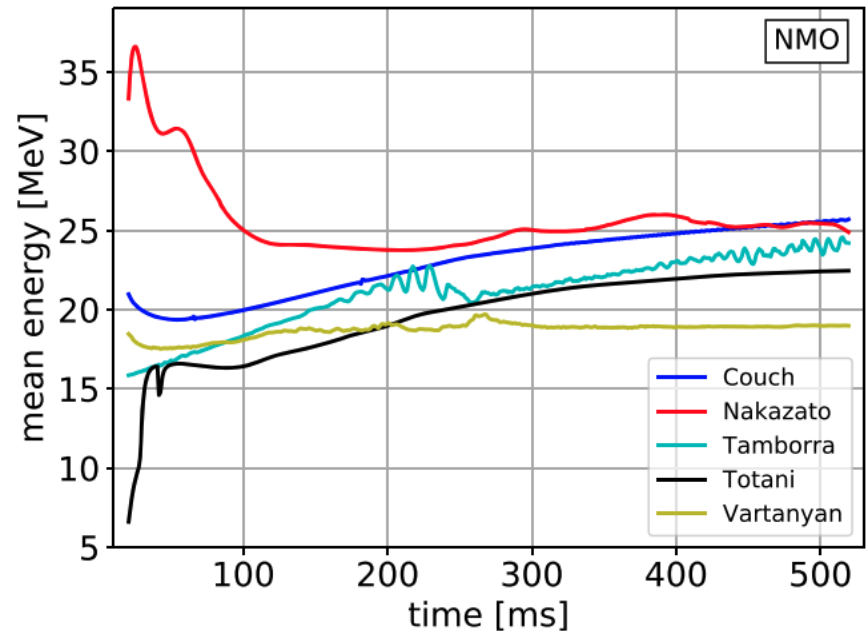
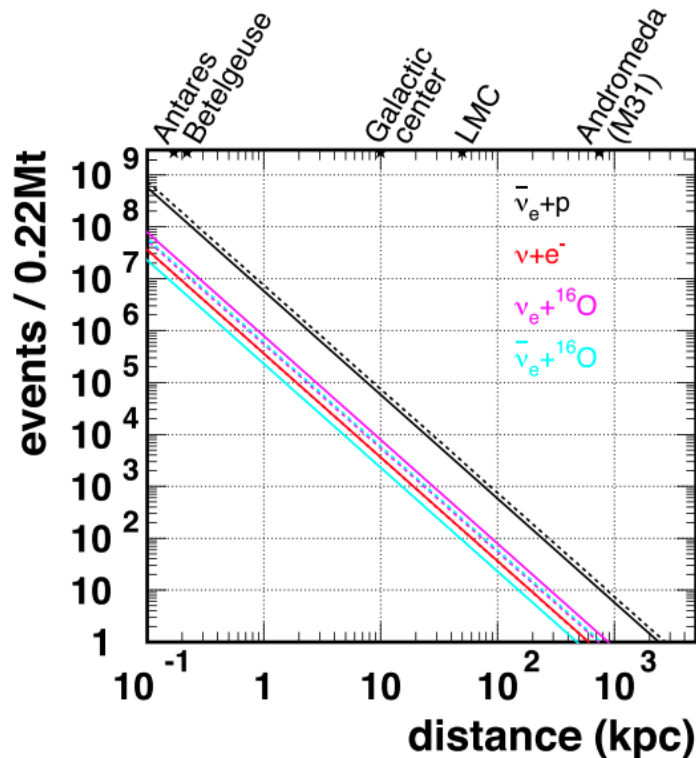
Sensitivity to Supernovae today



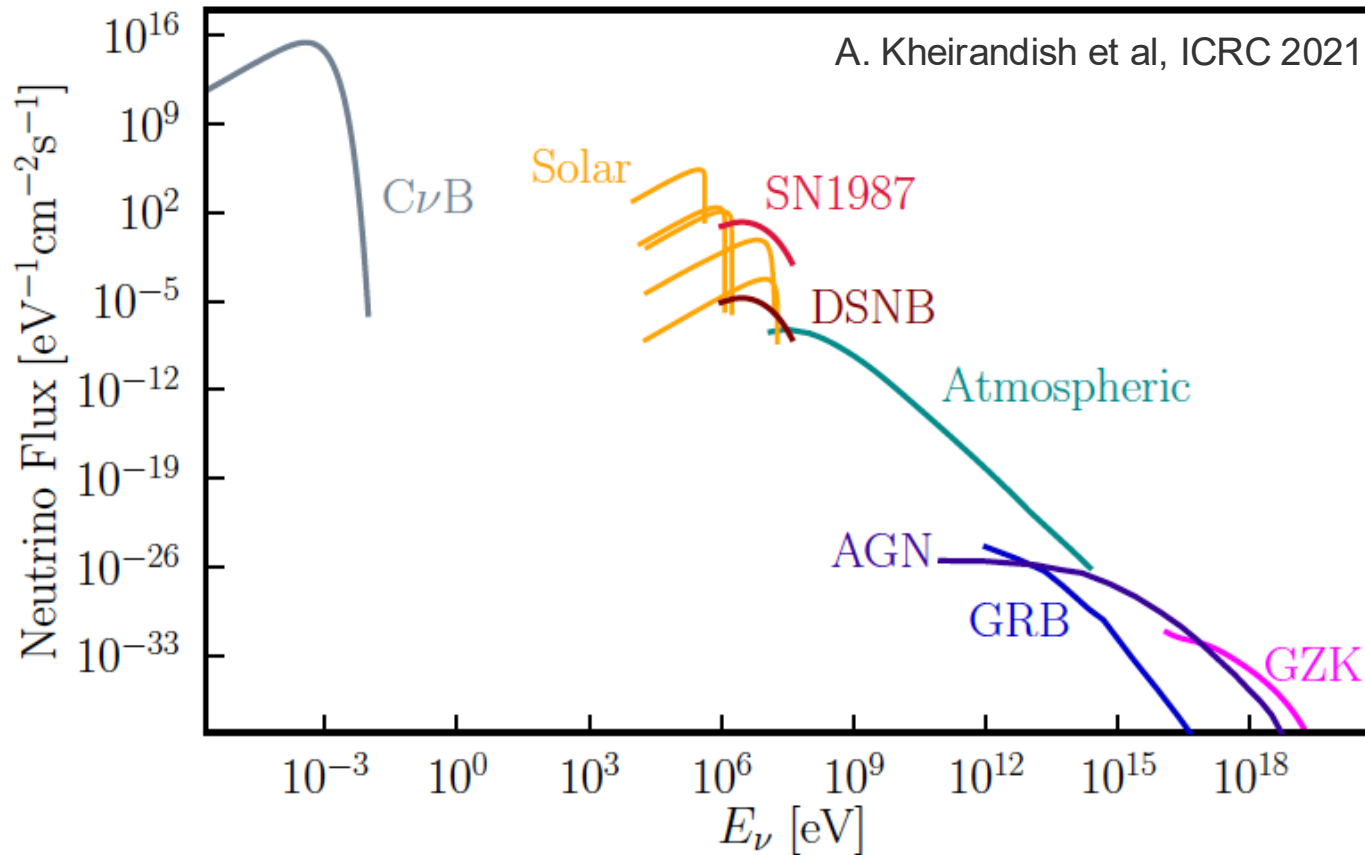
Super-Kamiokande



In the future: Hyper-Kamiokande (~2030)



Distinguish different supernova models



↔

Ideas: neutrino capture
on unstable nucleus
(e.g. tritium)

↔

↔

Neutrino capture on
stable nucleus
(Chlorine, Gallium)

↔

Water / ice
Cherenkov
detectors

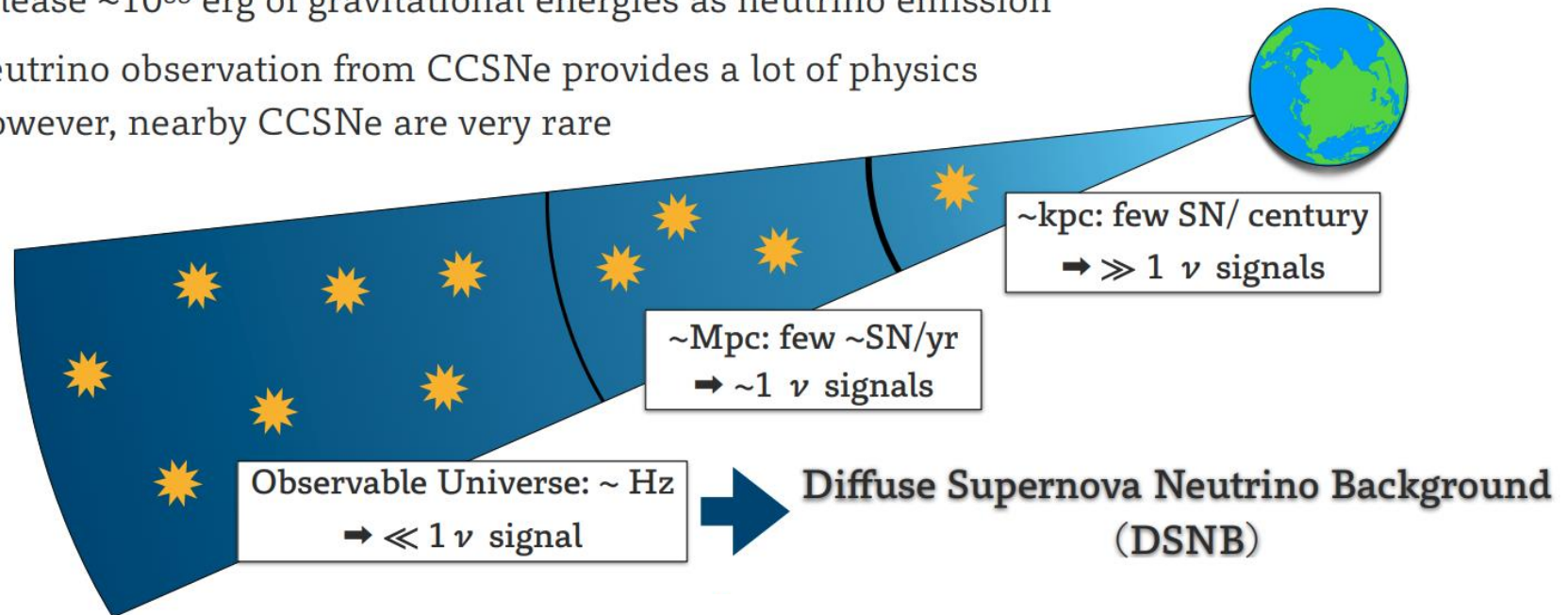
↔

Radio arrays,
cosmic-ray
detectors

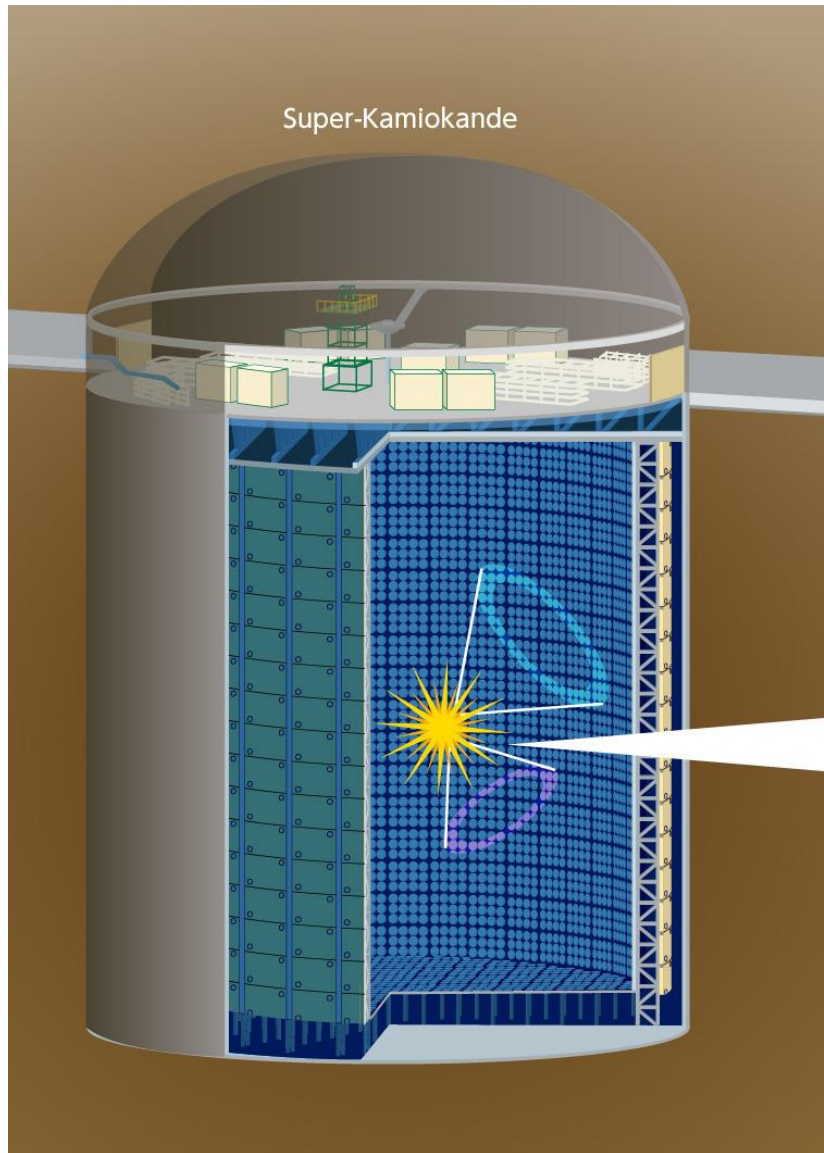
Diffuse Supernova Background (DSNB)

Core-Collapse Supernova (CCSN)

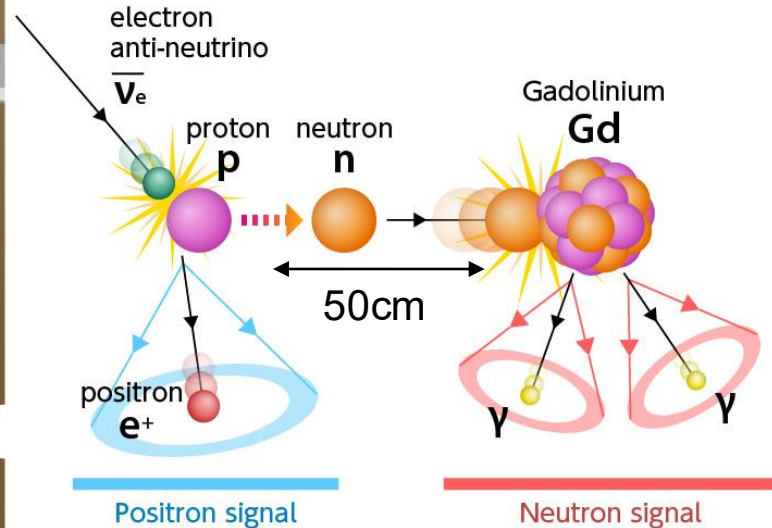
- Release $\sim 10^{53}$ erg of gravitational energies as neutrino emission
 - Neutrino observation from CCSNe provides a lot of physics
- However, nearby CCSNe are very rare



Super Kamiokande with Gadolinium



Gadolinium: highest affinity for capturing neutrons among all elements in nature.

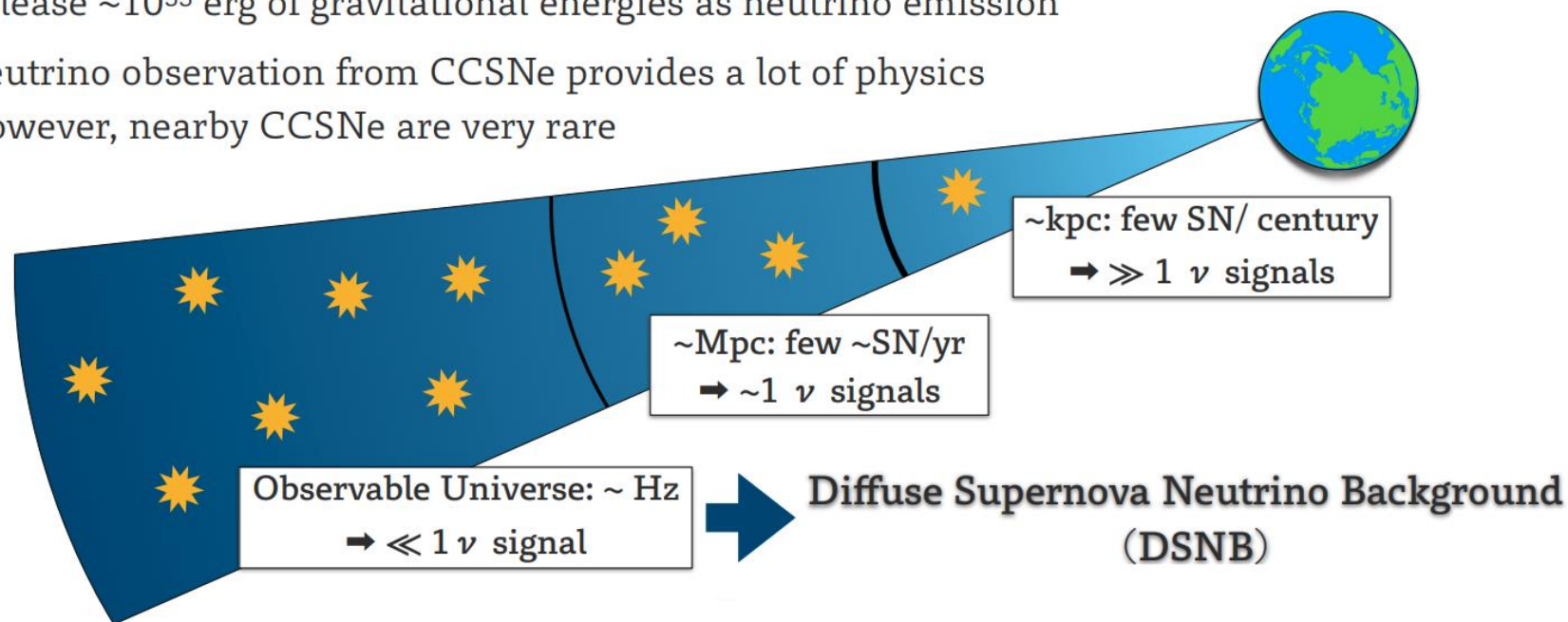


0.1% concentration of Gd: 90% of neutrons will be captured
→ highly efficient background suppression

Diffuse Supernova Background (DSNB)

Core-Collapse Supernova (CCSN)

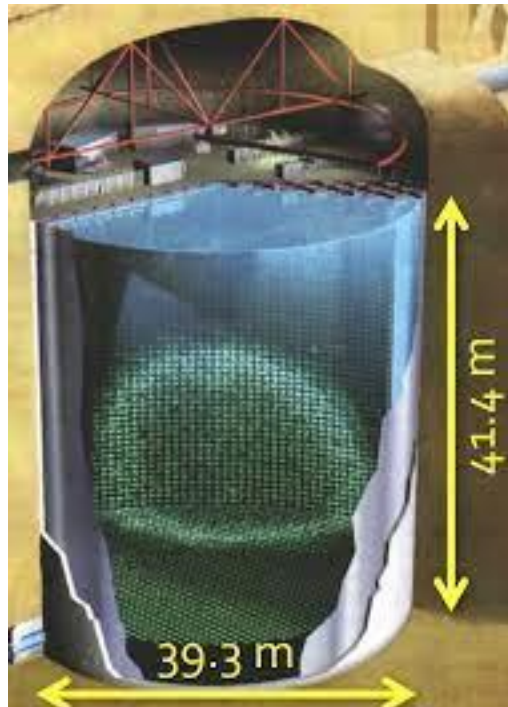
- Release $\sim 10^{53}$ erg of gravitational energies as neutrino emission
 - Neutrino observation from CCSNe provides a lot of physics
- However, nearby CCSNe are very rare



- Also, the latest update of DSNB search in SK-Gd using additional more condensed Gd-water data are exhibited
 - $\rightarrow$ There is no significant DSNB signal, however, some excess appears to be visible in the signal region, which is **2.3 σ tension from non-DSNB hypothesis**
- Looking forward to discovery of DSNB in the next decade !!

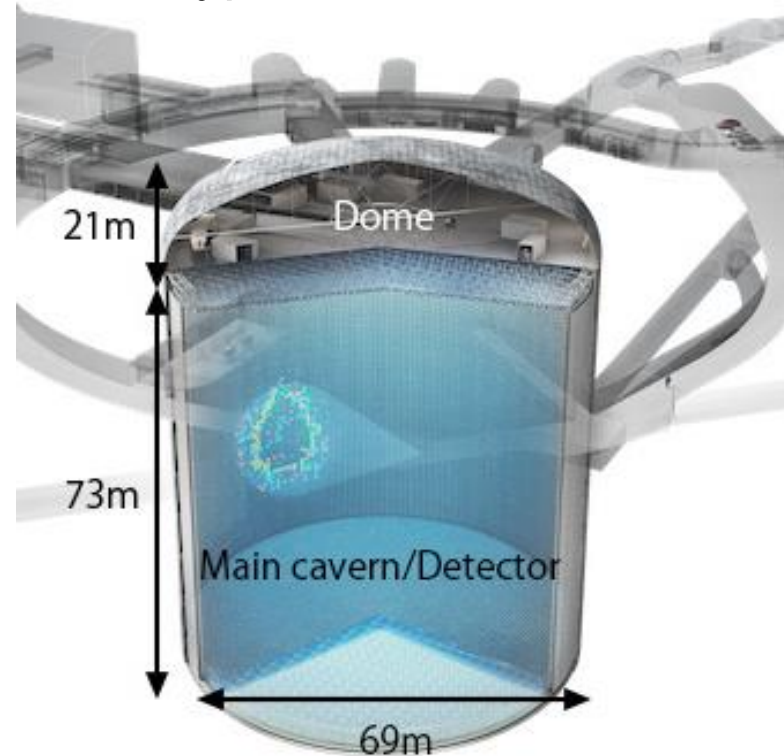
Next Generation Detector: Hyper Kamiokande

Super Kamiokande



Volume: 50kT

Hyper Kamiokande



Volume: 260kT

Data taking start in 2027



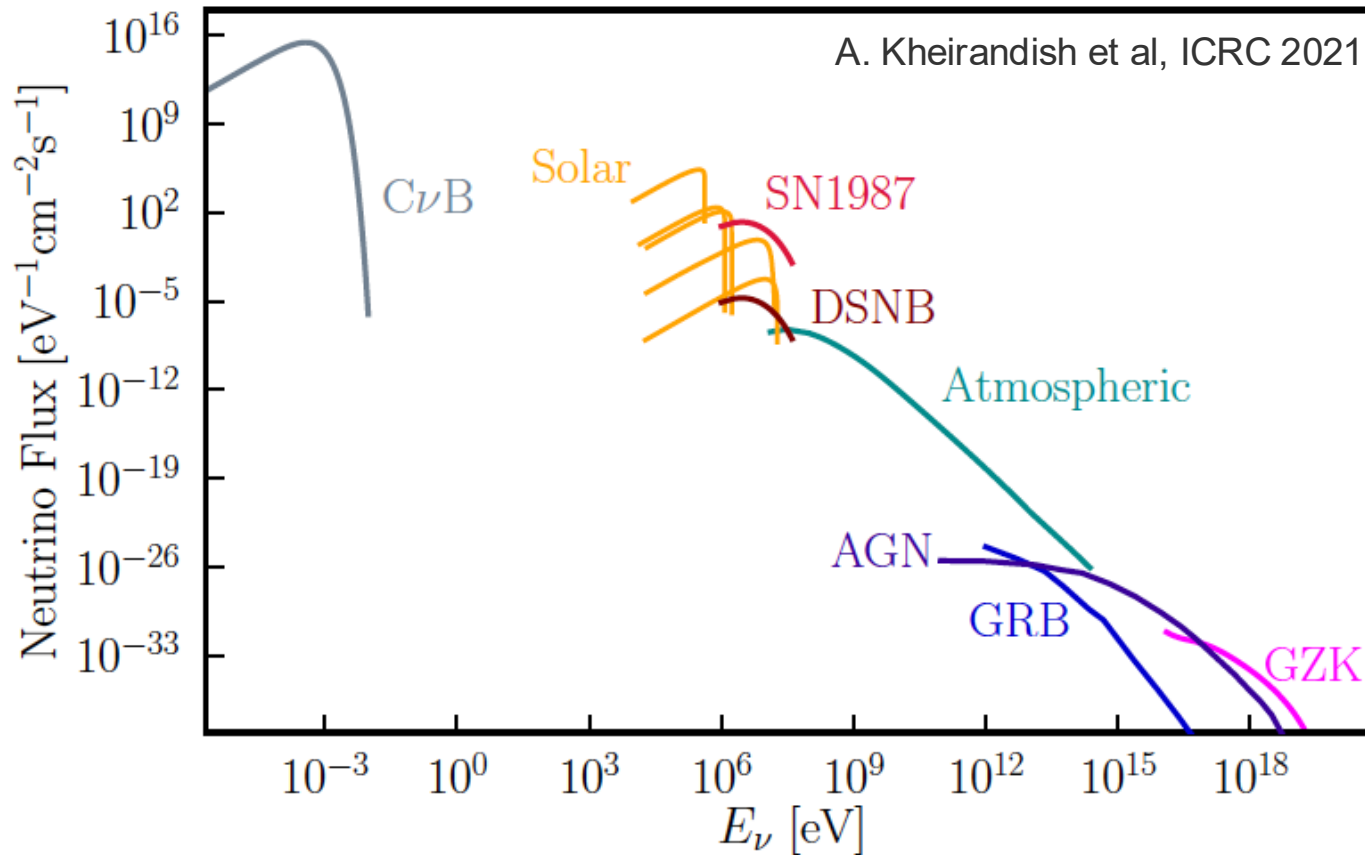
Take home message

A short burst of $\sim 10\text{MeV}$ neutrinos is expected from a core-collapse supernova and can be detected for Galactic source, expected only 1-3 times per 100 years. SNEWS connects many instruments and will send an alert to trigger observations.

Questions?!



<https://forms.gle/uHCYKcGDBP2bGRsC6>



↔

Ideas: neutrino capture
on unstable nucleus
(e.g. tritium)

↔

↔

Neutrino capture on
stable nucleus
(Chlorine, Gallium)

↔

↔

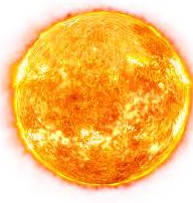
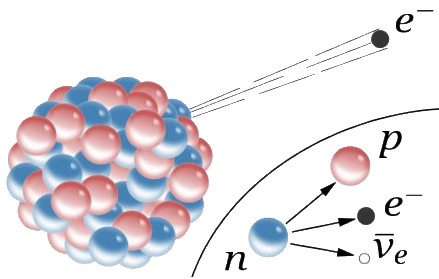
Water / ice
Cherenkov
detectors

↔

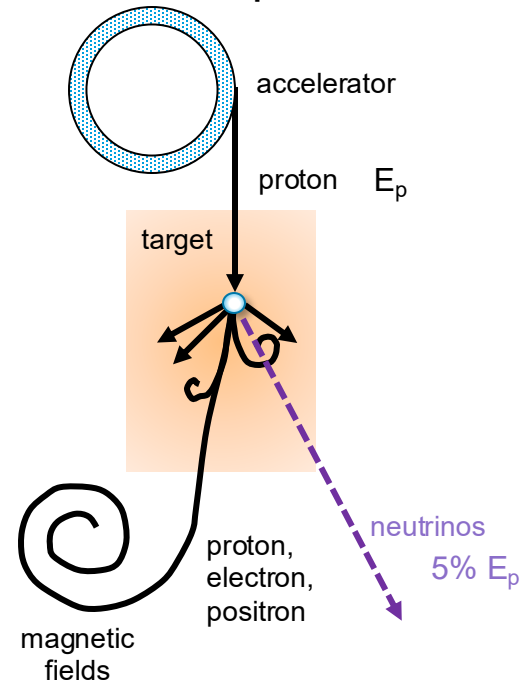
Radio arrays,
cosmic-ray
detectors

How are high-energy neutrinos produced?

MeV neutrinos from nuclear processes, (inverse) beta decay



TeV-PeV neutrinos from cosmic-ray “beam dumps”



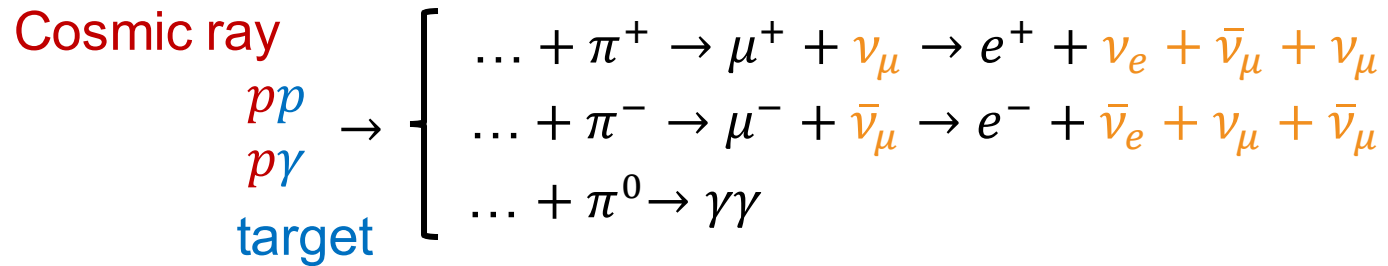
Two ingredients:

- Proton acceleration
- Target for interaction

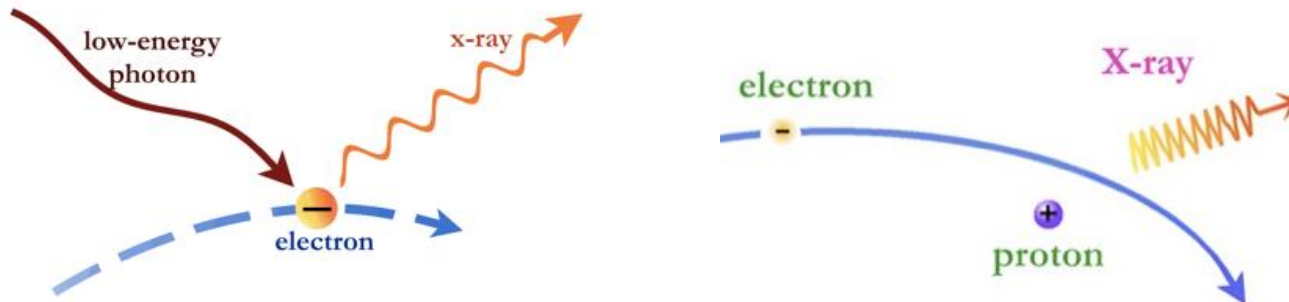
Neutrino Production Processes

$$\begin{array}{l}
 \text{Cosmic ray} \\
 \text{\textit{pp}} \\
 \text{\textit{p}\gamma} \\
 \text{target}
 \end{array}
 \rightarrow \left\{ \begin{array}{l}
 \dots + \pi^+ \rightarrow \mu^+ + \nu_\mu \rightarrow e^+ + \nu_e + \bar{\nu}_\mu + \nu_\mu \\
 \dots + \pi^- \rightarrow \mu^- + \bar{\nu}_\mu \rightarrow e^- + \bar{\nu}_e + \nu_\mu + \bar{\nu}_\mu \\
 \dots + \pi^0 \rightarrow \gamma\gamma
 \end{array} \right.$$

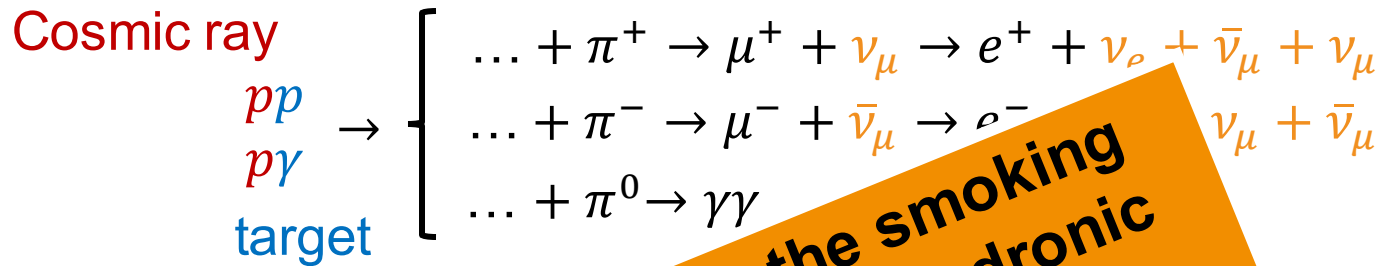
Neutrino Production Processes



Gamma-rays are not exclusively produced in hadronic processes



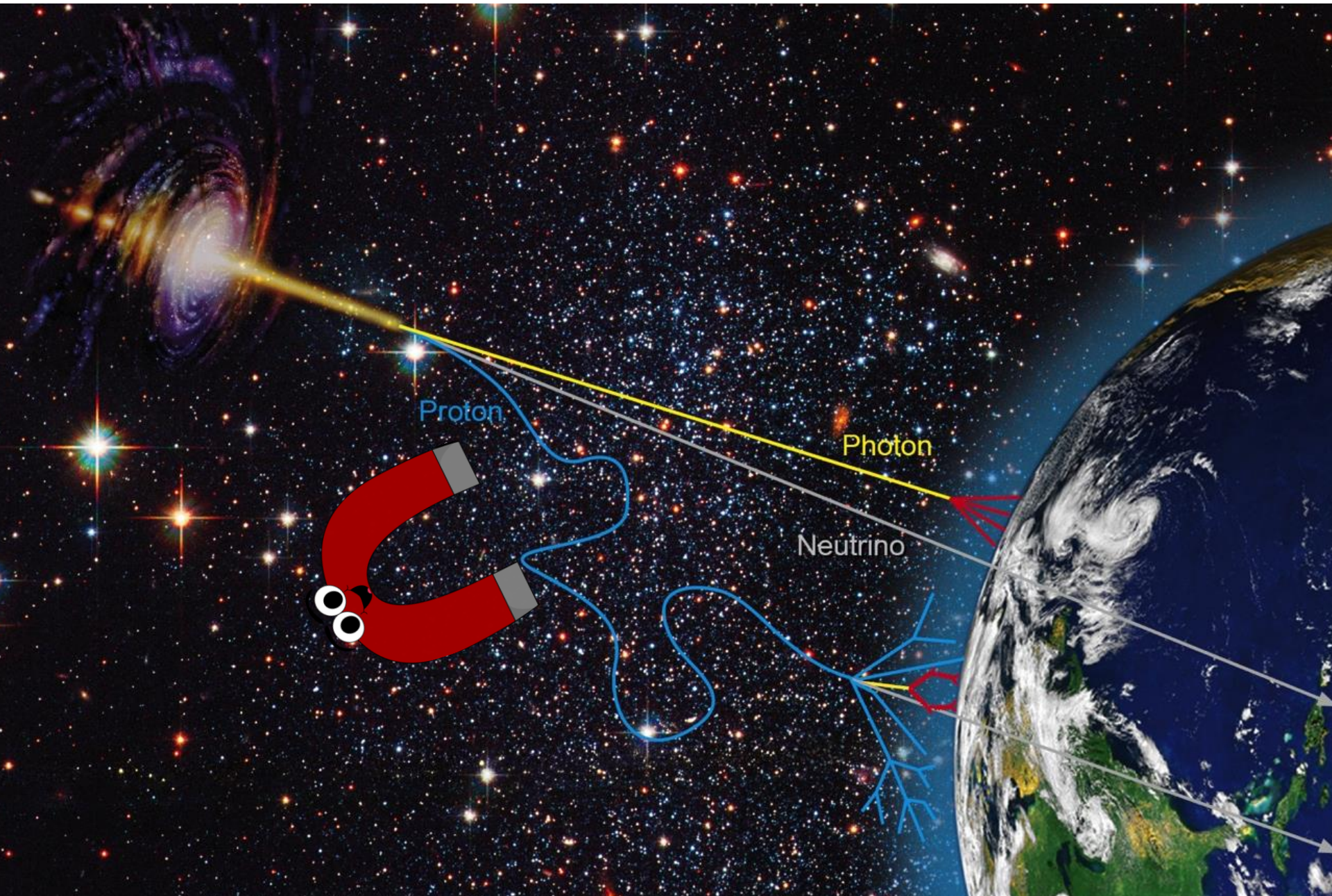
Neutrino Production Processes



**Neutrinos are the smoking
gun signature for hadronic
acceleration**



What are the Cosmic-Ray Sources?

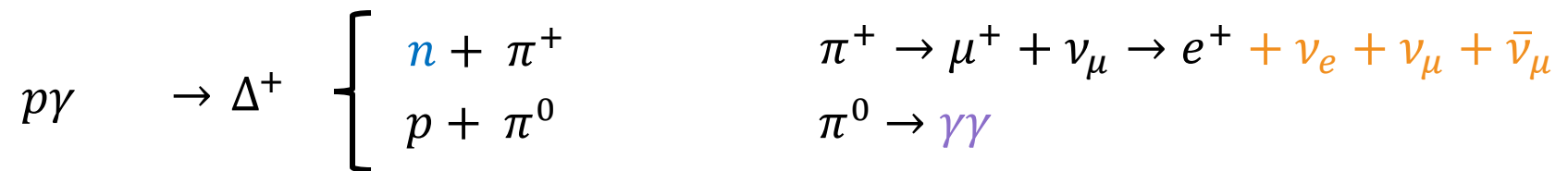


High-energy Neutrino Detectors

- Huge volumes necessary: $\sim 1\text{km}^3 \rightarrow$ need to use natural medium
- Transparent medium (Cherenkov emission of secondary charged particles)

Waxman-Bahcall Bound

- Protons are accelerated in the sources with an energy spectrum proportional to $dN_{\text{CR}}/dE_{\text{CR}} \propto E^{-2}$
- Energy production rate in CRs between 10^{19}eV and 10^{21}eV is $5 \times 10^{44} \text{ erg Mpc}^{-3} \text{ yr}^{-1}$
- All protons undergo photohadronic interactions ($p\gamma$), producing **neutrons**, **gamma rays** and **neutrinos**



Waxman-Bahcall Bound

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- Energy production rate in CRs between 10^{19}eV and 10^{21}eV is $5 \times 10^{44} \text{ erg Mpc}^{-3} \text{ yr}^{-1}$
- All protons undergo photohadronic interactions (py), producing neutrons, gamma rays and neutrinos
- The sources are optically thin to neutrons, which leave the sources and subsequently decay into protons, accounting for the UHECR flux observed at the Earth
- Energy flux of neutrinos cannot be greater than that of CR

$$E^2 \Phi < 3 \times 10^{-8} \text{ GeV s}^{-1} \text{ sr}^{-1} \text{ cm}^{-2}$$

1km³ volume needed to probe this flux



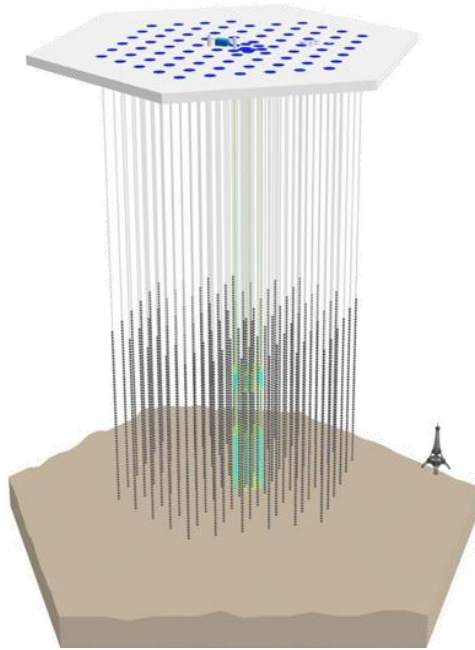
Take home message

Kubic-kilometer sized neutrino detectors are needed to detect cosmic neutrinos.

Those could reveal the sources of cosmic rays and deliver information from dense environments

High-energy Neutrino Detectors

IceCube (1 km³)

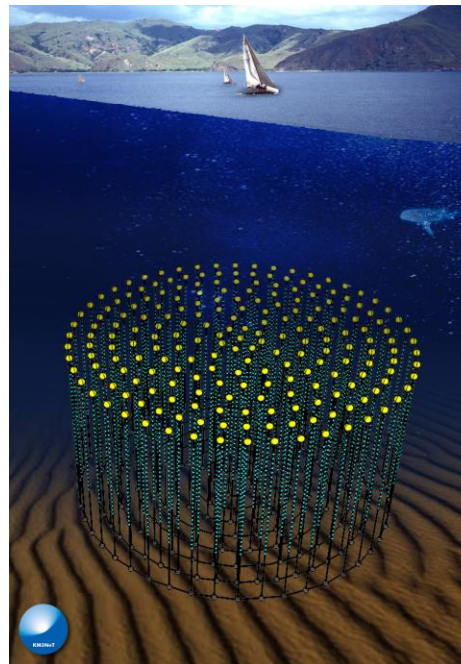


IceCube-Gen2 (8km³) planned

M.G. Aartsen *et al* 2017 *JINST* **12** P03012
 M.G. Aartsen *et al* *J.Phys.G*
 48 (2021) 6, 060501

KM3NeT (2 x 0.5 km³)

Under construction
 Predecessor ANTARES
 decommissioned after 14y

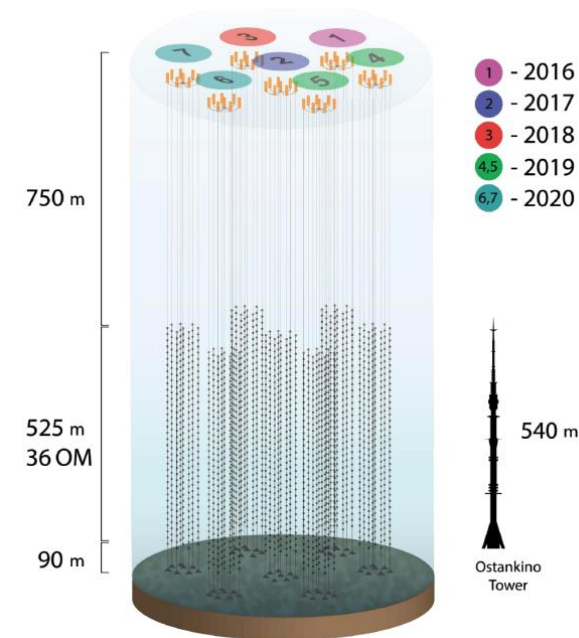


Two sites: ARCA for HE
 and ORCA for LE,
 roughly ~20% complete

S. Adrián-Martínez et al. arXiv:1601.07459

Baikal-GVD (1 km³)

Under construction



2022: 10 clusters
 2026: 18 clusters

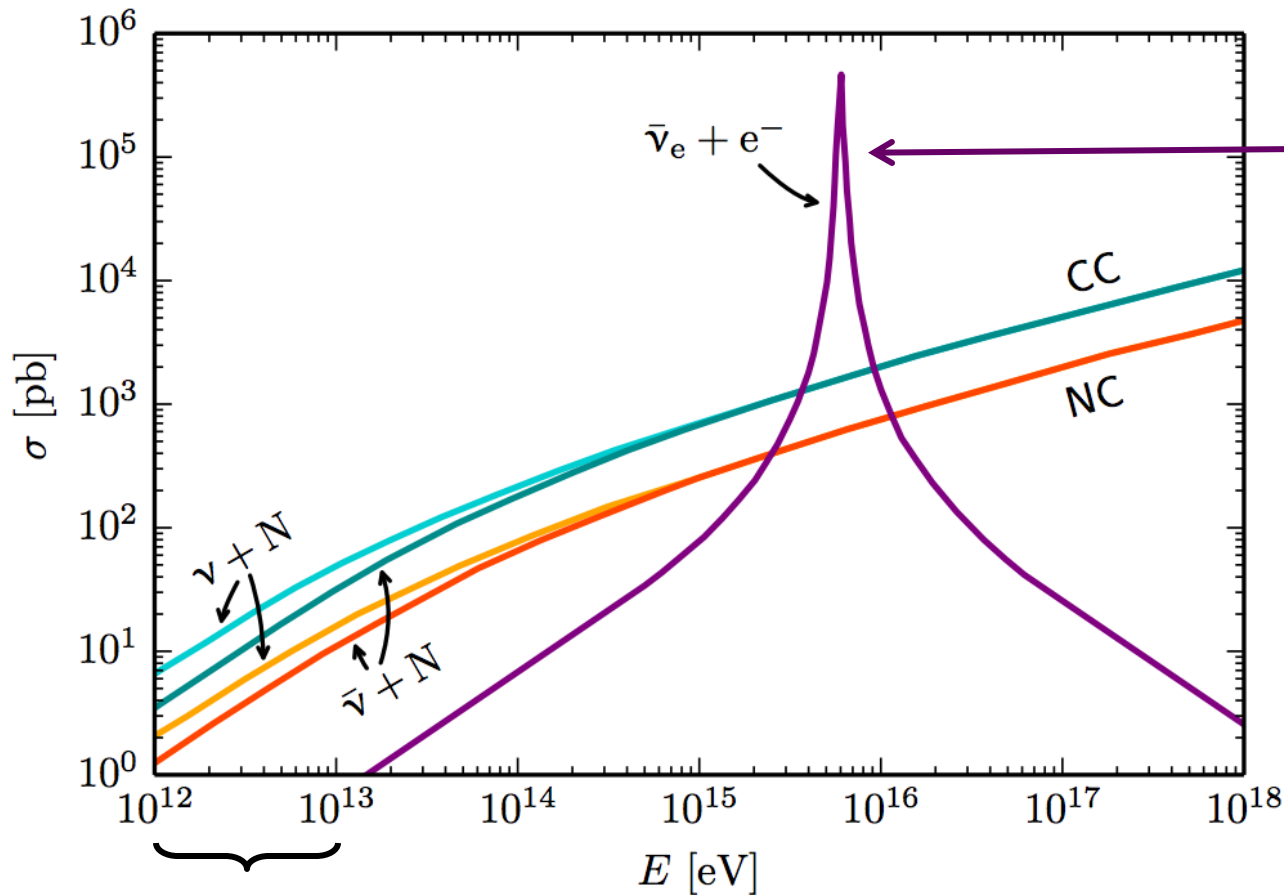
A.D. Avrorin et al.
 arXiv:2011.09209

Differences between ice and water detectors

Property	Lake Baikal	Mediterranean (ANTARES)	Antarctic Ice
Absorption length (m)	22	60	100
Effective Scattering (m)	480	265	25
Depth	1370	2475	2450
Noise	Quiet	^{40}K , bioluminescence	Quiet
Retrieve/redeploy	Yes	Yes	No

Long scattering length in Mediterranean implies better angular resolution; long absorption length for IceCube allows sparser instrumentation. Smaller depth implies larger atmospheric muon background.

Neutrino Cross-Sections



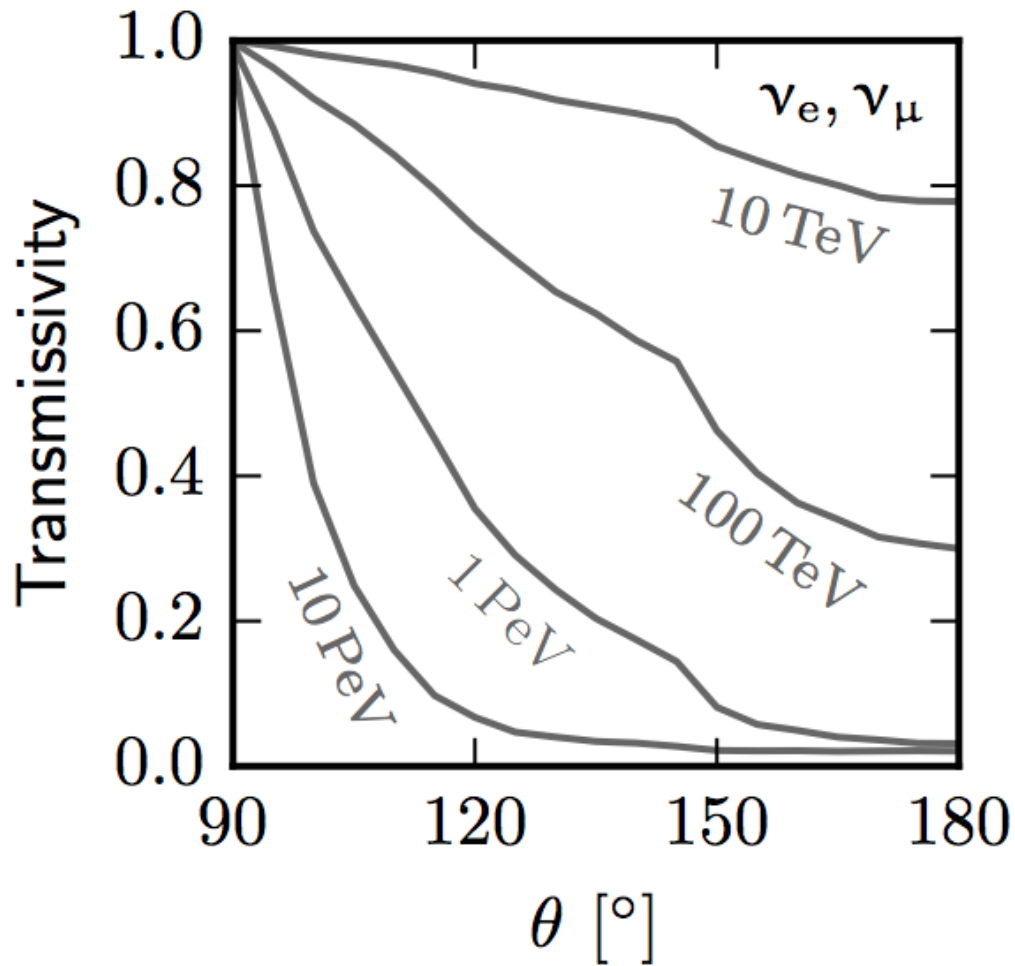
Linear with E ,
Smaller for
anti-neutrinos
due to helicity

damped by the W -
boson propagator

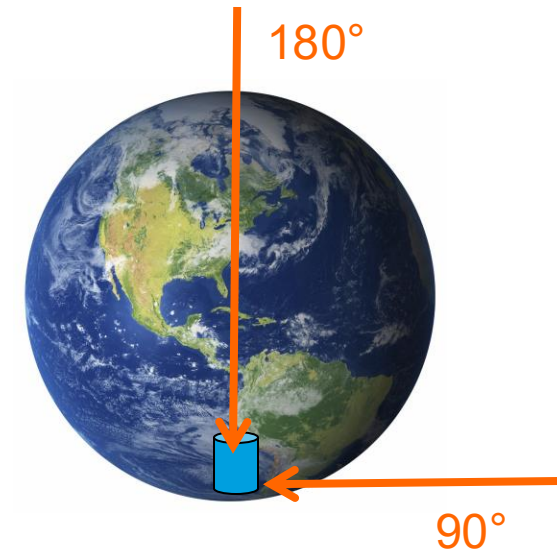
Glashow resonance:

resonance when the
center-of-mass
energy of the
system reaches the
mass of the
mediating boson.
For electrons at rest
and the mass of the
 $W^\pm$ boson (80 GeV),
the Glashow
resonance occurs at
a neutrino energy of
 6.3×10^{15} eV

Absorption of High-Energy Neutrinos in the Earth



For high-energies Earth becomes opaque to neutrinos





ICECUBE

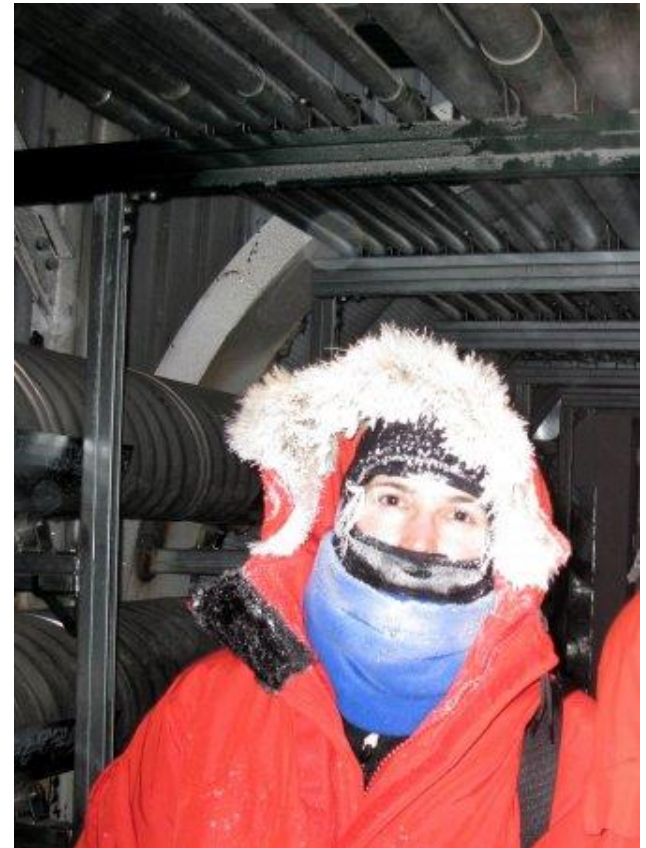
SOUTH POLE NEUTRINO OBSERVATORY



The South Pole

Elevation: 2,835 m

Average temperature: -28°C (summer), -60°C (winter)





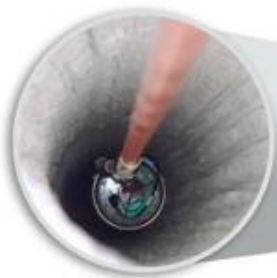
ICECUBE

SOUTH POLE NEUTRINO OBSERVATORY



IceCube Laboratory

Data is collected here and sent by satellite to the data warehouse at UW-Madison



Digital Optical Module (DOM)

5,160 DOMs deployed in the ice

50 m

IceTop

1450 m

2450 m

IceCube detector

86 strings of DOMs,
set 125 meters apart

DeepCore

Antarctic bedrock

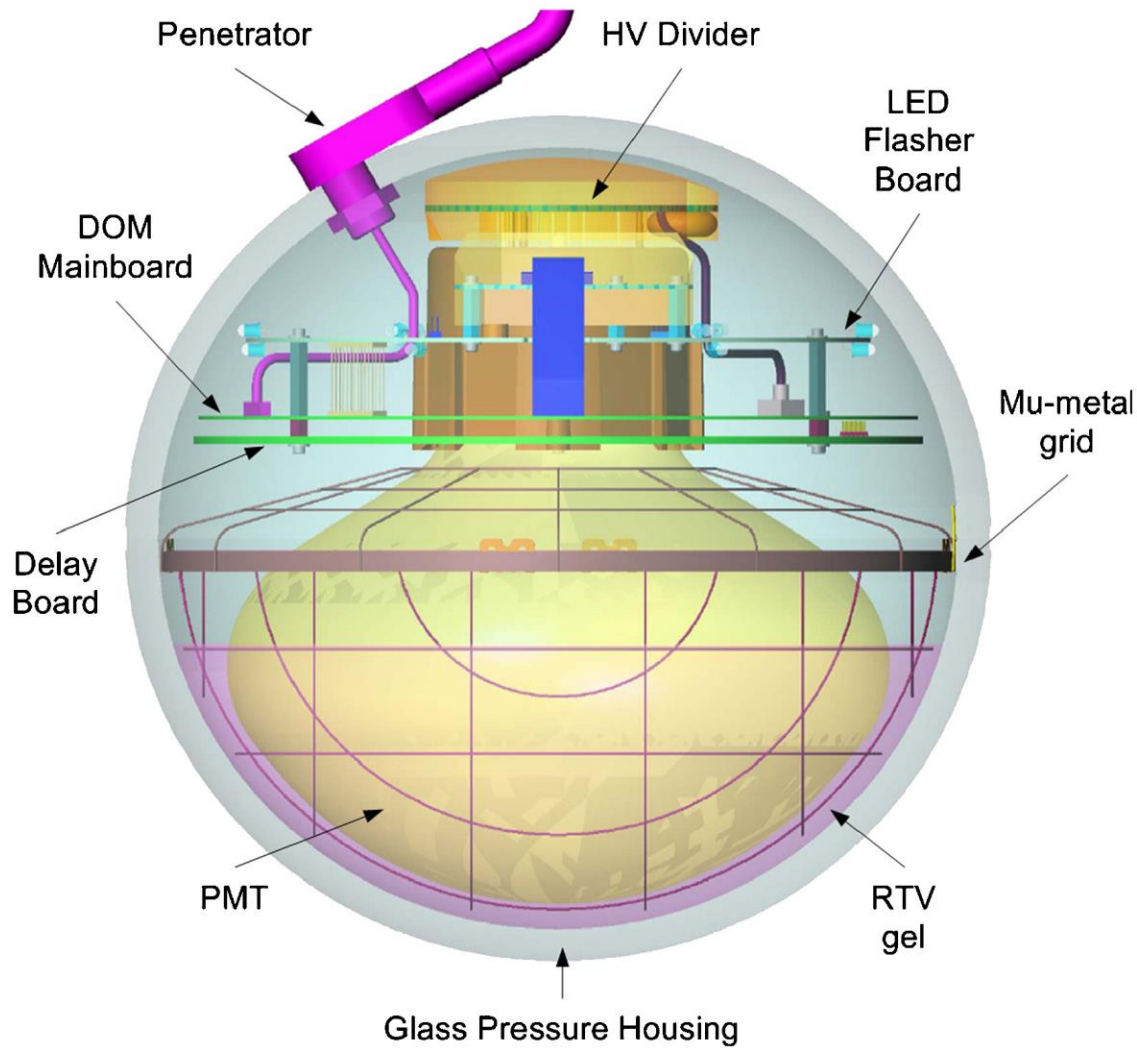
Amundsen-Scott South Pole Station, Antarctica

A National Science Foundation-managed research facility



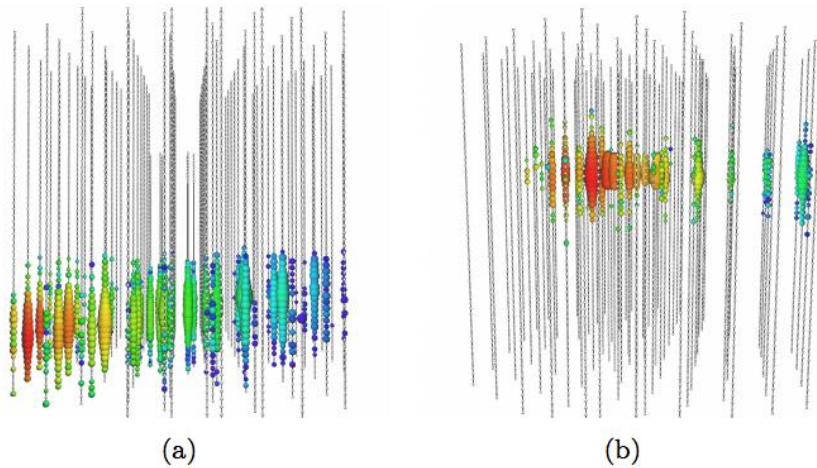
60 DOMs
on each
string

DOMs
are 17
meters
apart

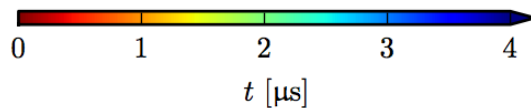
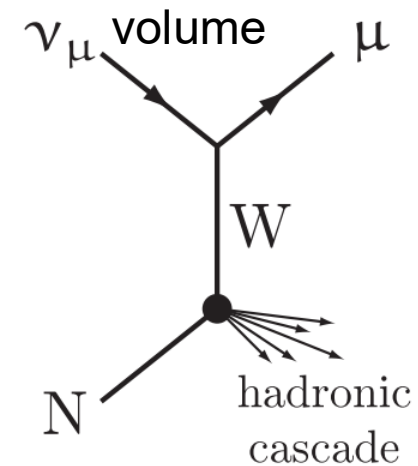


5160 DOMs
deployed
Failure rate: $<0.3\%$

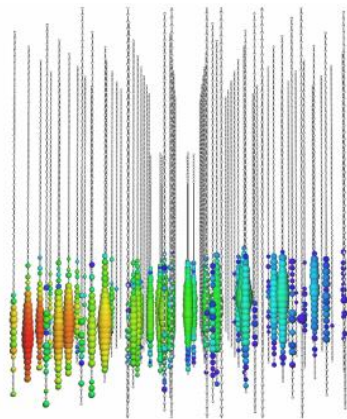
Event Signatures



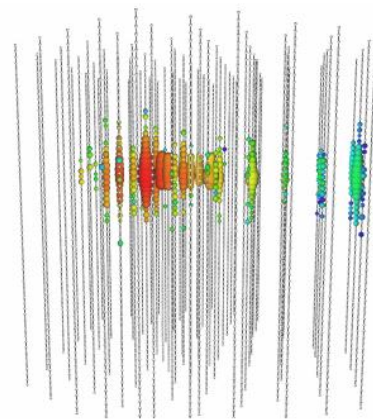
Charged current interaction of muon neutrino outside / inside the detector



Event Signatures

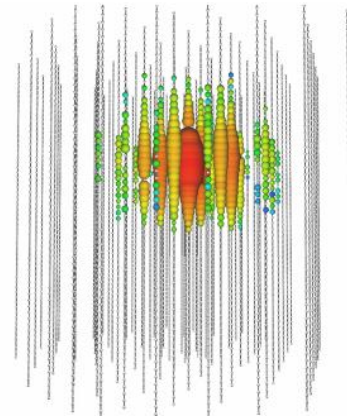


(a)



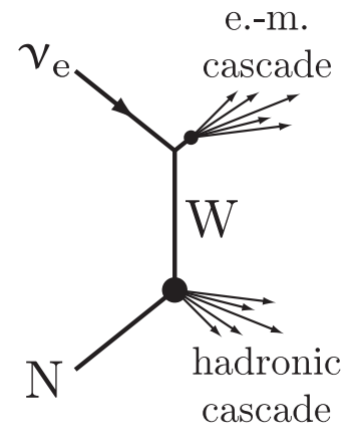
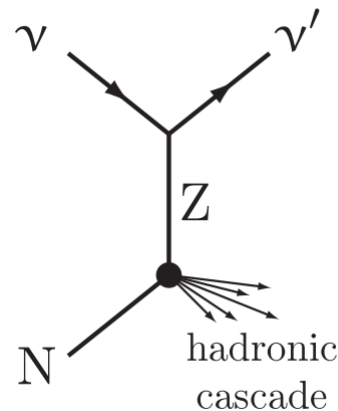
(b)

- a) through-going muon track $E \sim 140$ TeV
- b) Starting muon track $E \sim 70$ TeV
- c) **Shower event $E \sim 1$ PeV**

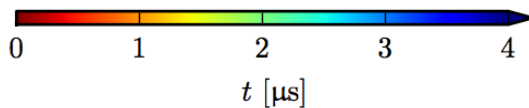


(c)

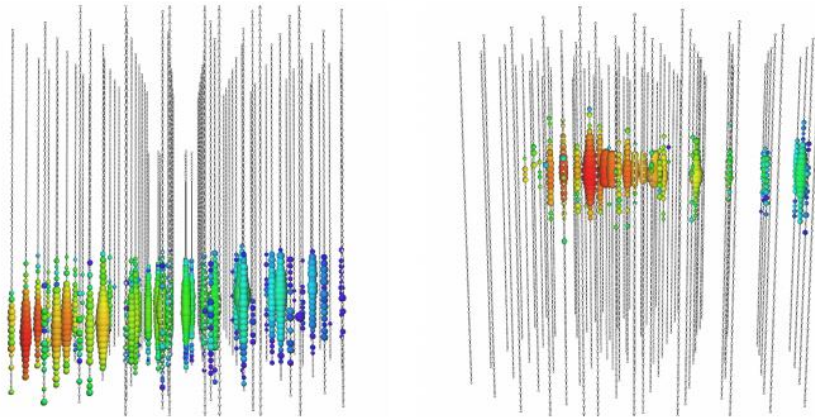
Neutral current or electron neutrino
charged current interaction



Cannot
distinguish
between
showers (size
few meters)



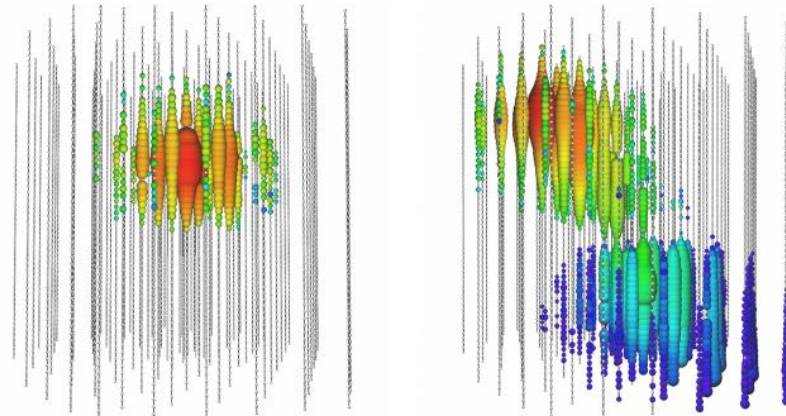
Event Signatures



(a)

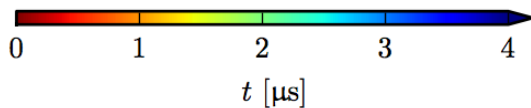
(b)

- a) through-going muon track $E \sim 140 \text{ TeV}$
- b) Starting muon track $E \sim 70 \text{ TeV}$
- c) Shower event $E \sim 1 \text{ PeV}$
- d) **“double bang” event $E \sim 200 \text{ PeV}$**

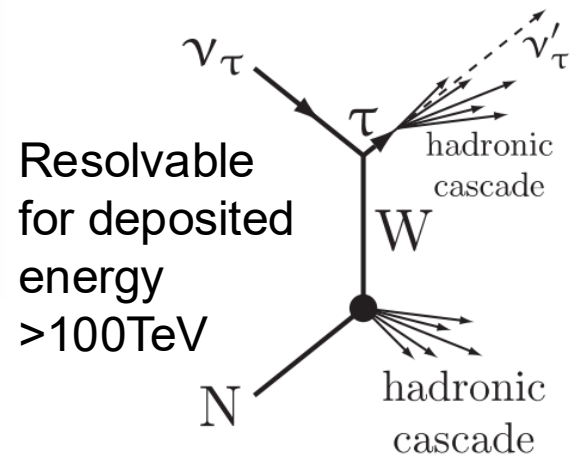


(c)

(d)

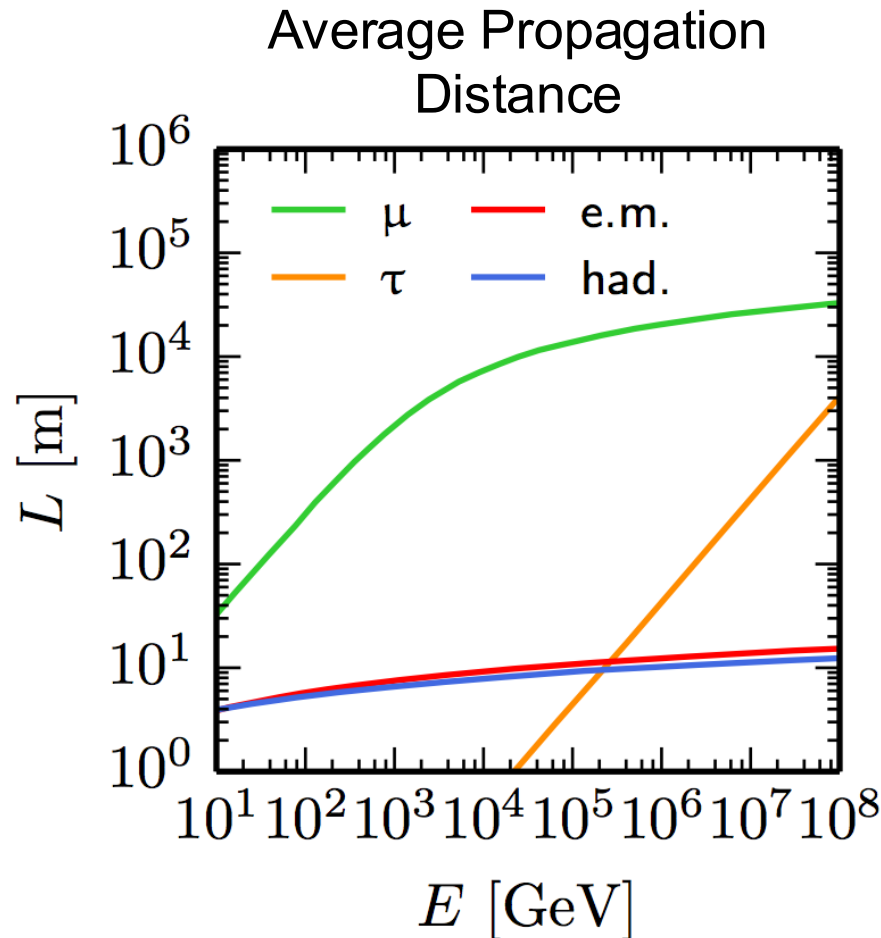


Tau neutrino charged current interaction



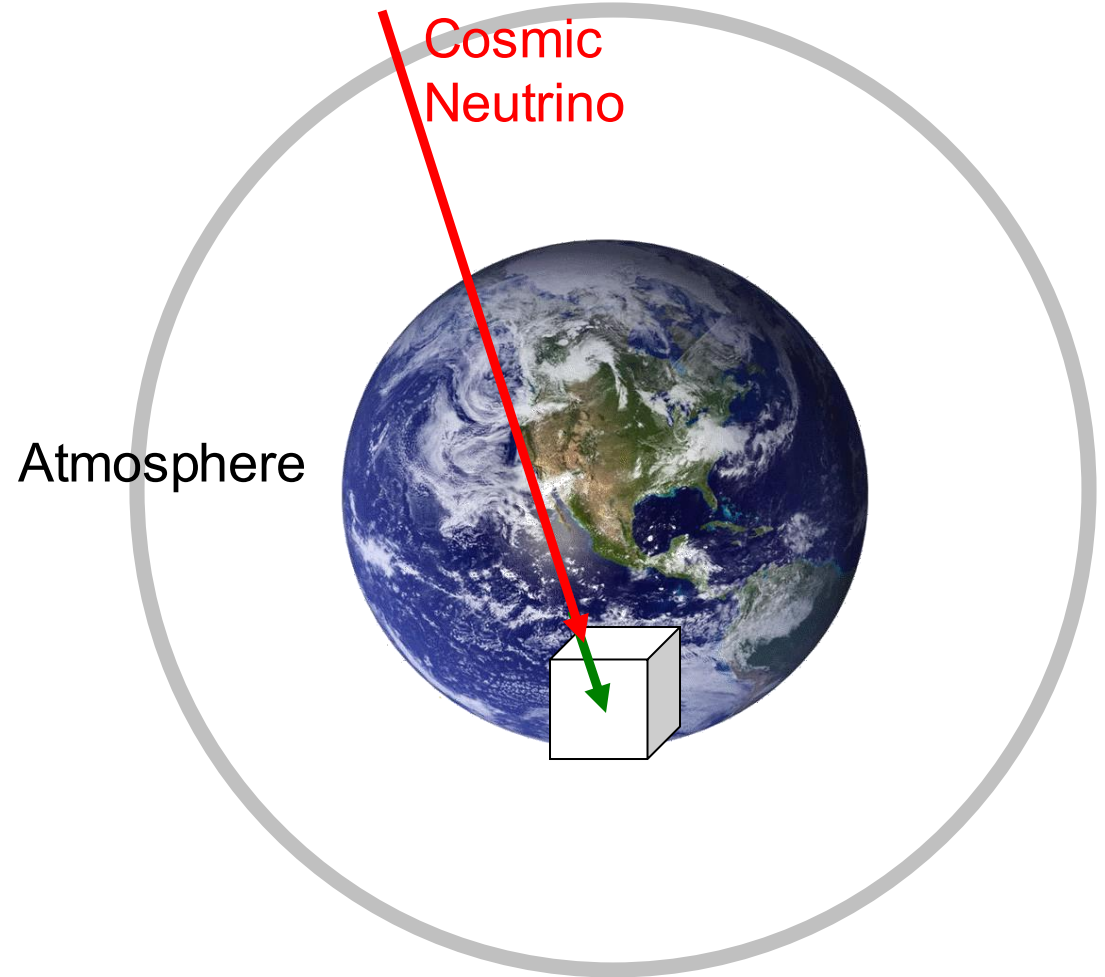
Only for very large energies the two showers can be separated (otherwise signature c)

Secondary Particle Interactions in Ice

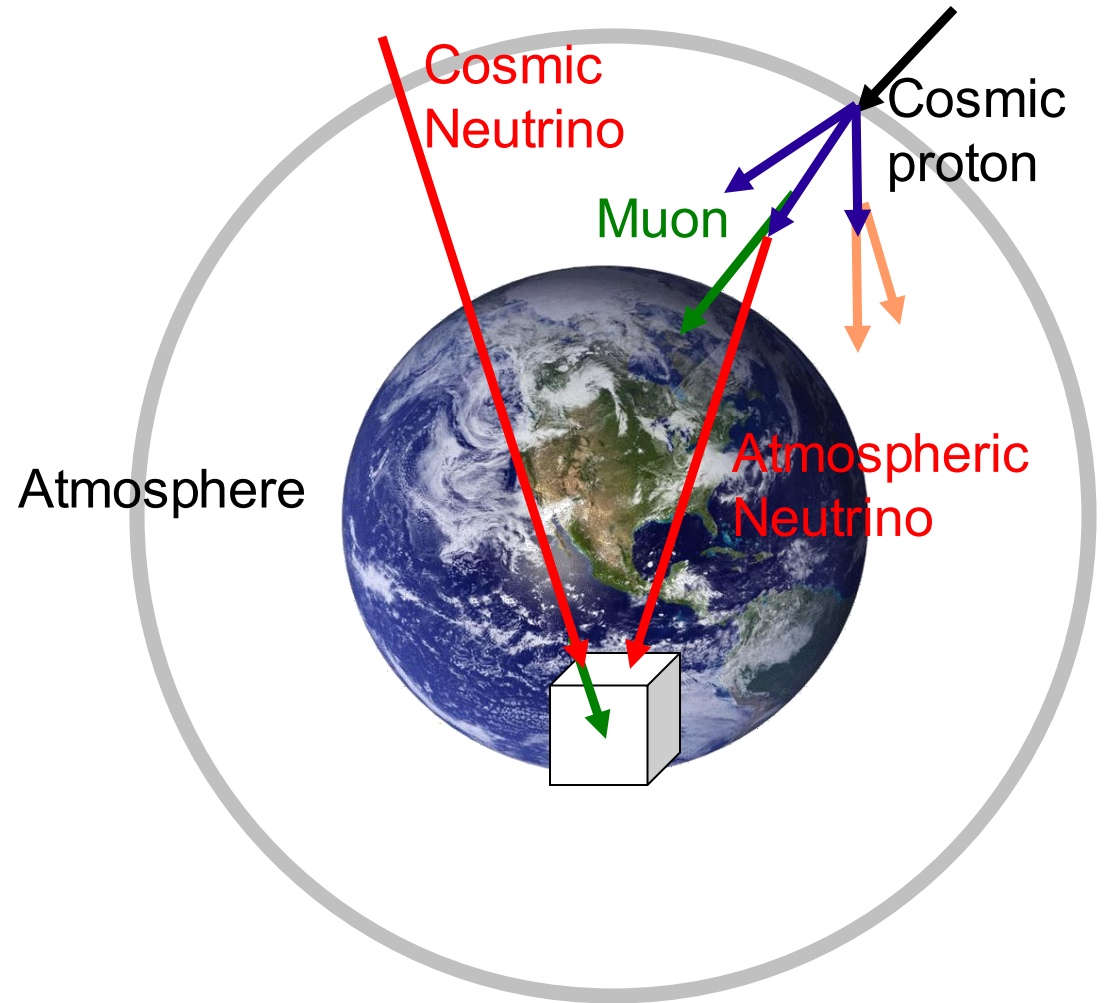


In the TeV range muons travel 1-10km, while showers have a small extension of few meters. Tau needs to reach >100TeV energies to reach detectable length of >10 meters.

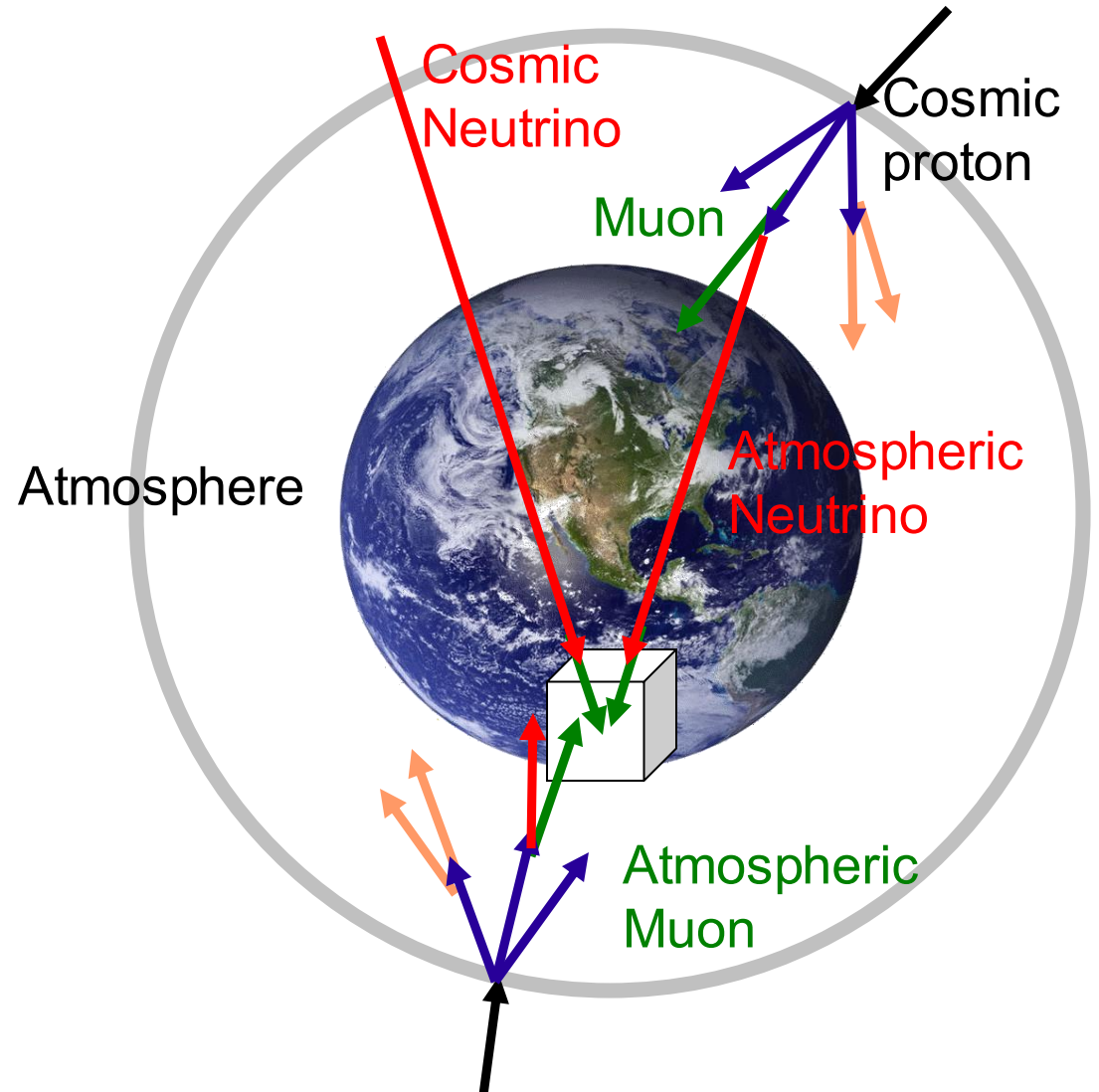
Background in Search for Cosmic Neutrinos



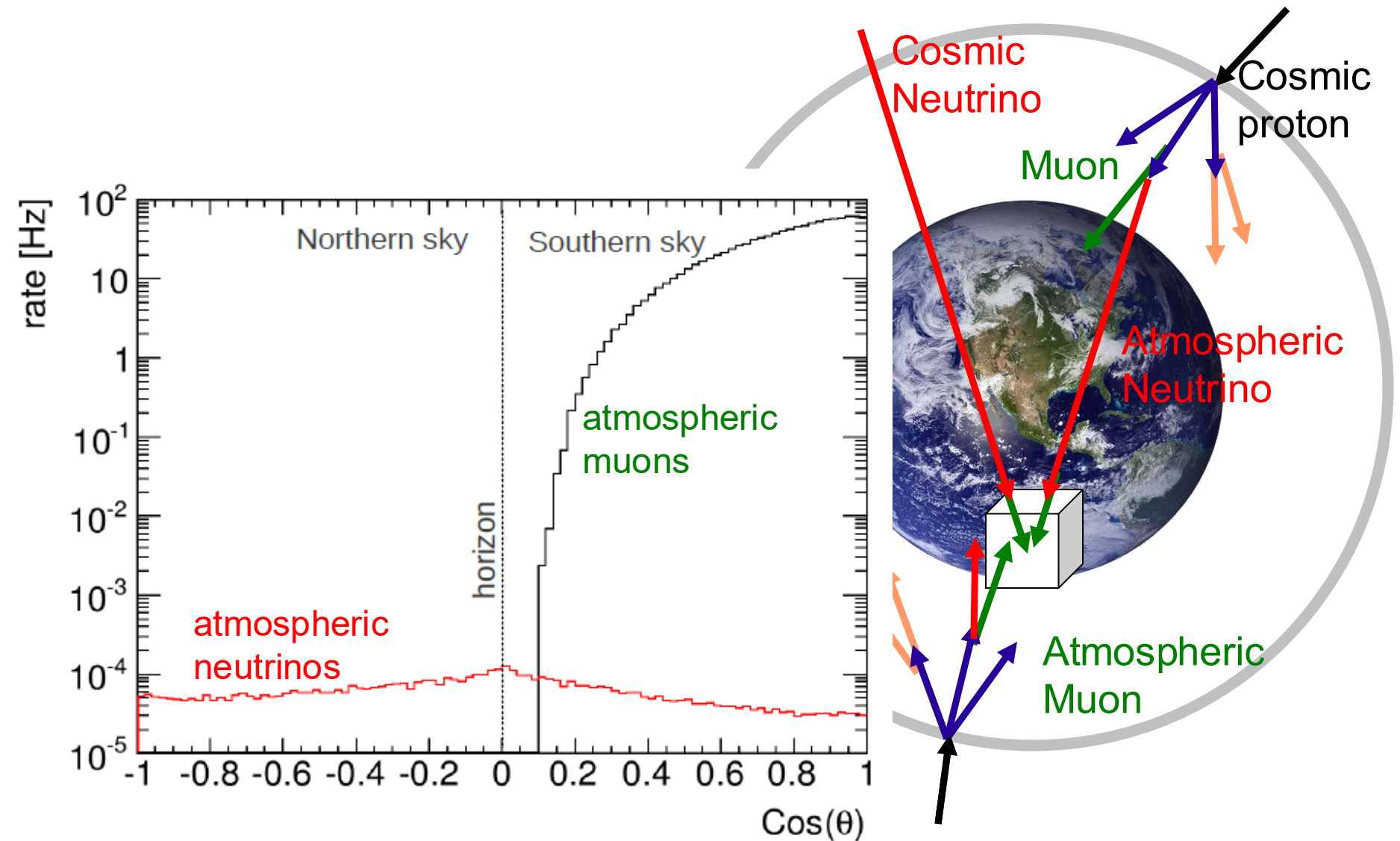
Background in Search for Cosmic Neutrinos



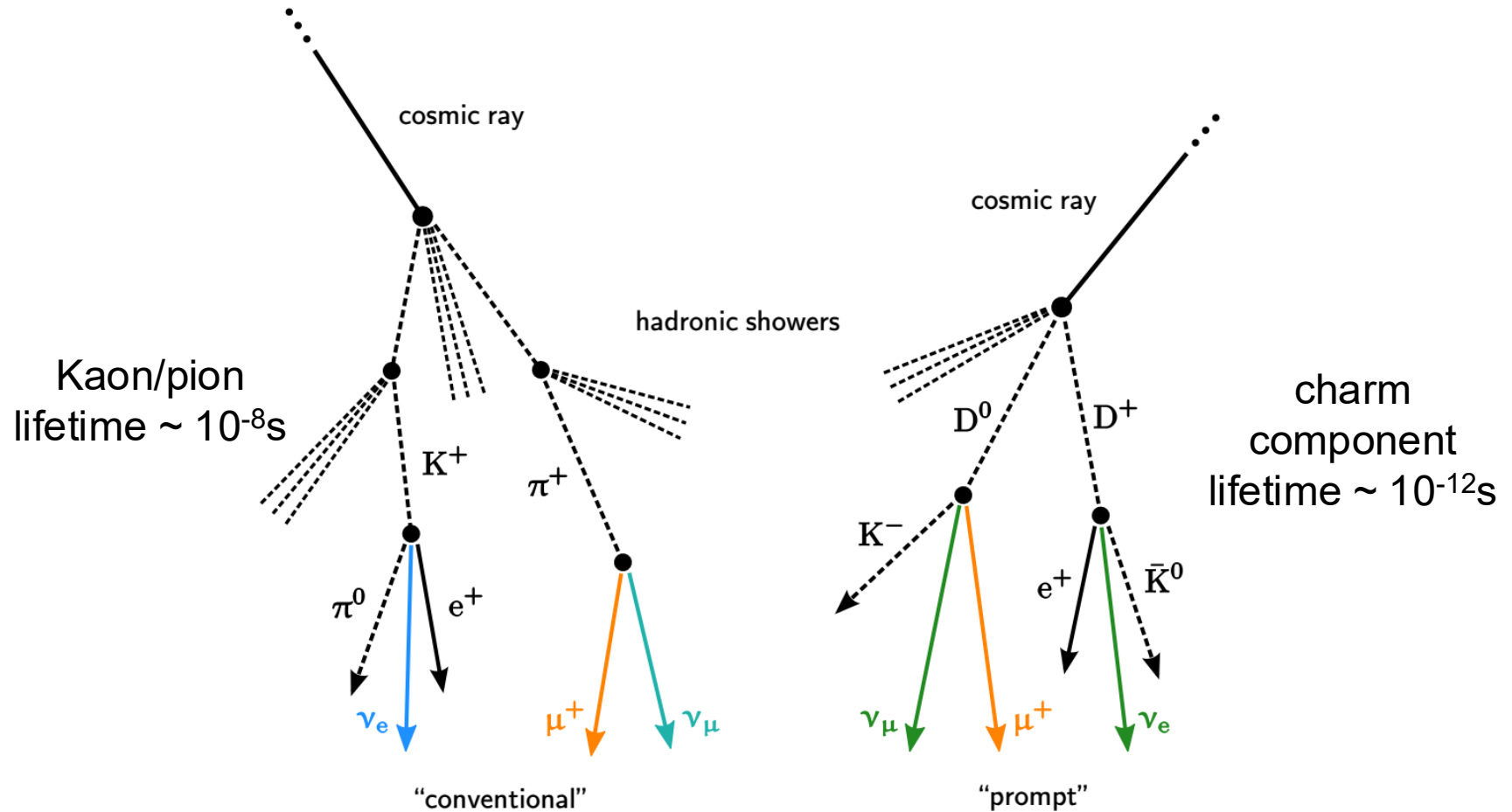
Background in Search for Cosmic Neutrinos



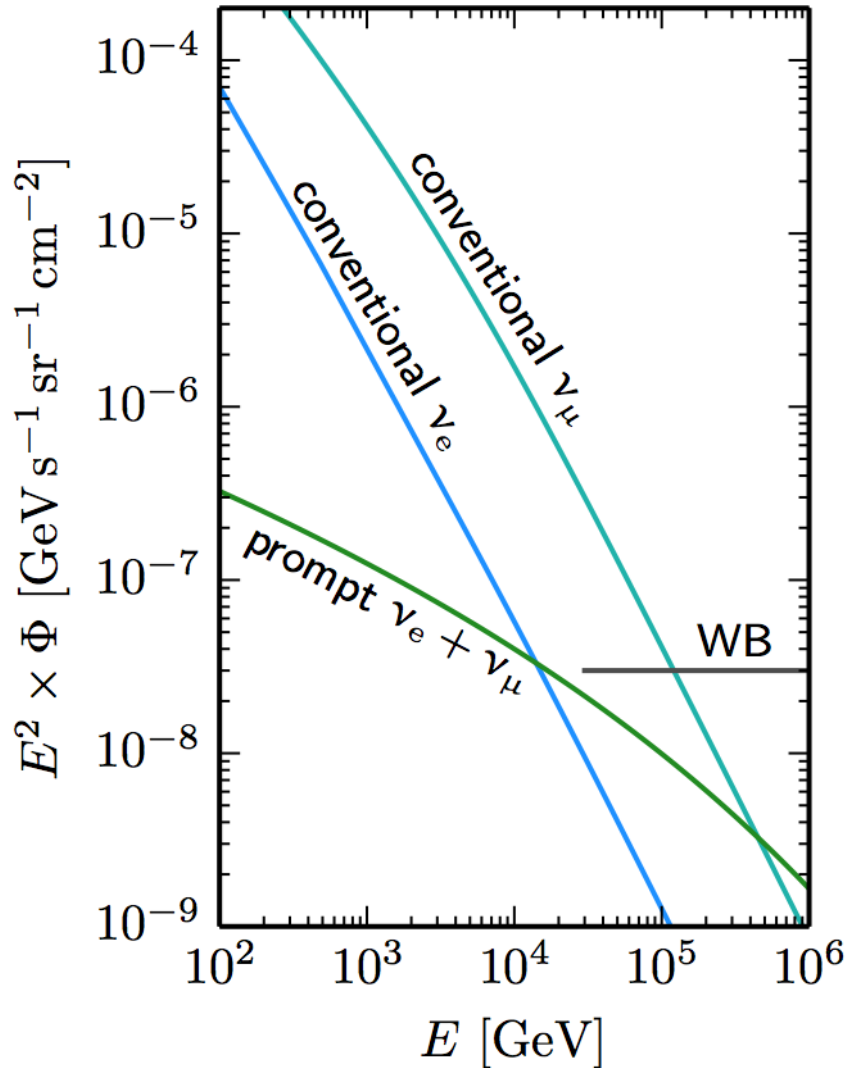
Background in Search for Cosmic Neutrinos



Atmospheric Neutrinos - Production



Atmospheric Neutrinos - Spectrum



Kaons and pions (conventional)

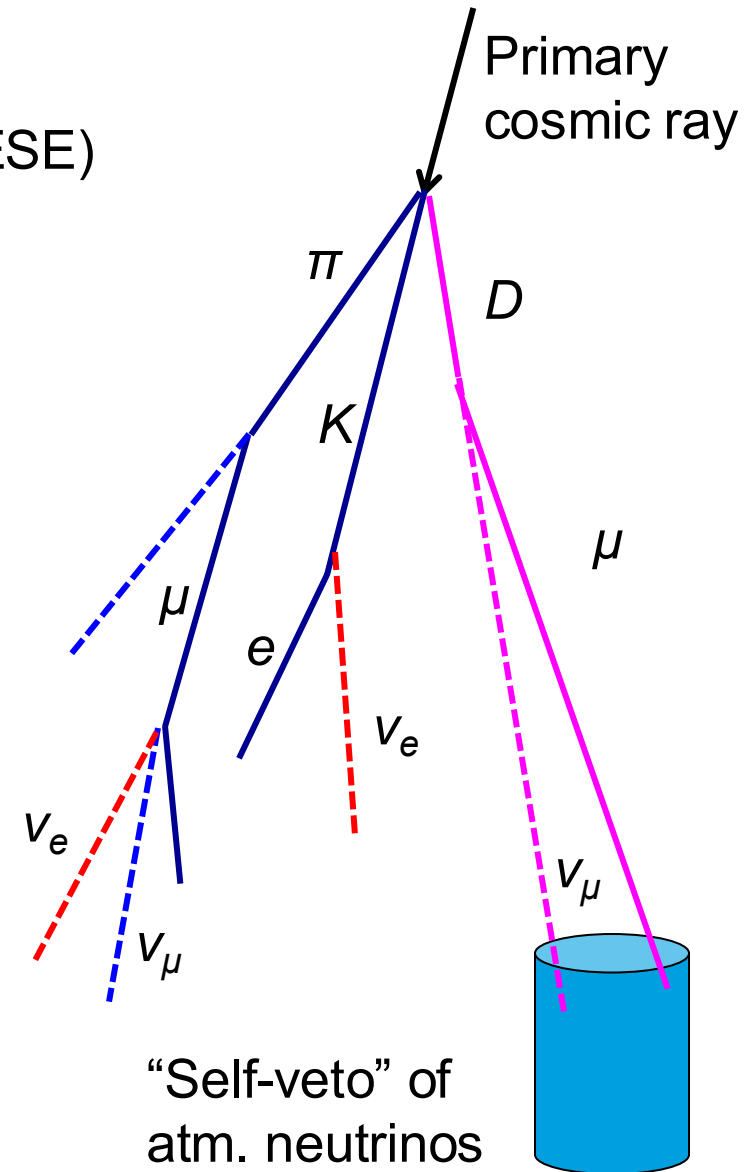
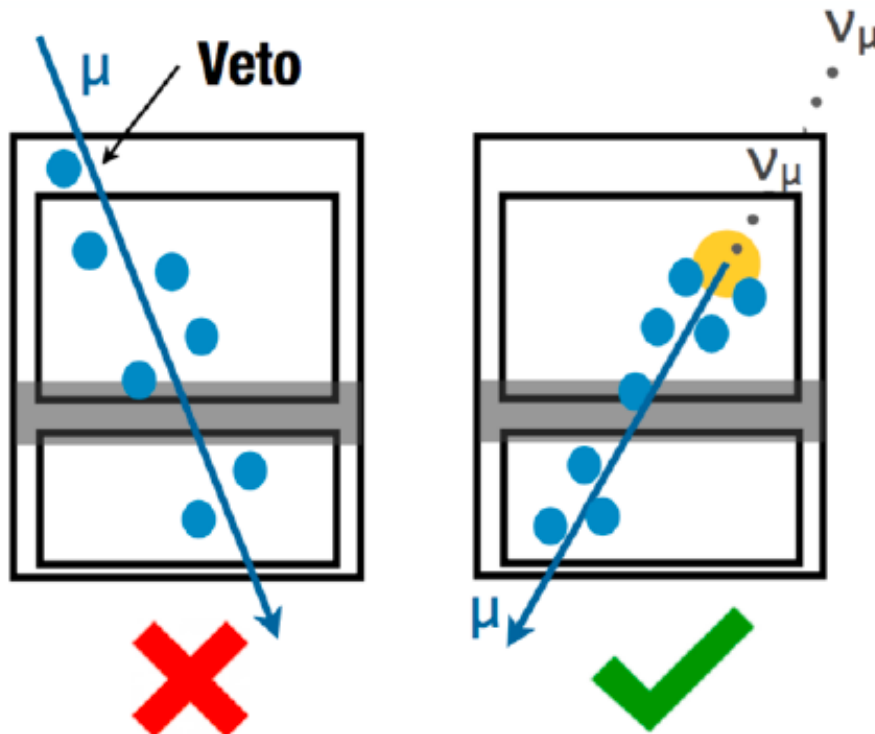
- predominantly decay into muon neutrinos rather than electron neutrinos
- are more likely to interact in the atmosphere rather than to decay into neutrinos if the atmosphere is denser

Charm (prompt)

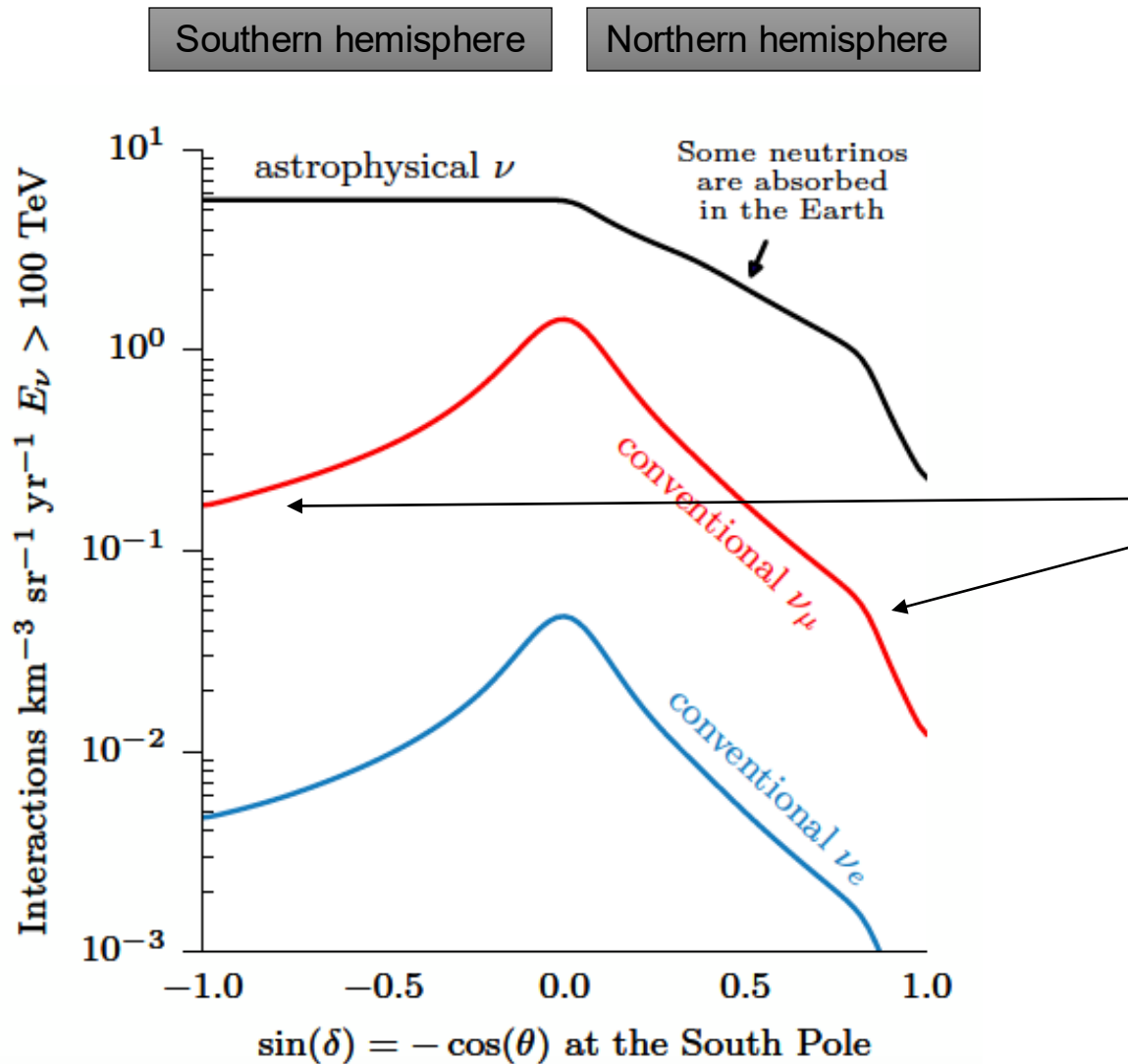
- Decay before interacting

Vetoing Atmospheric Muons and Neutrinos

Selection of high-energy starting events (HESE)



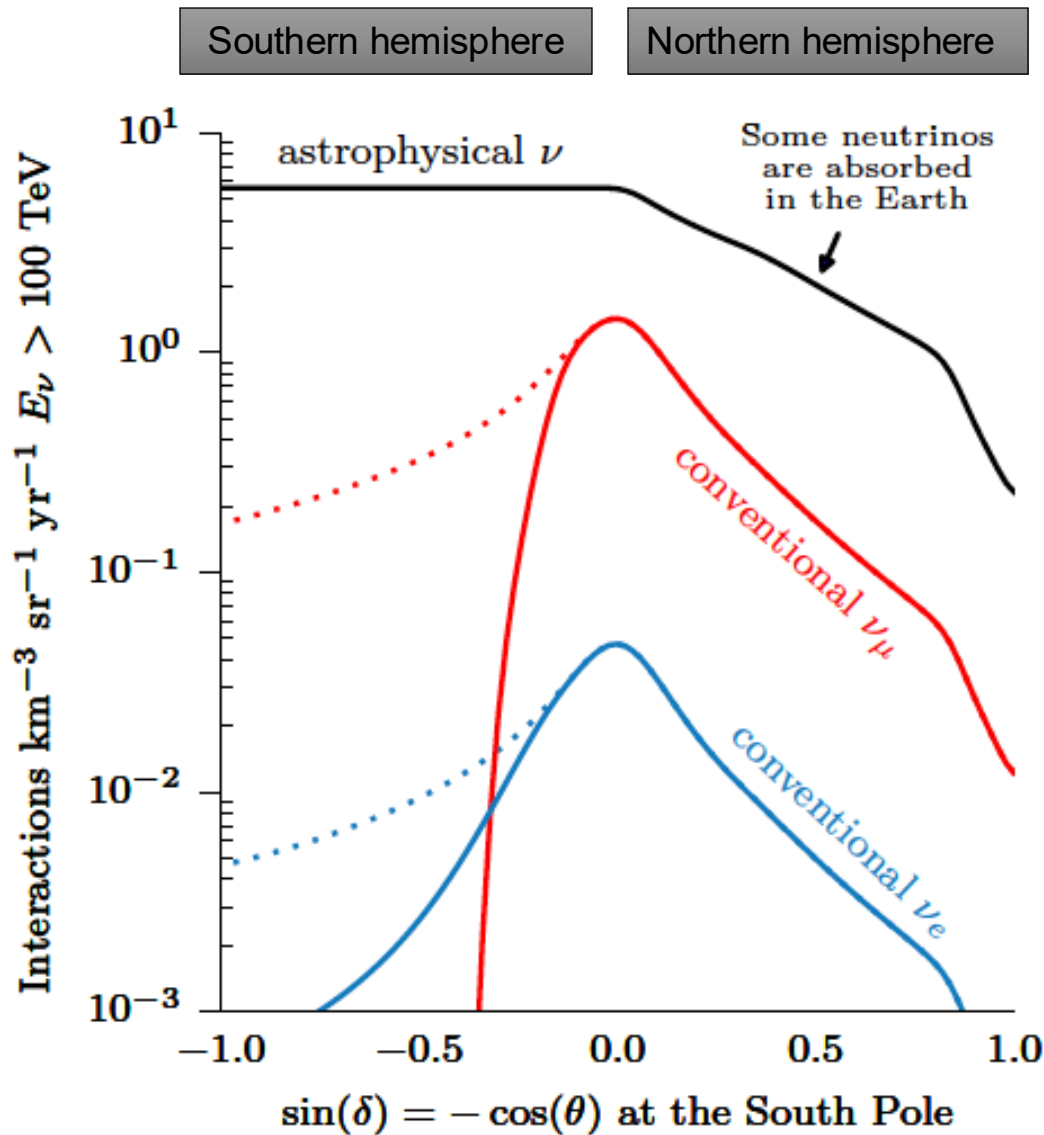
Zenith Distribution



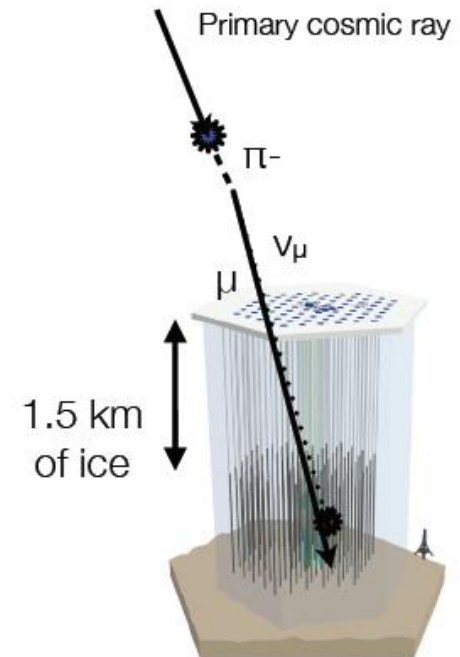
from vertical directions
denser parts of the
atmosphere are reached
more quickly, leading to a
suppression of the
neutrino flux

Schönert, Gaisser, Resconi,
Schulz, Phys. Rev. D,
79:043009 (2009)
Gaisser, Jero, Karle, van Santen,
Phys. Rev. D, 90:023009 (2014)

Zenith Distribution

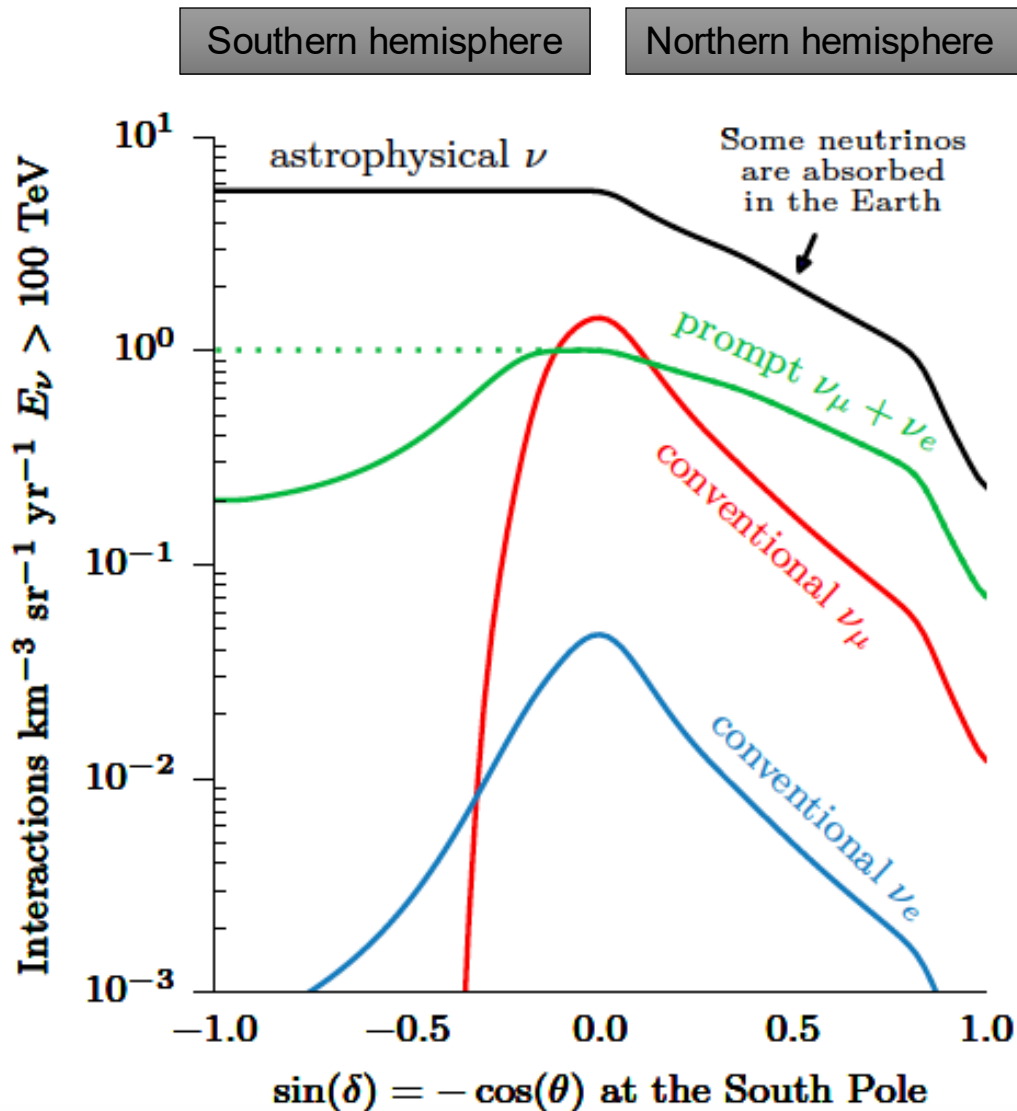


active muon veto
removes down-going
atmospheric neutrinos

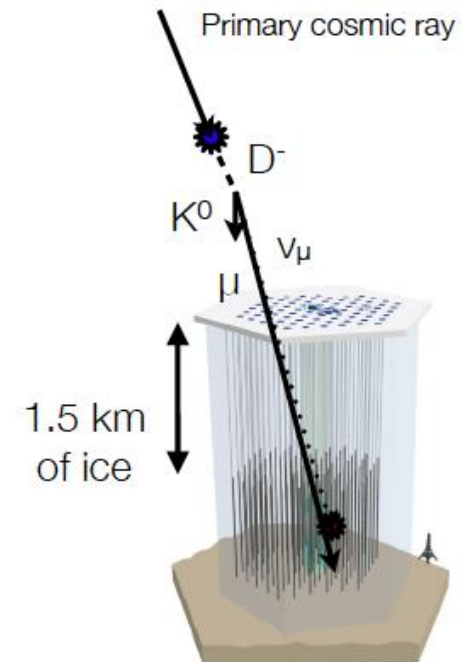


Schönert, Gaisser, Resconi,
Schulz, Phys. Rev. D,
79:043009 (2009)
Gaisser, Jero, Karle, van Santen,
Phys. Rev. D, 90:023009 (2014)

Zenith Distribution

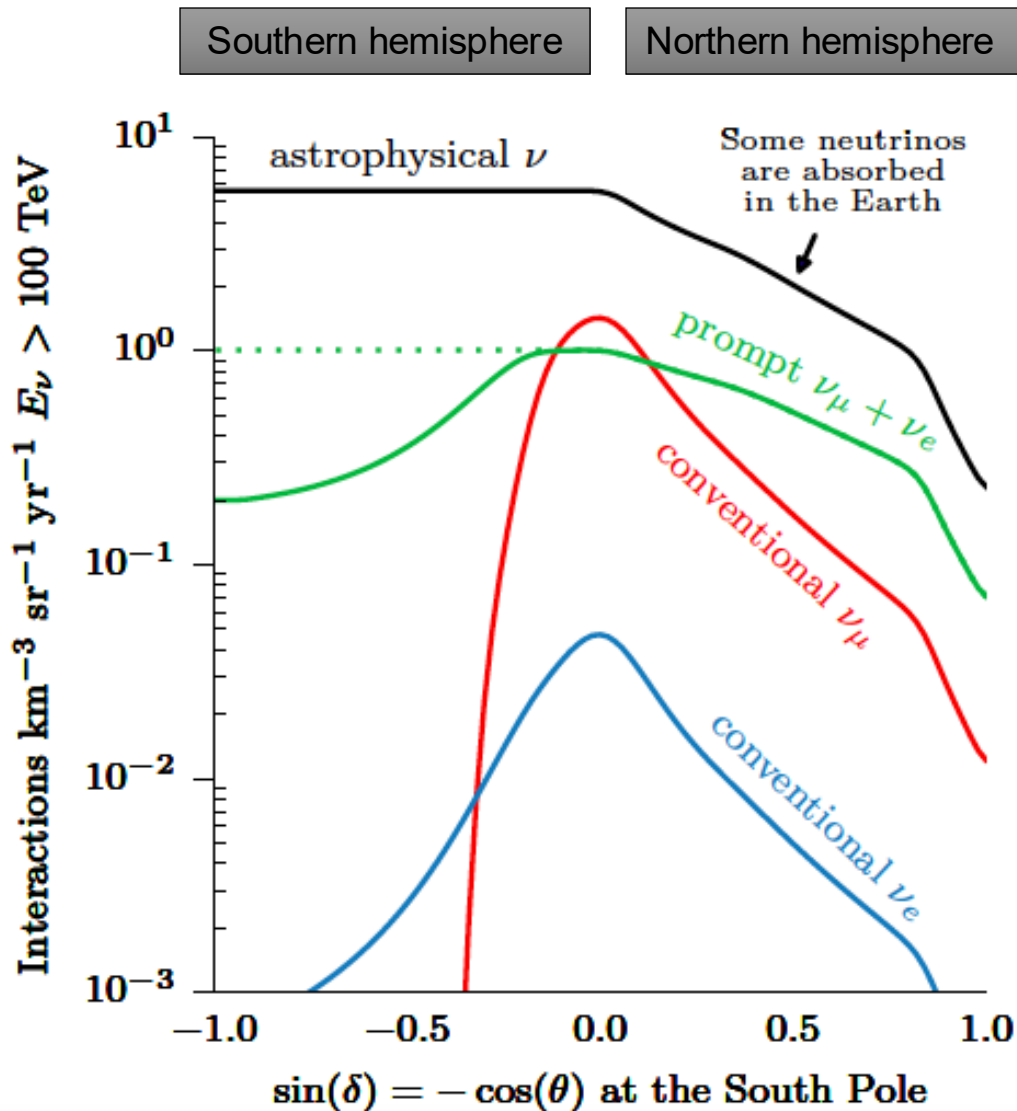


Prompt atmospheric neutrinos are vetoed too



Schönert, Gaisser, Resconi,
Schulz, Phys. Rev. D,
79:043009 (2009)
Gaisser, Jero, Karle, van Santen,
Phys. Rev. D, 90:023009 (2014)

Zenith Distribution



The zenith distributions of high-energy astrophysical and atmospheric neutrinos are fundamentally different.

Schönert, Gaisser, Resconi, Schulz, Phys. Rev. D, 79:043009 (2009)
 Gaisser, Jero, Karle, van Santen, Phys. Rev. D, 90:023009 (2014)



Take home message

Water or ice Cherenkov detectors can reach km^3 size.

**Atmospheric muons and neutrinos are a background for the search of cosmic neutrinos
→ using outer detector layers as veto can reduce this background**



<https://forms.gle/dr1h4NkN1zip3z3XA>

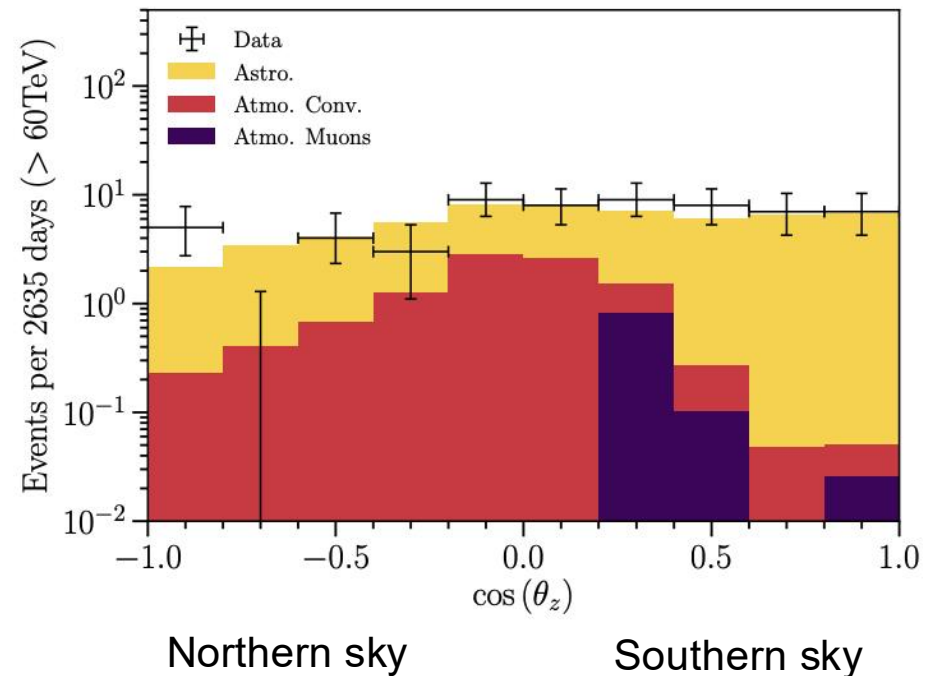
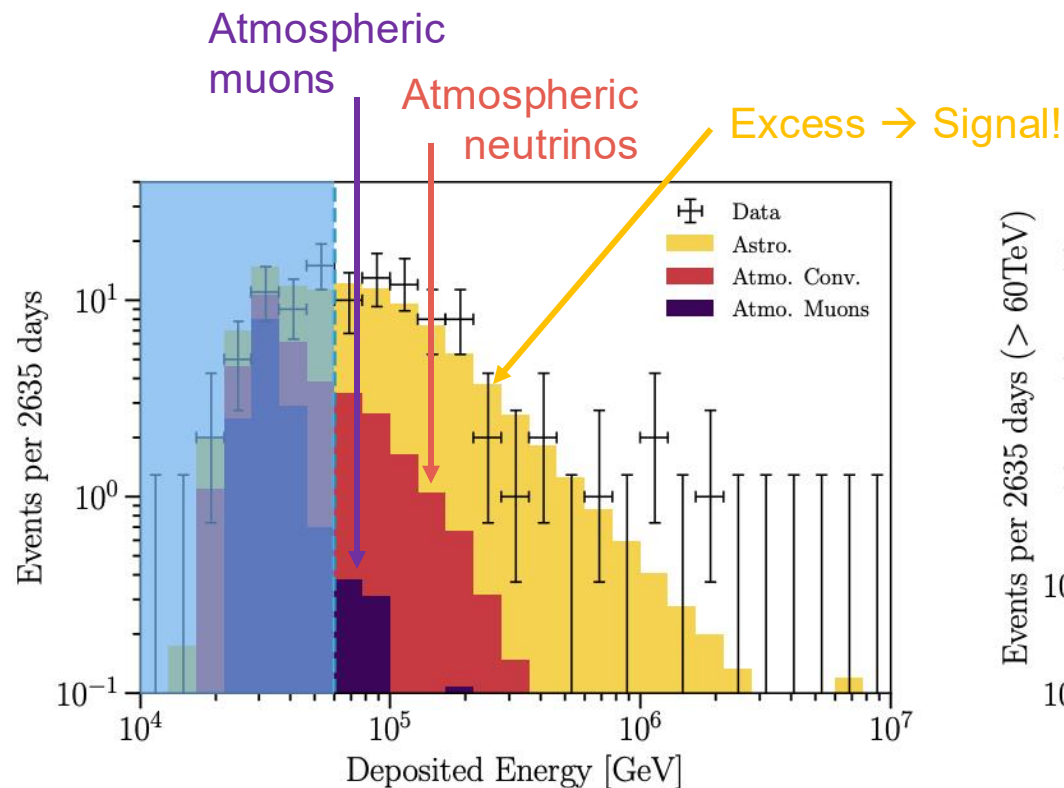
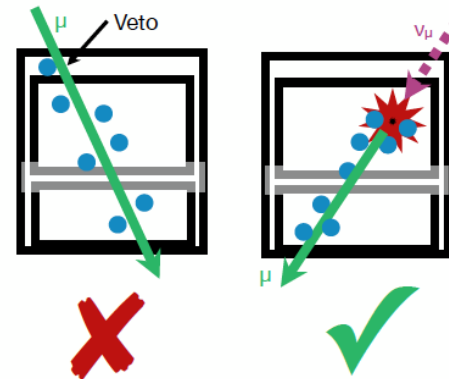
Results from water / ice Cherenkov Telescopes

- Astronomy
 - Diffuse Flux
 - Milky Way
 - Source Candidates
- Neutrino as a particle
 - Neutrino Oscillations
 - Neutrino Cross Section
- New physics
 - Sterile Neutrinos
 - Dark Matter

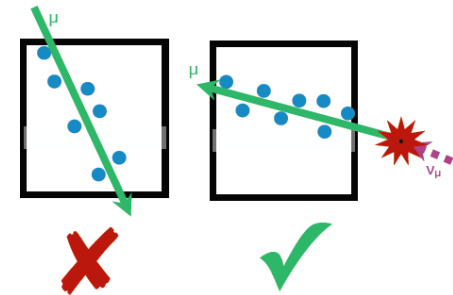
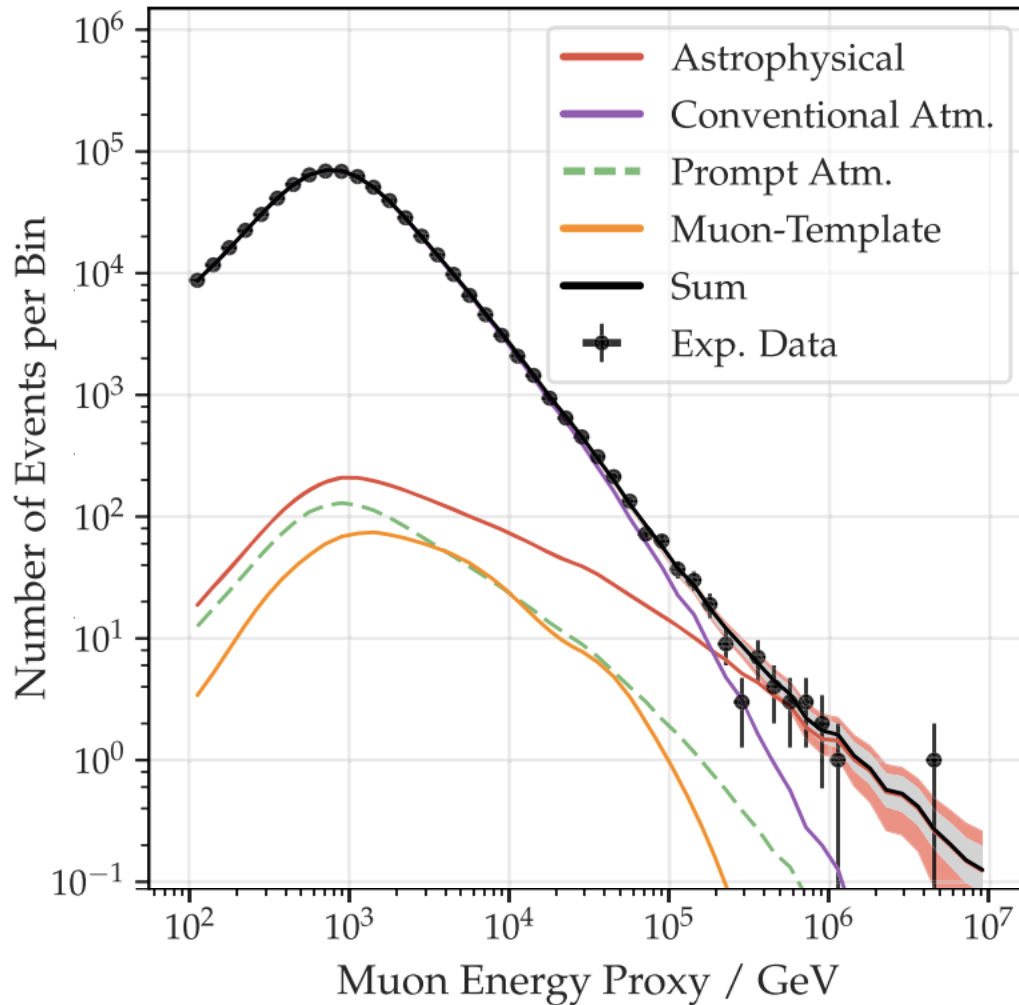
Results from water / ice Cherenkov Telescopes

- Astronomy
 - **Diffuse Flux**
 - Milky Way
 - Source Candidates
- Neutrino as a particle
 - Neutrino Oscillations
 - Neutrino Cross Section
- New physics
 - Sterile Neutrinos
 - Dark Matter

Discovery of High-Energy Astrophysical Neutrinos

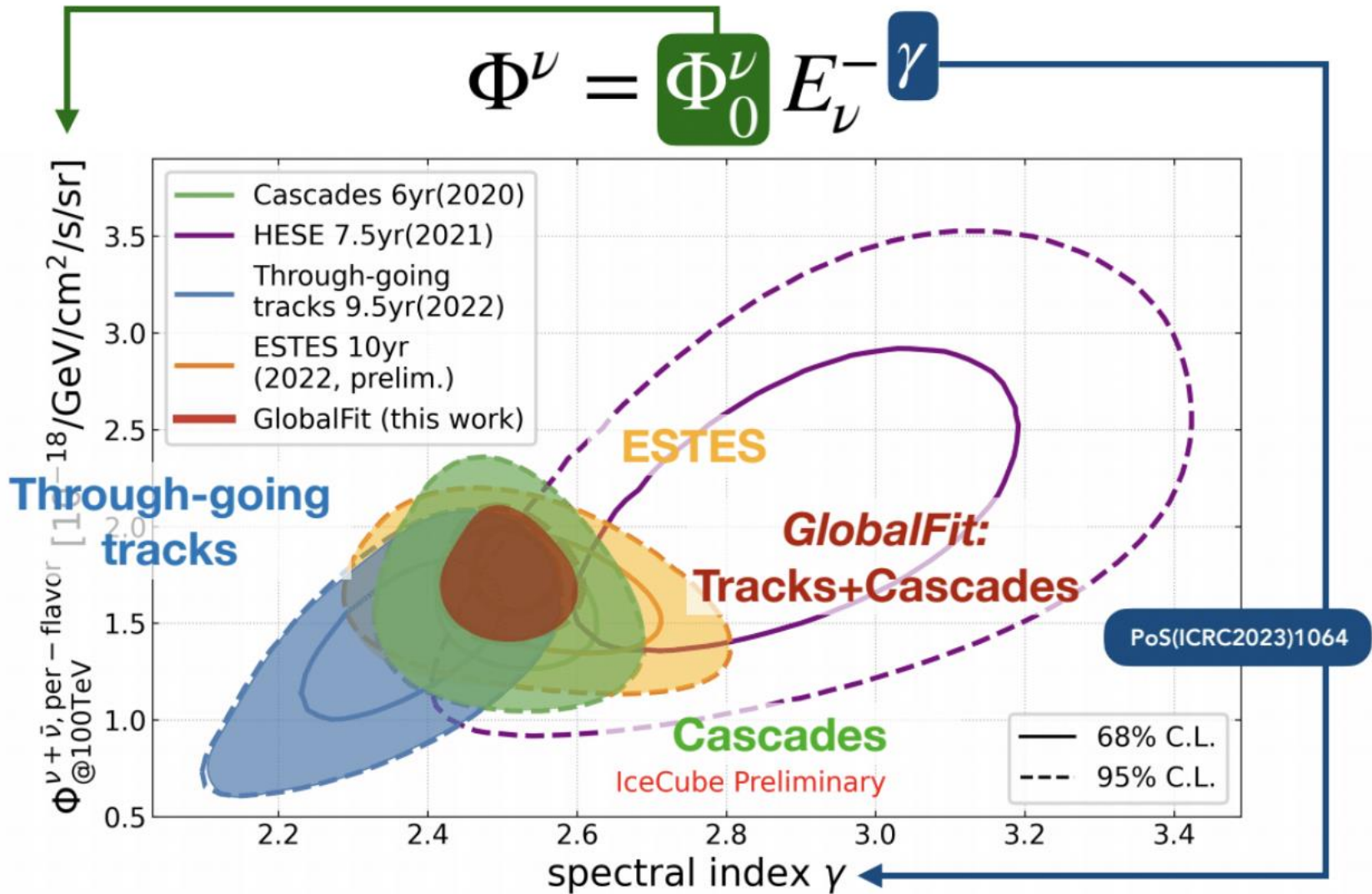


Diffuse flux seen in different channels

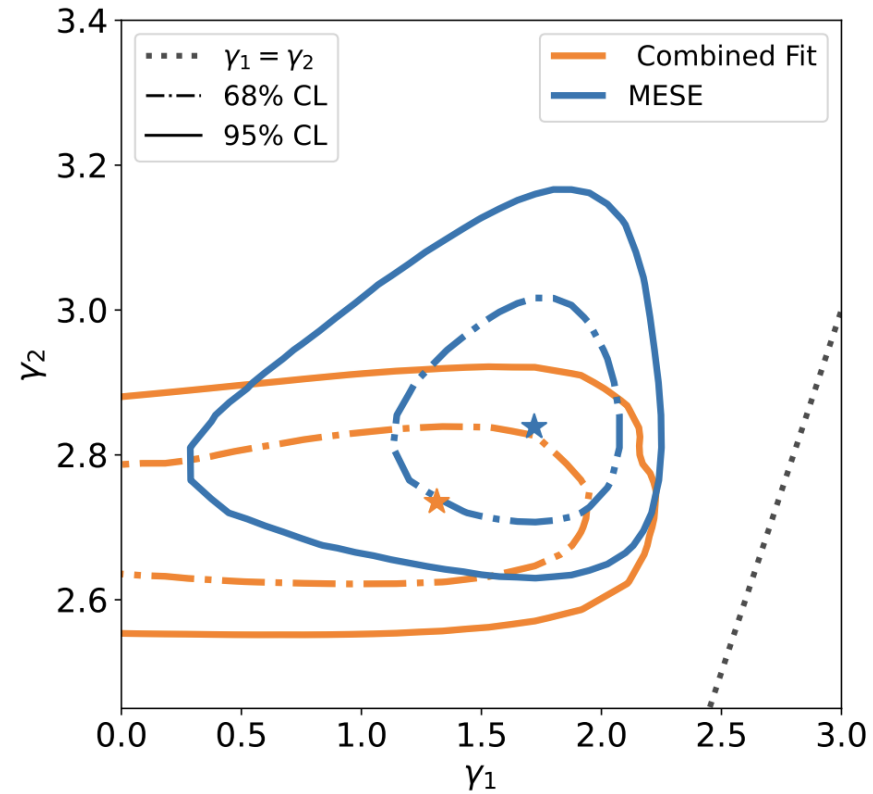
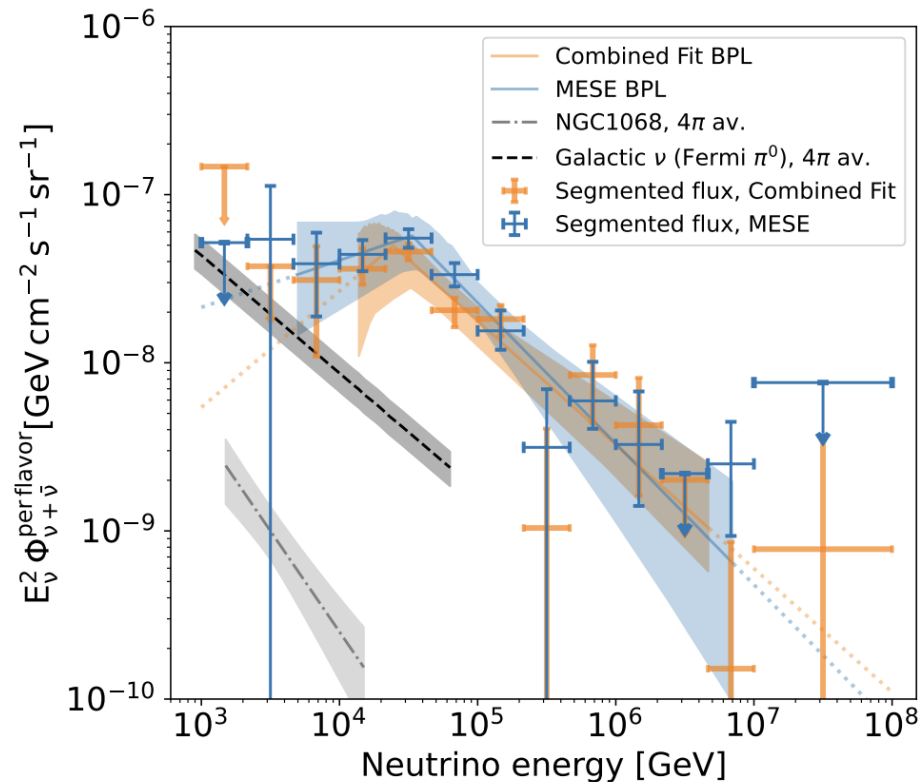


- Selected horizontal and up-going muon tracks
- Sensitive to astrophysical neutrinos above ~ 100 TeV

Spectrum of diffuse flux



Spectrum of diffuse flux: not a power law?



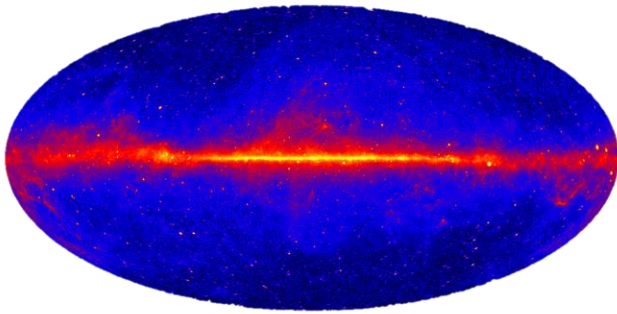
**What is the origin of the
diffuse flux?**

Results from water / ice Cherenkov Telescopes

- Astronomy
 - Diffuse Flux
 - **Milky Way**
 - Source Candidates
- Neutrino as a particle
 - Neutrino Oscillations
 - Neutrino Cross Section
- New physics
 - Sterile Neutrinos
 - Dark Matter

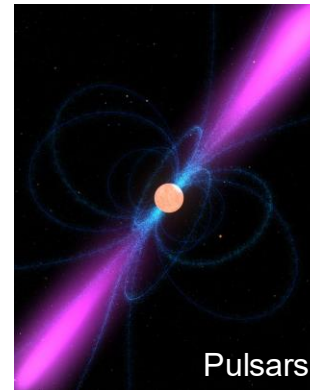
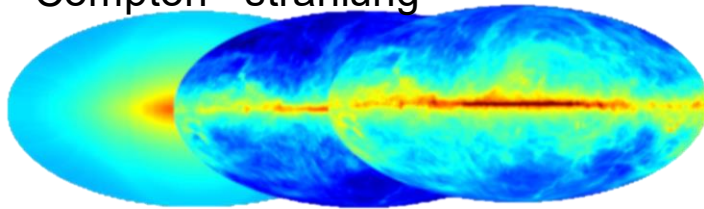
Galactic Contribution

GeV gamma-ray sky by Fermi-LAT

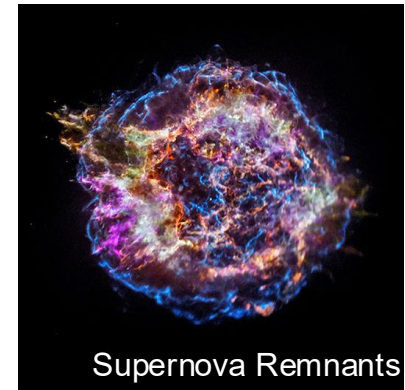


Cosmic rays propagate through the Galaxy and interact with photons and gas

Inverse Compton Bremsstrahlung π^0 decay



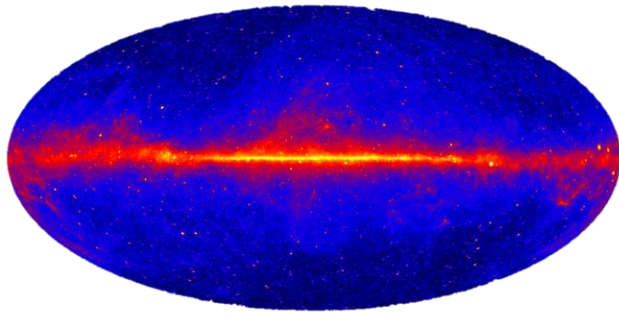
Pulsars



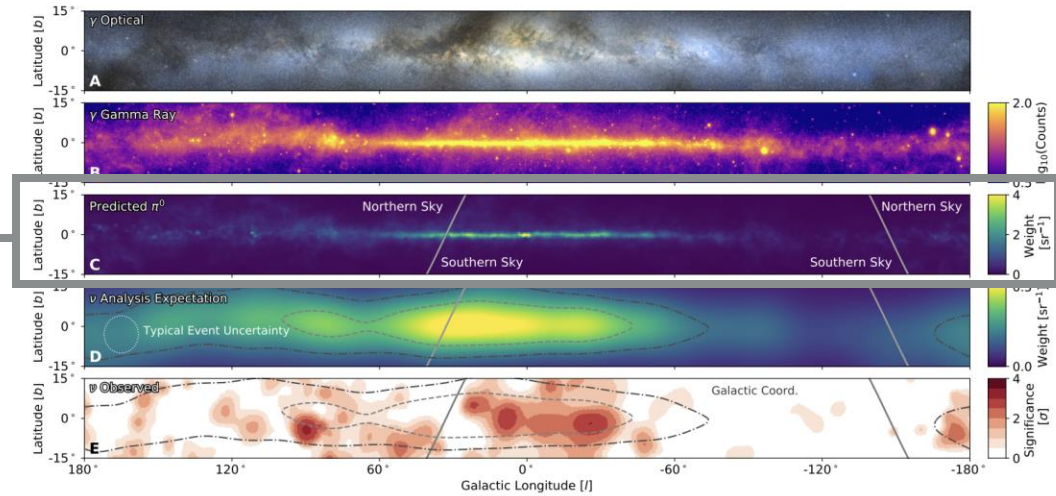
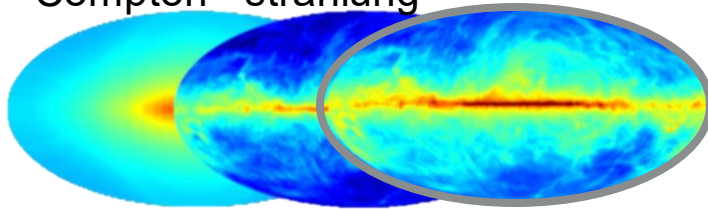
Supernova Remnants

Galactic Contribution

GeV gamma-ray sky by Fermi-LAT

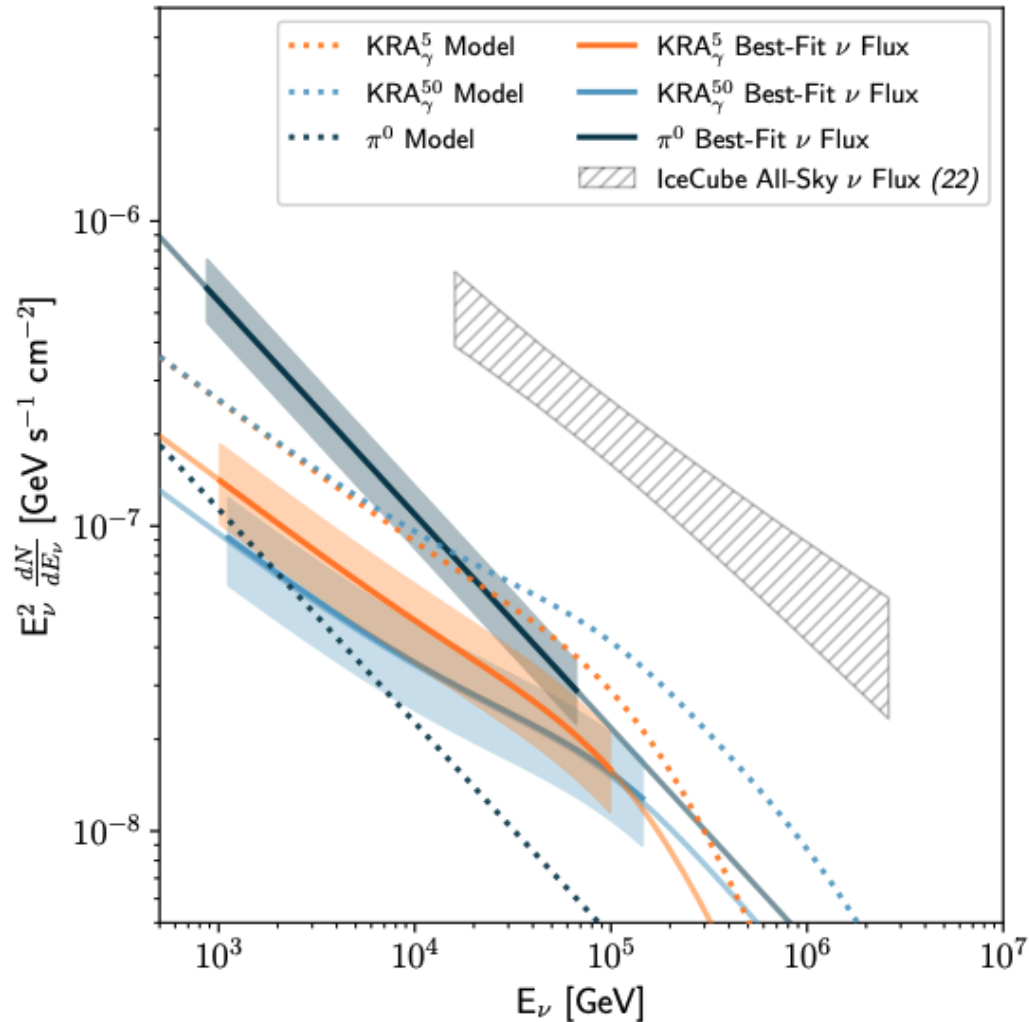


Inverse Compton Bremsstrahlung π^0 decay



First detection of galactic plane neutrino flux thanks to gamma-ray template fit

Galactic Neutrinos



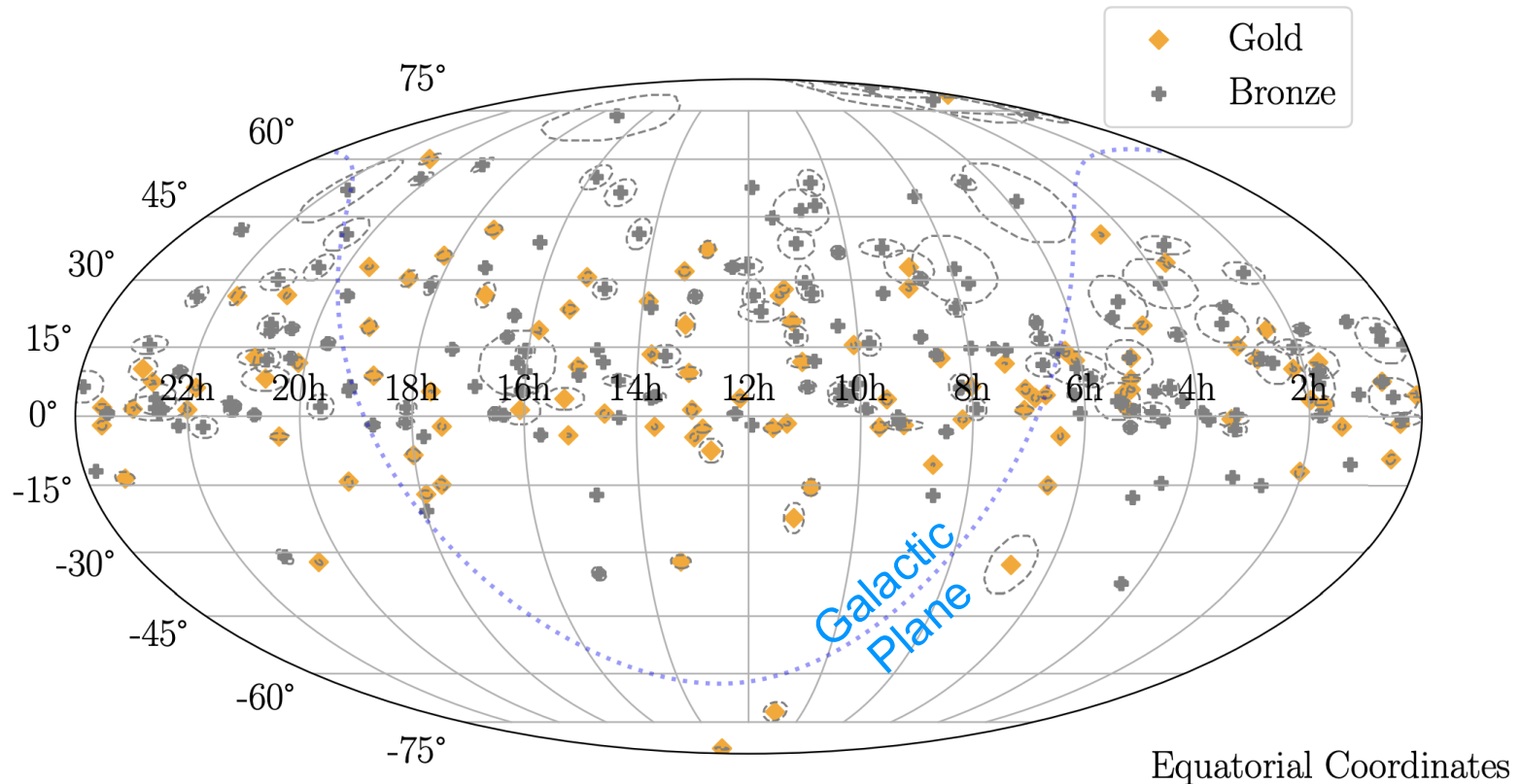
Galactic neutrinos
make up ~10% of
diffuse neutrinos

Results from water / ice Cherenkov Telescopes

- Astronomy
 - Diffuse Flux
 - Milky Way
 - **Source Candidates**
- Neutrino as a particle
 - Neutrino Oscillations
 - Neutrino Cross Section
- New physics
 - Sterile Neutrinos
 - Dark Matter

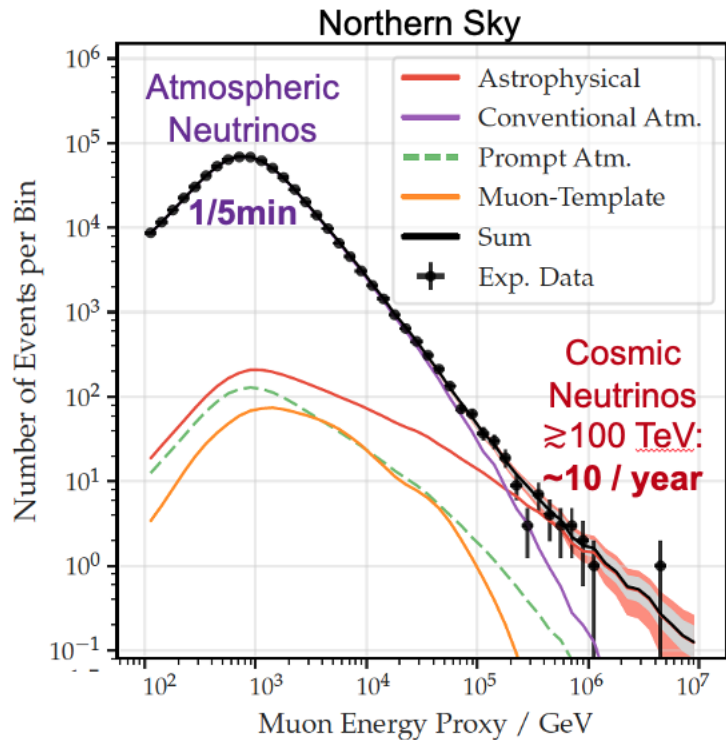
Neutrino Sky Map

IceCube neutrinos with high ($>30\%$) probability to be of cosmic origin



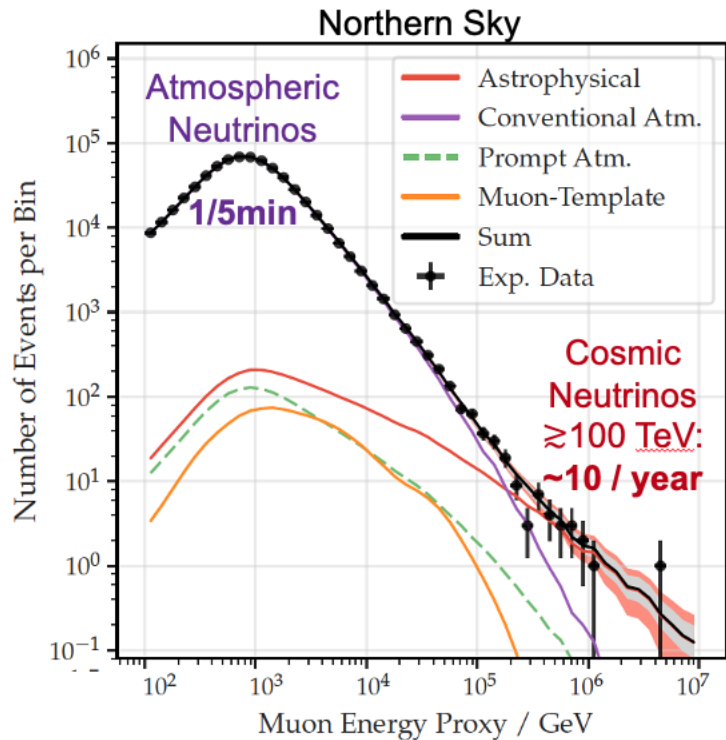
Neutrinos alone do not reveal the sources (yet)

Search for Extragalactic Sources: Strategies



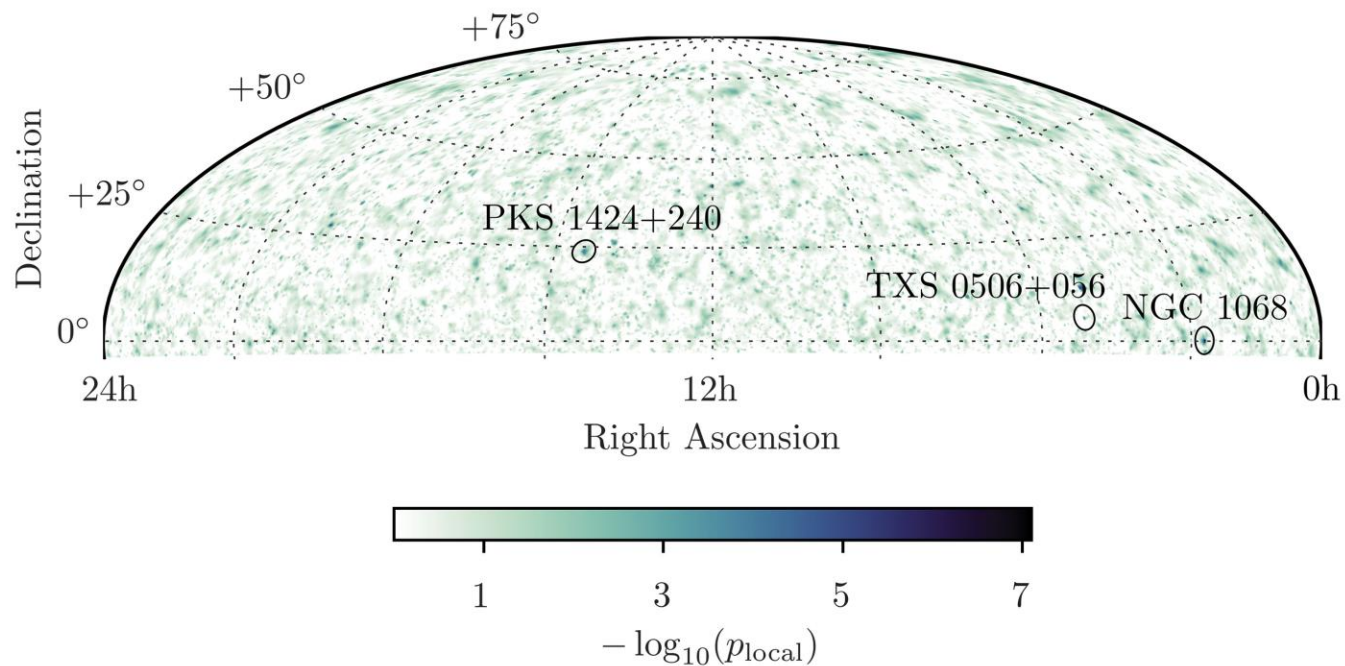
1. Look for hotspots in the neutrino sky $\rightarrow$ identify source candidates
2. Start from EM source catalog $\rightarrow$ look for neutrinos from source population
3. Focus on high-energy neutrinos with high signal probability $\rightarrow$ look for EM counterparts

Search for Extragalactic Sources: Strategies



1. Look for hotspots in the neutrino sky
→ identify source candidates
2. Start from EM source catalog → look for neutrinos from source population
3. Focus on high-energy neutrinos with high signal probability → look for EM counterparts

Search for hotspots in the neutrino sky



Maximum Likelihood Technique

$$\mathcal{L}(n_S, \gamma) = \prod_{\substack{i \\ \text{Neutrino event } i}} \left(\overbrace{\frac{n_S}{N} \mathcal{S}(|\mathbf{x}_S - \mathbf{x}_i|, E_i; \gamma)}^{\text{Signal probability}} + \left(1 - \frac{n_S}{N}\right) \underbrace{\mathcal{B}(\sin \delta_i, E_i)}_{\text{background probability}} \right)$$

Test Statistic:

$$\mathcal{TS} = 2 \log (\mathcal{L}(\hat{n}_S, \hat{\gamma}) / \mathcal{L}(n_S = 0))$$

Maximum Likelihood Technique

$$\mathcal{S} = \frac{1}{2\pi\sigma_i^2} e^{-\frac{|\mathbf{x}_S - \mathbf{x}_i|^2}{2\sigma_i^2}} \times \mathcal{E}_S(E_i, \sin \delta_i; \gamma)$$

Signal probability

$$\mathcal{L}(n_S, \gamma) = \prod_i \left(\frac{n_S}{N} \overbrace{\mathcal{S}(|\mathbf{x}_S - \mathbf{x}_i|, E_i; \gamma)}^{\text{Signal probability}} + \left(1 - \frac{n_S}{N}\right) \mathcal{B}(\sin \delta_i, E_i) \right)$$

$\mathbf{x}_S$: source position

$\mathbf{x}_i$: neutrino direction

δ_i : neutrino declination

σ_i : neutrino declination

E_i : neutrino energy

γ : spectral index

Maximum Likelihood Technique

$$\mathcal{S} = \frac{1}{2\pi\sigma_i^2} e^{-\frac{|\mathbf{x}_S - \mathbf{x}_i|^2}{2\sigma_i^2}} \times \mathcal{E}_S(E_i, \sin \delta_i, \gamma)$$

Signal probability

Spectral index
(free parameter
in fit)

$$\mathcal{L}(n_S, \gamma) = \prod_i \left(\frac{n_S}{N} \mathcal{S}(|\mathbf{x}_S - \mathbf{x}_i|, E_i; \gamma) + \left(1 - \frac{n_S}{N}\right) \mathcal{B}(\sin \delta_i, E_i) \right)$$

number of events
from source (**free
parameter in fit**)

$\mathbf{x}_S$: source position

$\mathbf{x}_i$: neutrino direction

δ_i : neutrino declination

σ_i : neutrino declination

E_i : neutrino energy

γ : spectral index

Maximum Likelihood Technique

$$\mathcal{S} = \frac{1}{2\pi\sigma_i^2} e^{-\frac{|\mathbf{x}_S - \mathbf{x}_i|^2}{2\sigma_i^2}} \times \mathcal{E}_S(E_i, \sin \delta_i, \gamma)$$

Signal probability

Spectral index
(free parameter
in fit)

$$\mathcal{L}(n_S, \gamma) = \prod_i \left(\frac{n_S}{N} \mathcal{S}(|\mathbf{x}_S - \mathbf{x}_i|, E_i; \gamma) + \left(1 - \frac{n_S}{N}\right) \mathcal{B}(\sin \delta_i, E_i) \right)$$

number of events
from source (**free
parameter in fit**)

background
probability

$$\mathcal{B} = \frac{\mathcal{P}_B(\sin \delta_i)}{2\pi} \times \mathcal{E}_B(E_i, \sin \delta_i)$$

$\mathbf{x}_S$: source position

$\mathbf{x}_i$: neutrino direction

δ_i : neutrino declination

σ_i : neutrino declination

E_i : neutrino energy

γ : spectral index

Maximum Likelihood Technique

$$\mathcal{L}(n_S, \gamma) = \prod_i \left(\frac{n_S}{N} \underbrace{\mathcal{S}(|\mathbf{x}_S - \mathbf{x}_i|, E_i; \gamma)}_{\text{Signal probability}} + \left(1 - \frac{n_S}{N}\right) \underbrace{\mathcal{B}(\sin \delta_i, E_i)}_{\text{background probability}} \right)$$

Spectral index
(free parameter
in fit)

$\mathcal{S} = \frac{1}{2\pi\sigma_i^2} e^{-\frac{|\mathbf{x}_S - \mathbf{x}_i|^2}{2\sigma_i^2}} \times \mathcal{E}_S(E_i, \sin \delta_i, \gamma)$

$\mathcal{B} = \frac{\mathcal{P}_B(\sin \delta_i)}{2\pi} \times \mathcal{E}_B(E_i, \sin \delta_i)$

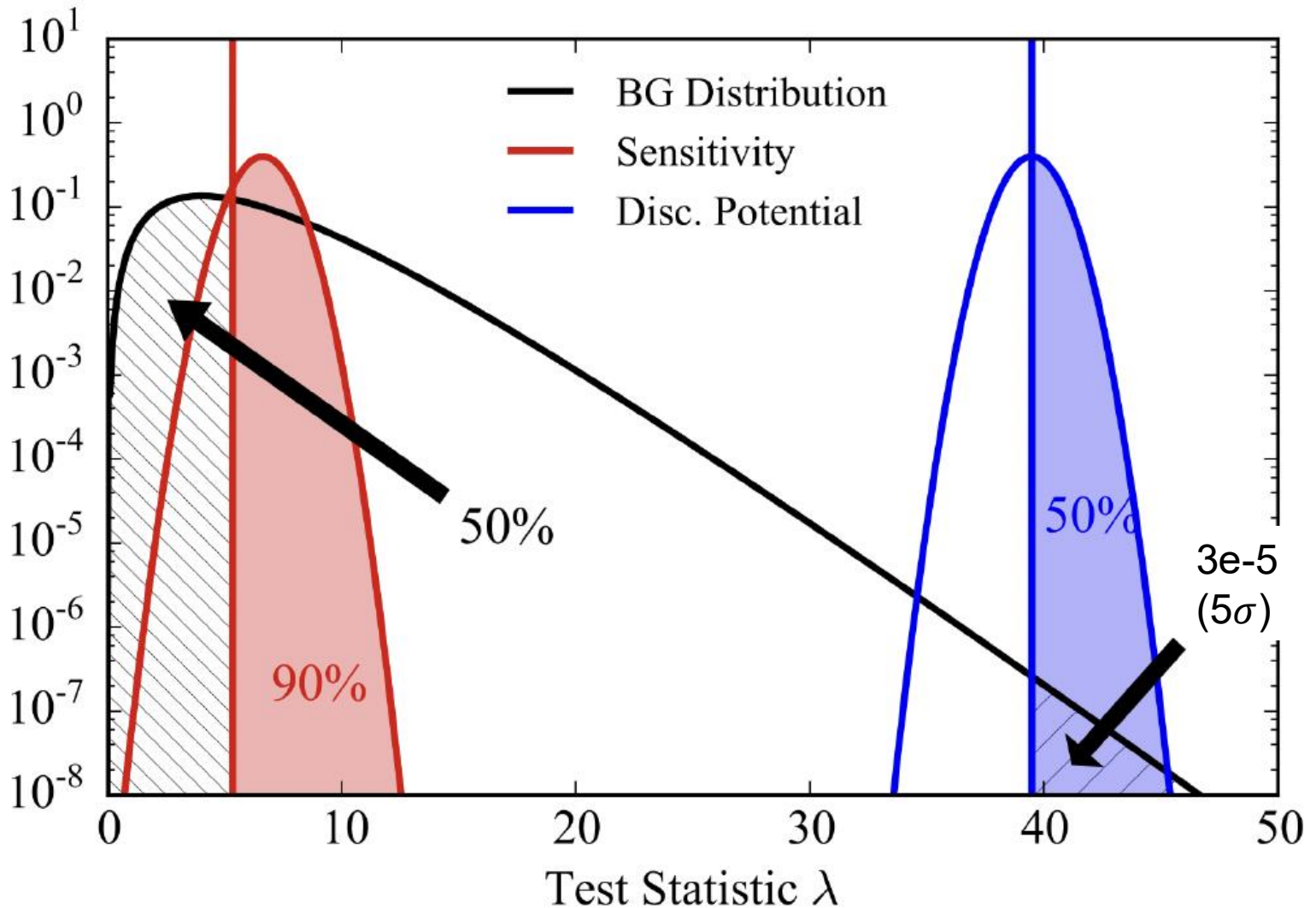
Diagram annotations:

- A blue circle around $\frac{n_S}{N}$ in the likelihood equation has an arrow pointing to a box labeled "number of events from source (**free parameter in fit**)".
- A blue circle around $\frac{n_S}{N}$ in the background term has an arrow pointing to the same box.
- A green bracket under the signal term is labeled "Signal probability".
- A red bracket under the background term is labeled "background probability".
- A blue circle around γ in the signal term has an arrow pointing to a box labeled "Spectral index (free parameter in fit)".

Test Statistic:

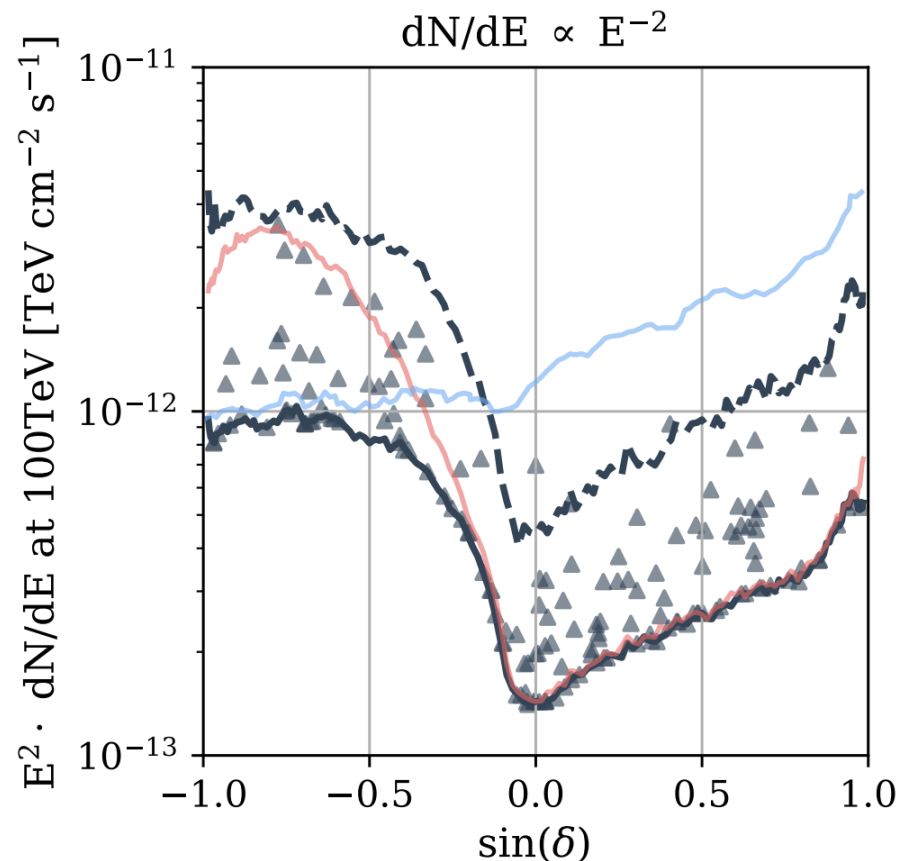
$$\mathcal{TS} = 2 \log (\mathcal{L}(\hat{n}_S, \hat{\gamma}) / \mathcal{L}(n_S = 0))$$

Definition of Sensitivity



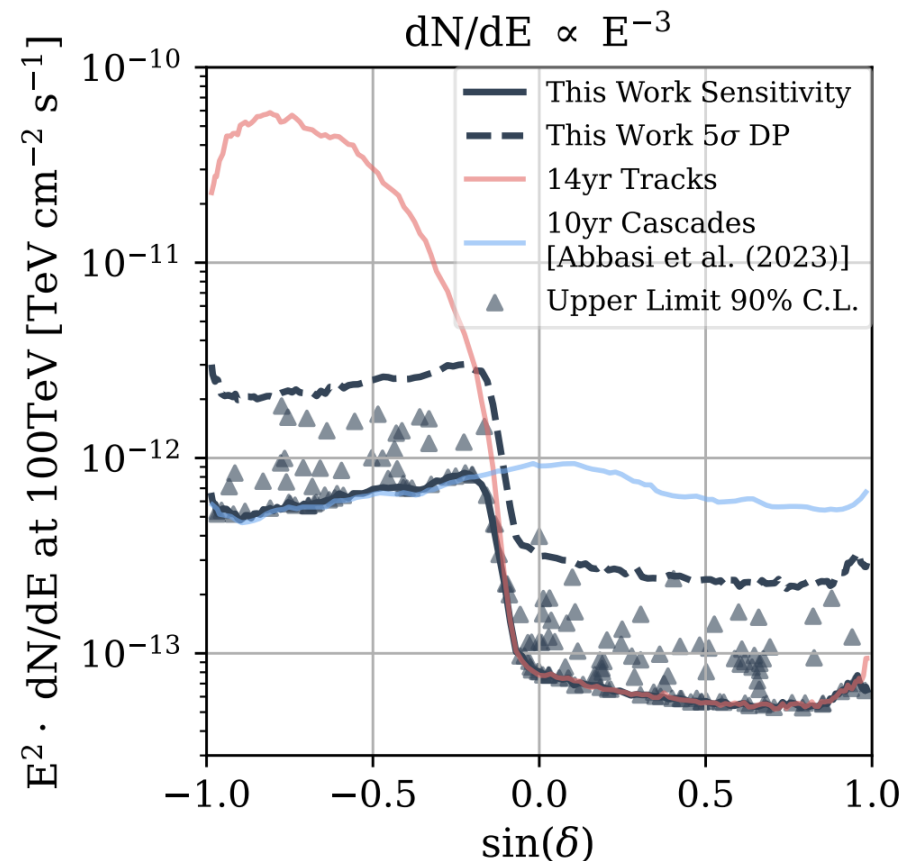
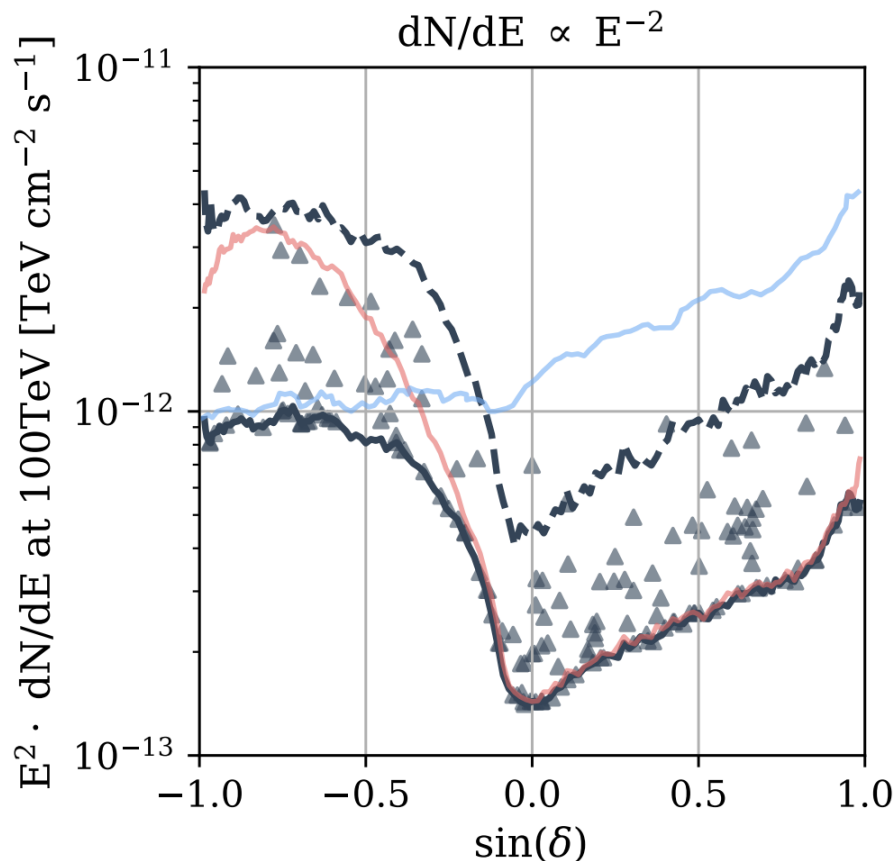
Point Source Sensitivity

Signal trials: determine how much signal needs to be injected to end have 90% of the signal trials yield a TS greater than the median TS derived from background trials

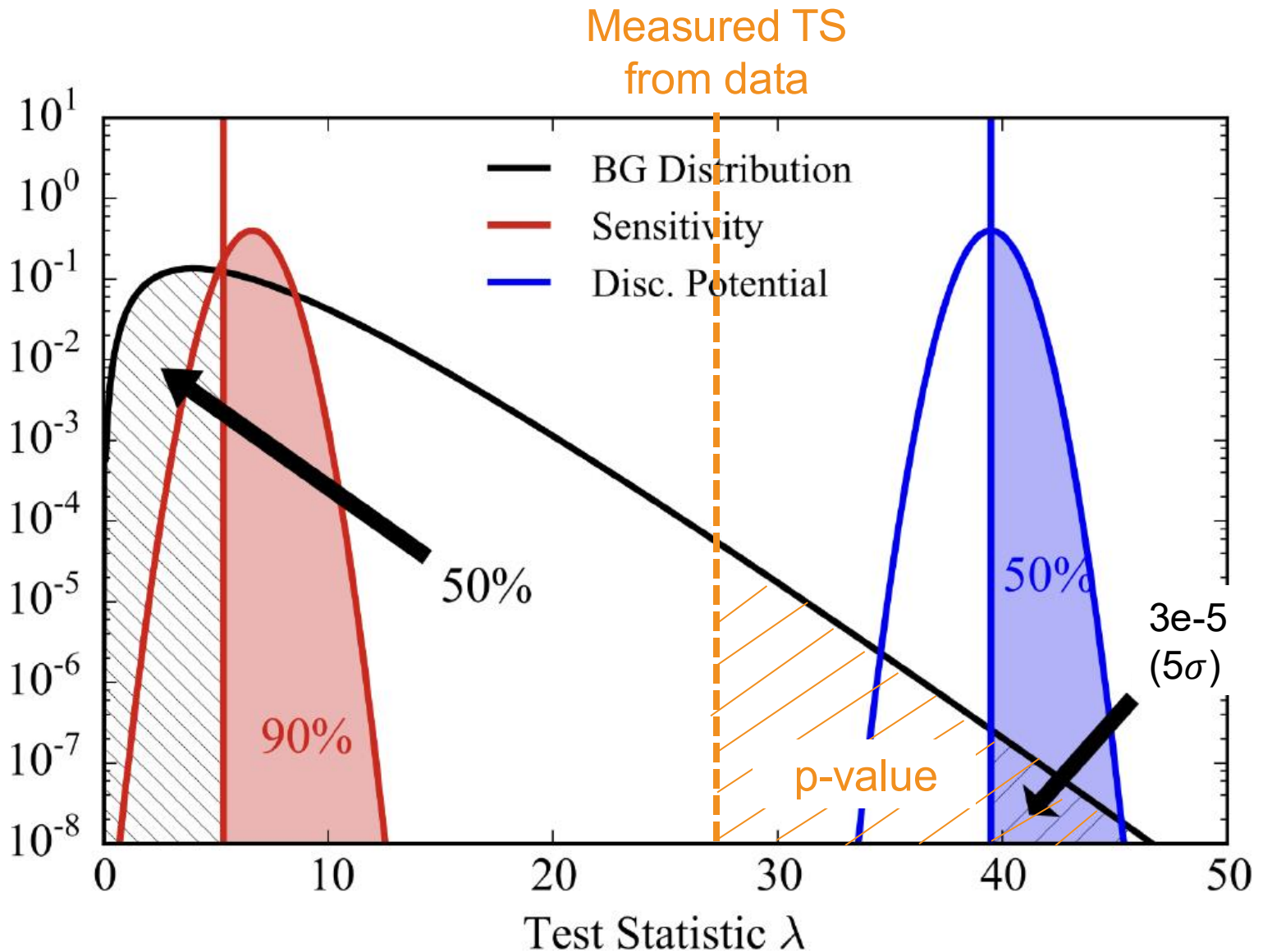


Point Source Sensitivity

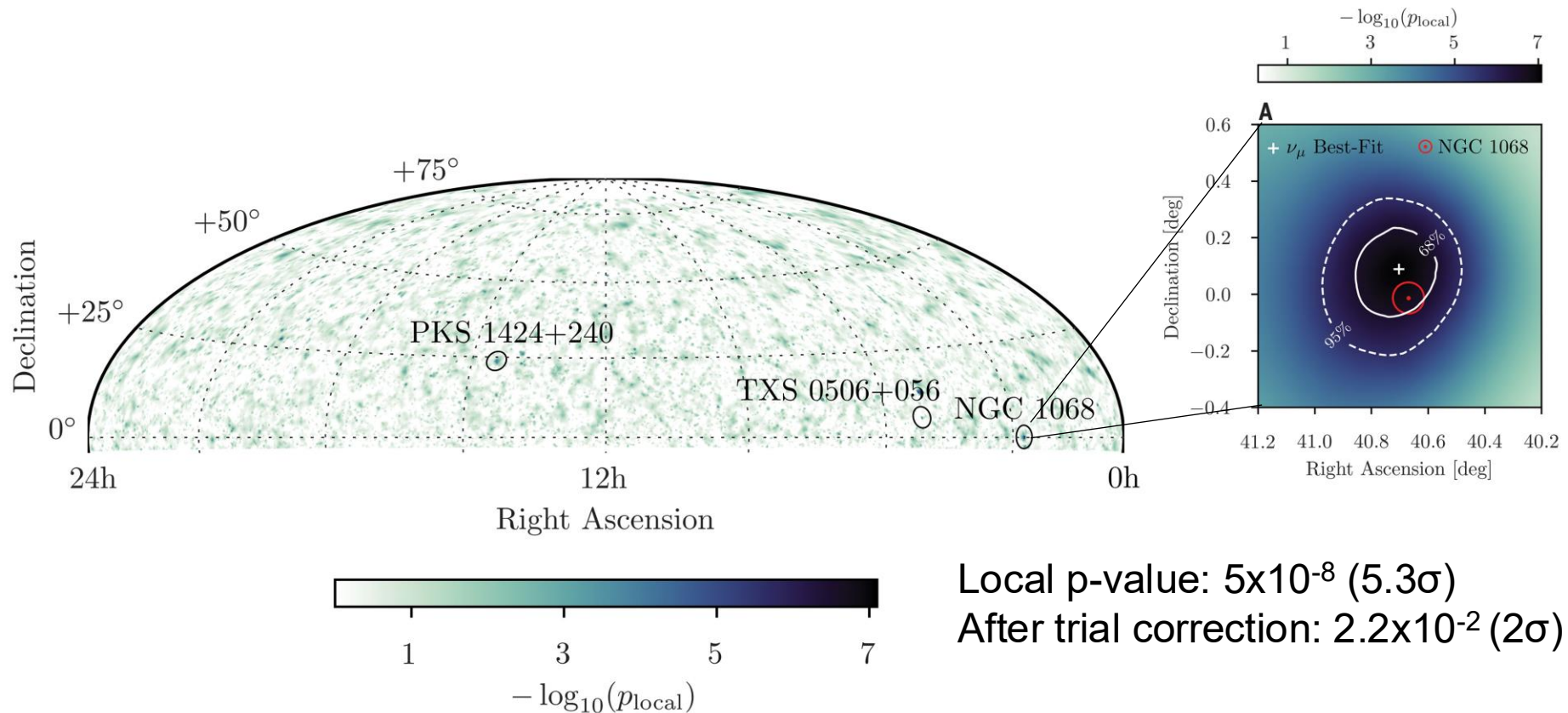
Signal trials: determine how much signal needs to be injected to end have 90% of the signal trials yield a TS greater than the median TS derived from background trials



Definition of p-value

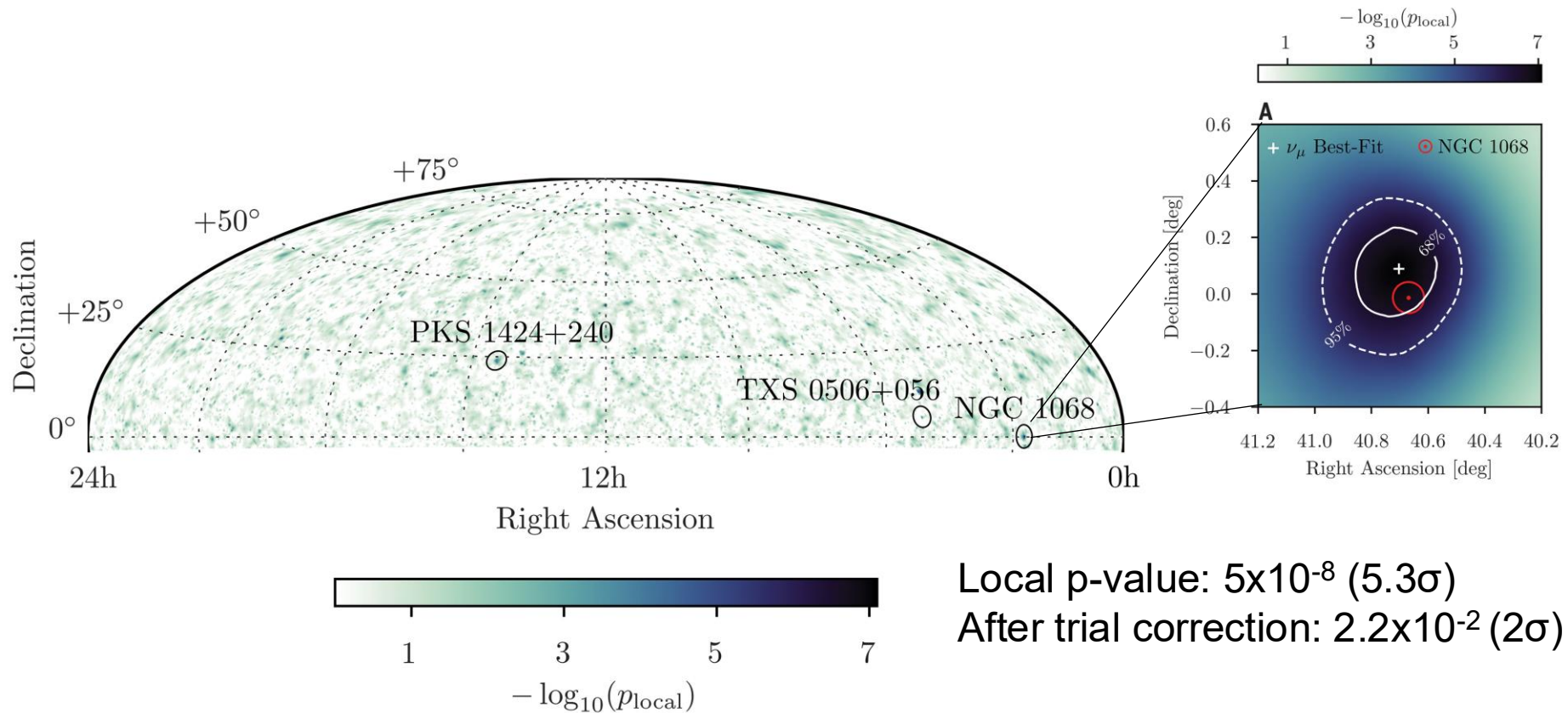


Search for hotspots in the neutrino sky



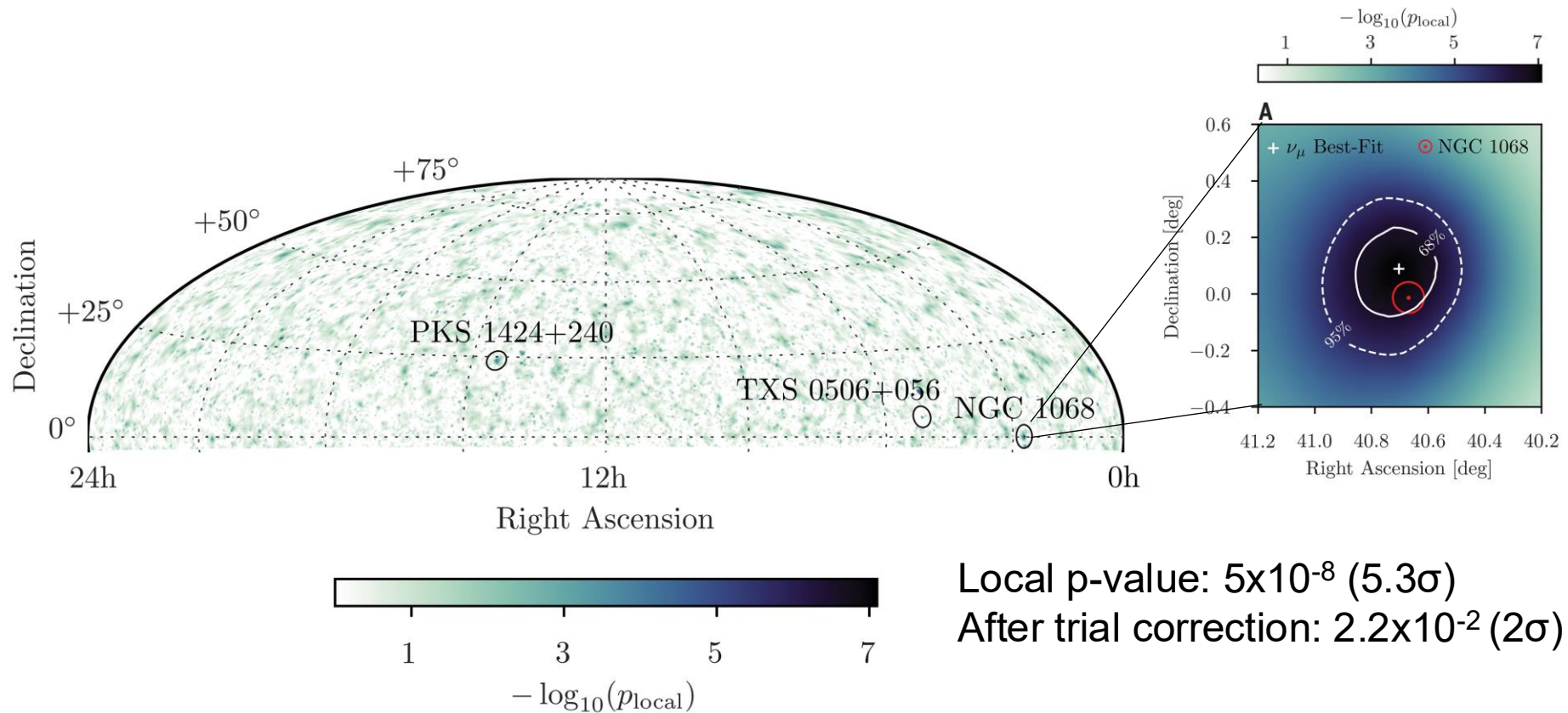
The post-trial probability is estimated by comparing the p-value of the hottest spot in the data with a distribution of hottest spots in a large number of background trials.

Search for hotspots in the neutrino sky



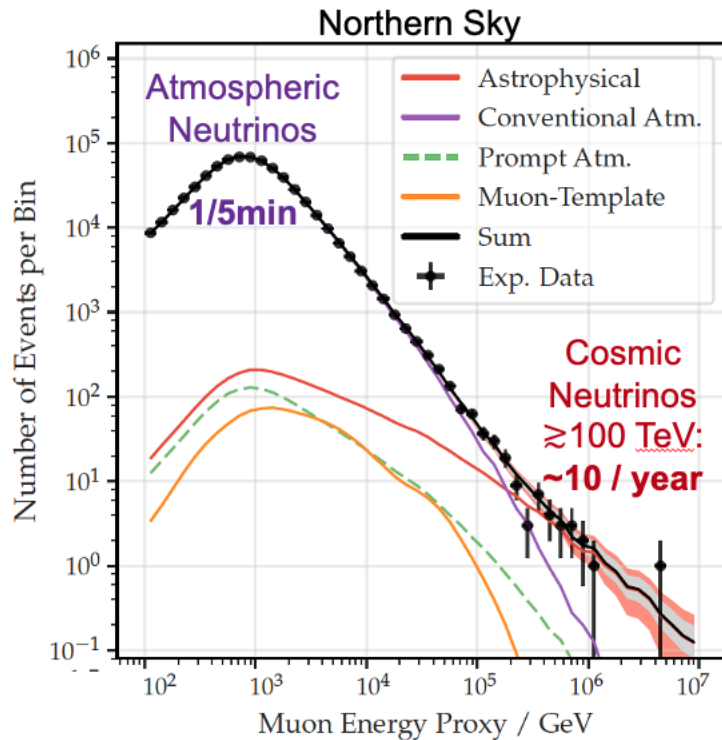
Challenge: Atmospheric background, large trial factor

Search for hotspots in the neutrino sky



Solution: Use predefined source lists to reduce trials

Search for Extragalactic Sources: Strategies



1. Look for hotspots in the neutrino sky $\rightarrow$ identify source candidates
2. Start from EM source catalog $\rightarrow$ look for neutrinos from source population
3. Focus on high-energy neutrinos with high signal probability $\rightarrow$ look for EM counterparts

Predefined source list

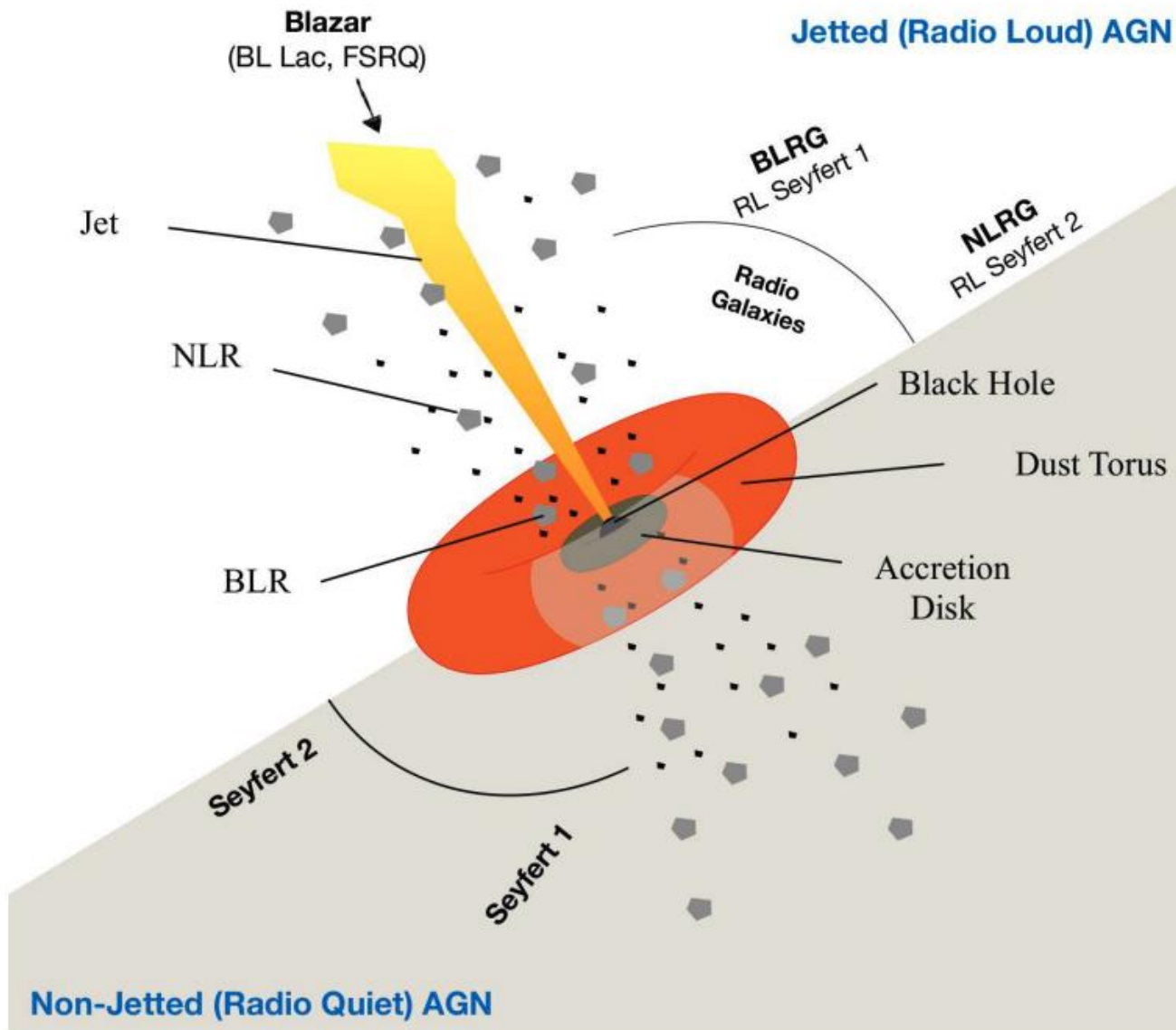
110 sources based on gamma-ray properties and weighted with neutrino search sensitivity

Why 110 sources: 5σ p-value before trials would remain above 4σ after trials

Extragal. Sources:

1. 4FGL sources weighted according to the integral gamma-ray flux > 1 GeV divided by the sensitivity neutrino flux for this analysis at the respective source declination
 2. 5% highest-weighted BL Lacs and flat spectrum radio quasars (FSRQs) are each selected
 3. minimum weighted integral flux from the combined selection of BL Lac and FSRQs is used as a flux threshold to include unidentified blazars and AGN
 4. Starburst galaxies in 4FGL (8 sources) are included in addition
- **87 extragalactic gamma-ray sources**

Unified AGN Model



Predefined source list

110 sources based on gamma-ray properties and weighted with neutrino search sensitivity

Why 110 sources: 5 σ p-value before trials would remain above 4 σ after trials

Galactic Sources:

1. VHE gamma-ray sources from TeVCat and gammaCat
 2. Convert gamma-ray spectra to equivalent neutrino fluxes, assuming a purely hadronic origin: $E_\nu = 2E_\gamma$
 3. Keep sources with predicted energy fluxes $> 50\%$ of neutrino sensitivity (assuming E^{-2} spectrum)
- **12 Galactic gamma-ray sources**

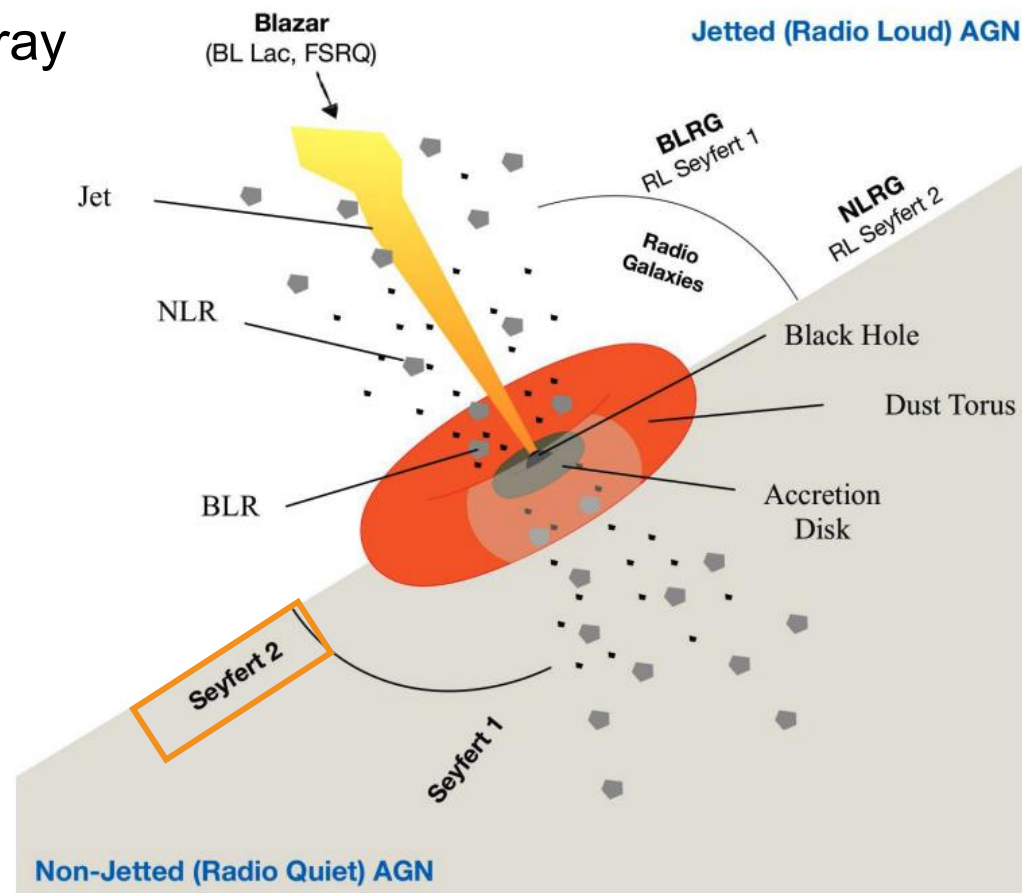
Predefined source list: Results

110 sources based on gamma-ray properties and weighted with neutrino search sensitivity

Most significant candidate:

NGC 1068 (M77), 4.2σ

- Nearby ($M=14\text{Mpc}$) Seyfert 2 galaxy
- AGN and star-forming activity



Combining gamma-ray source list with neutrino data
allowed neutrino source detection

Predefined source list: Results

110 sources based on gamma-ray properties and weighted with neutrino search sensitivity

Most significant candidate:

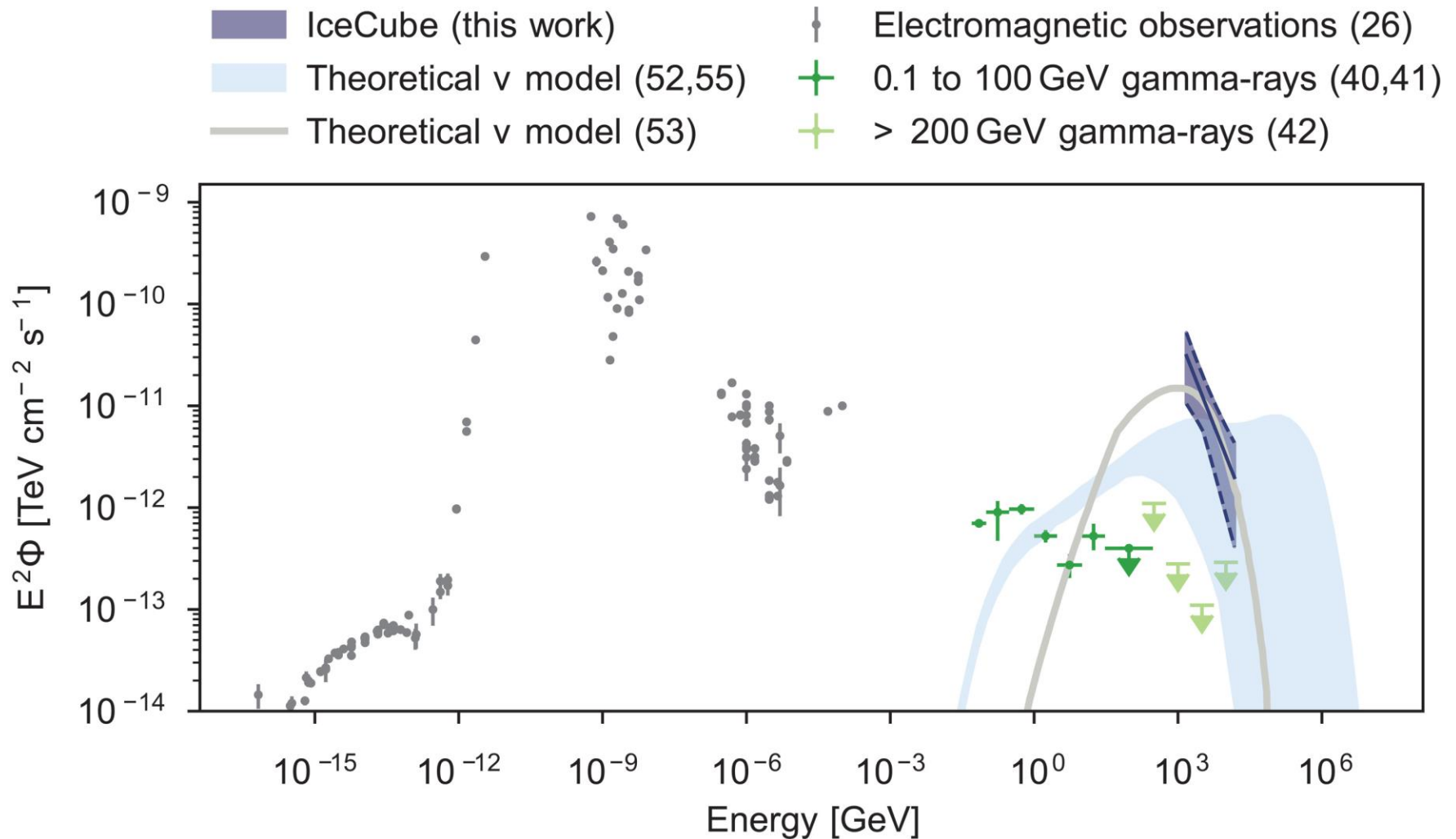
NGC 1068 (M77), 4.2σ

- Nearby ($M=14\text{Mpc}$) Seyfert 2 galaxy
- AGN and star-forming activity



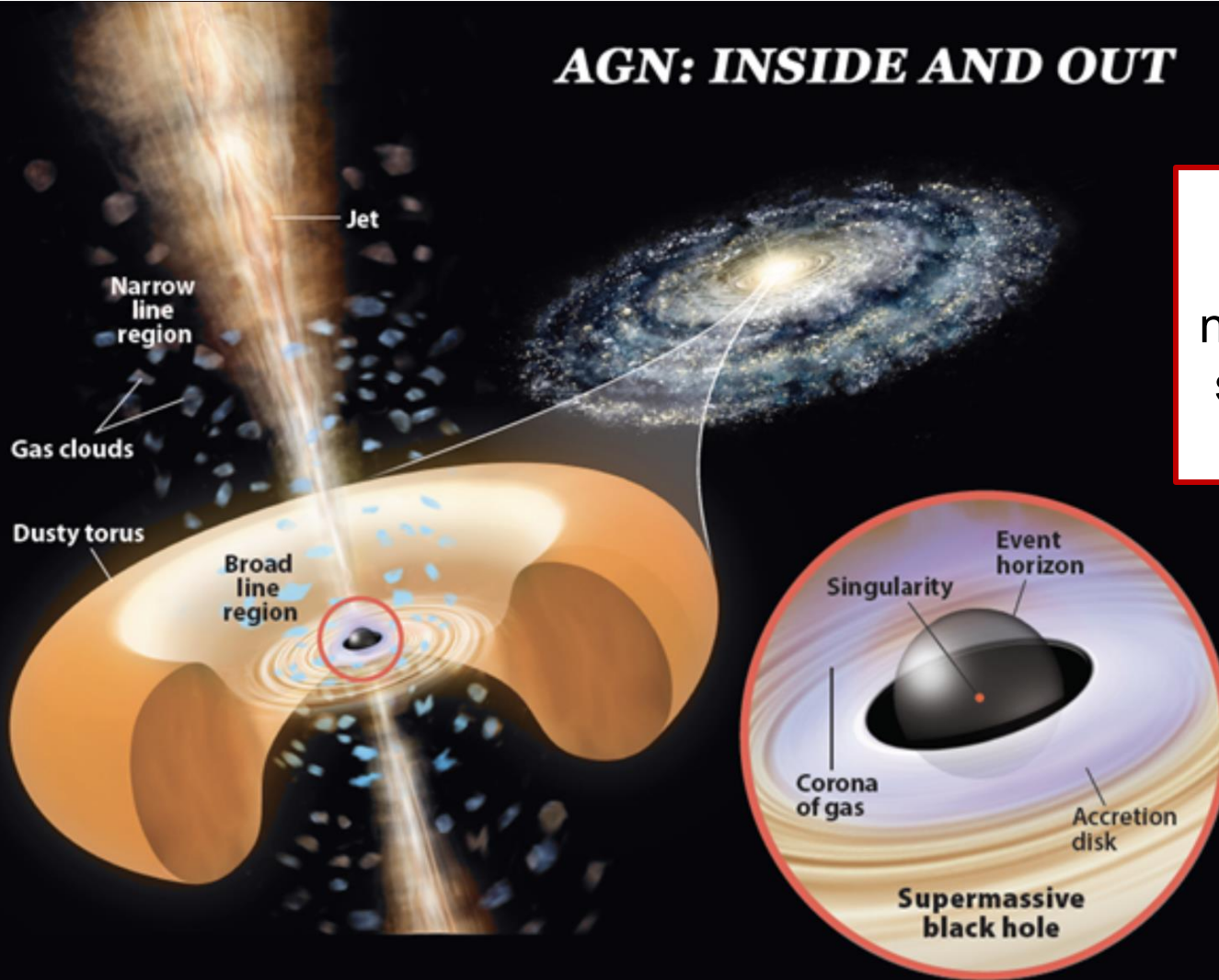
Combining gamma-ray source list with neutrino data
allowed neutrino source detection

Complete Multi-wavelength data of NGC 1068



Gamma rays need to be absorbed

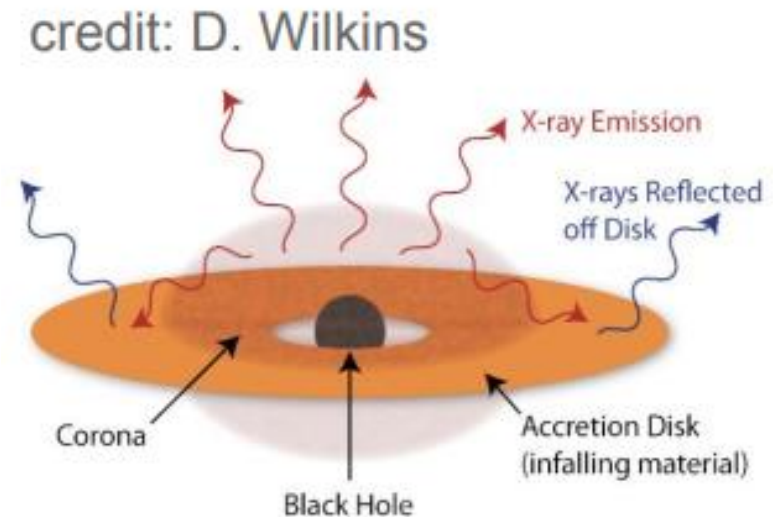
AGN: INSIDE AND OUT



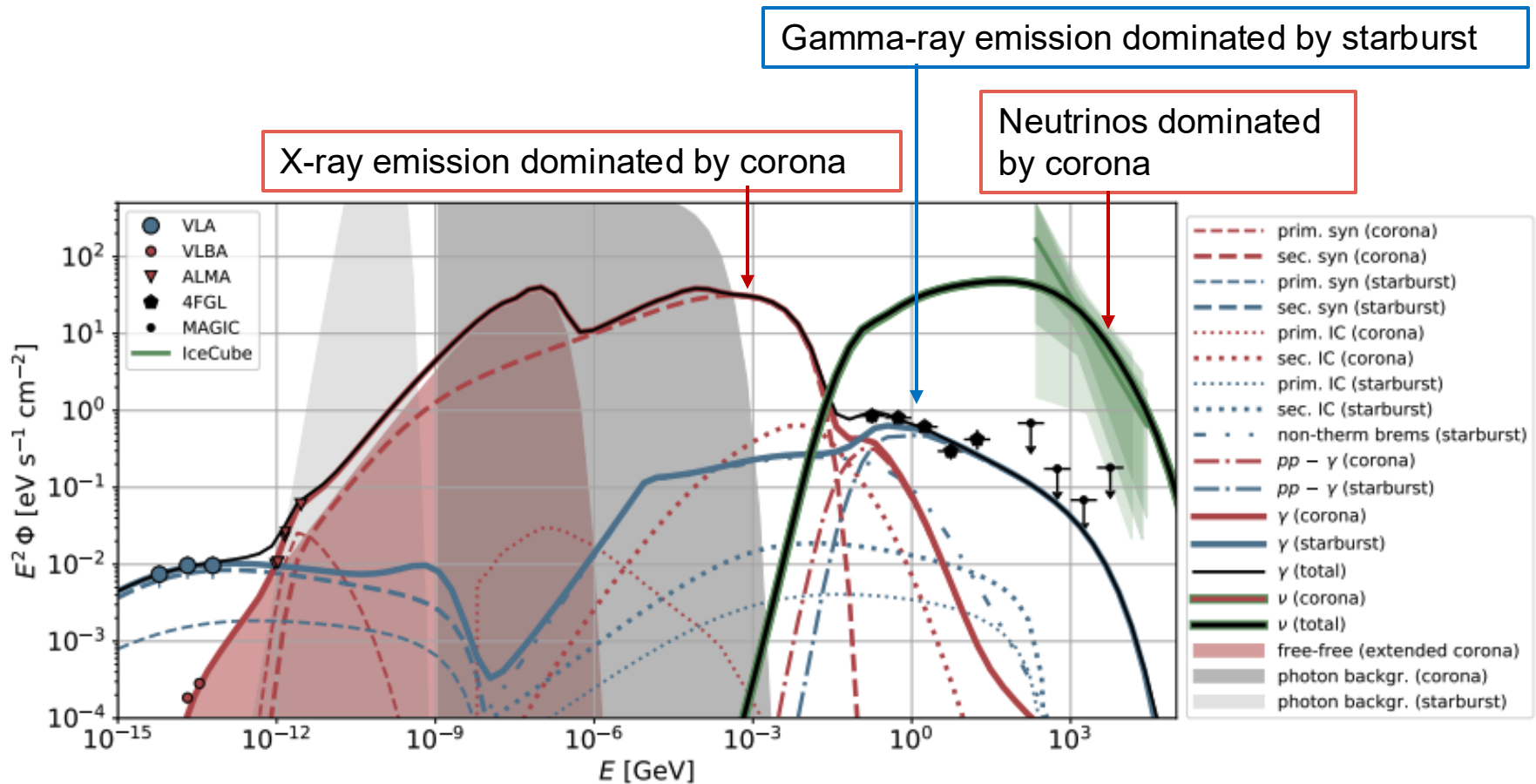
Lack of gamma rays places neutrino production site in the heart of the galaxy

What is the corona?

- region of very **hot plasma** (10^8 – 10^9 K) close to accretion disk (few **gravitational radii**)
- Accretion disk emits thermal **optical/UV photons**.
- In the corona: low-energy photons undergo **inverse Compton scattering** with energetic electrons → **X-rays**
- X-ray emission is observed in AGN spectra (**power-law component**)
- X-ray flux changes on timescales of hours to days → consistent with a **very compact region**

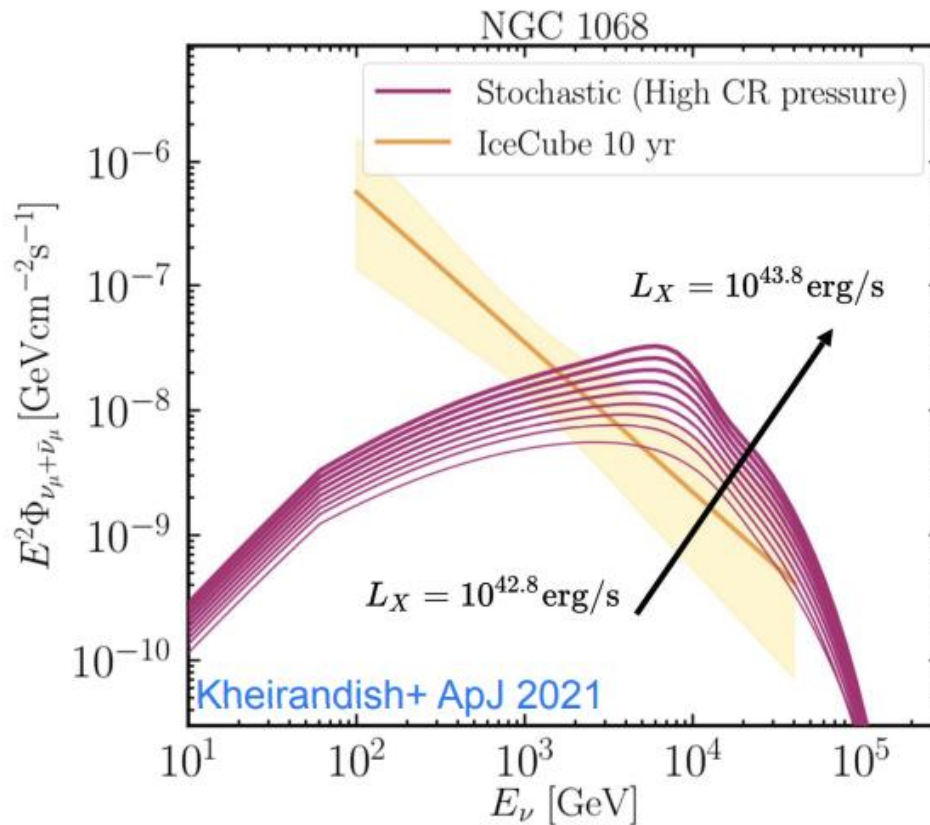


Detailed Modeling of Starburst-AGN Composite

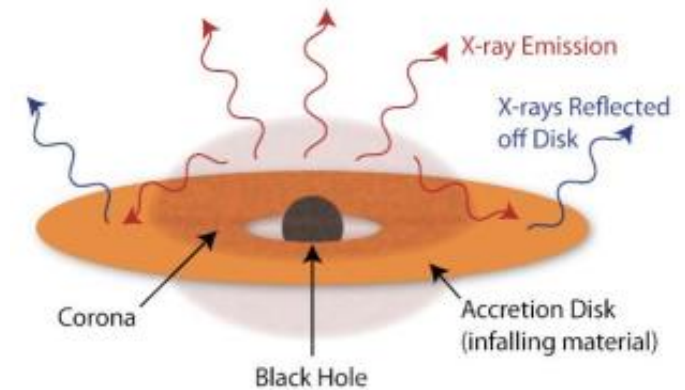


More Seyferts?

Assumption: Neutrino flux scales with intrinsic X-ray flux (2–10 keV)



credit: D. Wilkins



Stacking Search

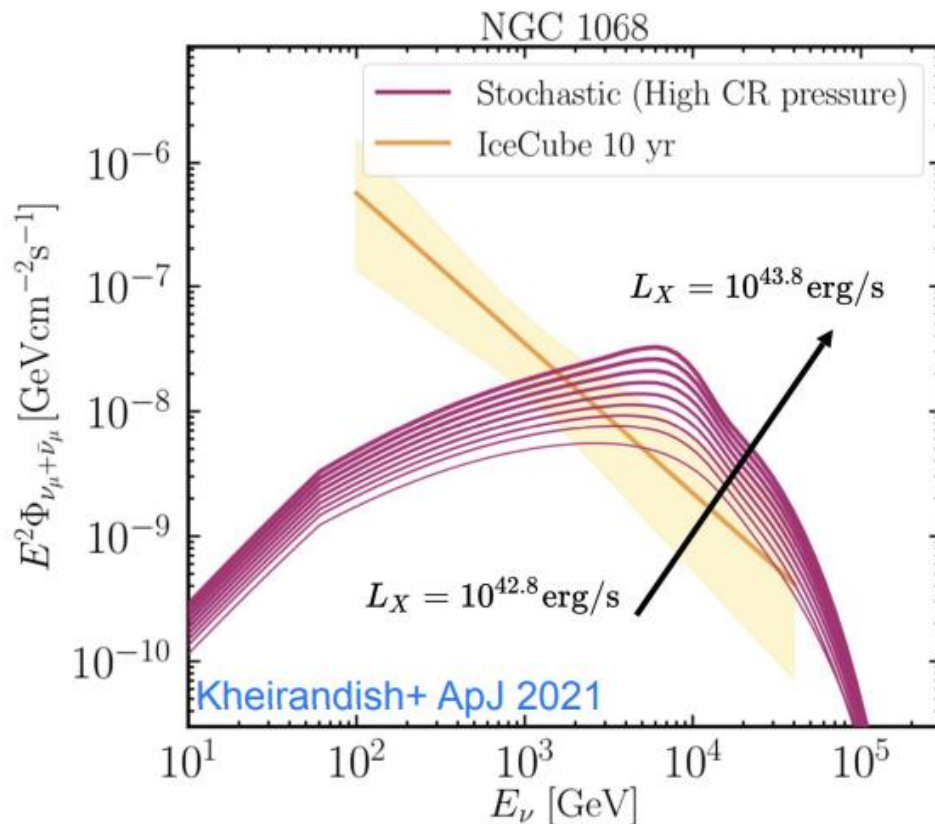
- Use catalog of sources (here: Seyferts)
 - Assume neutrino spectral shape (often power-law, here: disk corona model prediction)
 - Assume weighting scheme between neutrino sources (here: neutrino flux follows X-ray flux)
- only one fit parameter for entire population: overall normalization

Potential problem: if assumption (weighting) does not hold, sensitivity is reduced

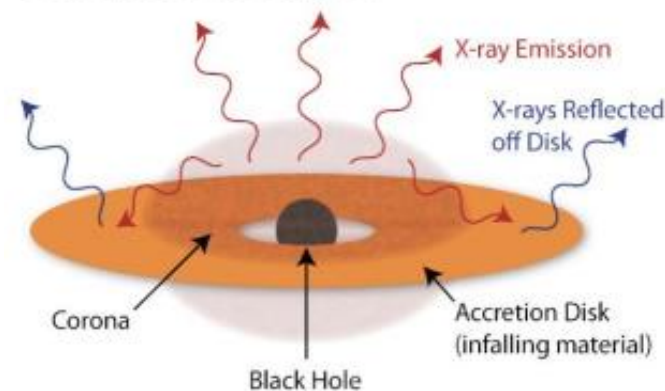
Alternative: fit each individually and combine results with binomial test

More Seyferts: Stacking results

Assumption: Neutrino production in disk corona, intrinsic X-ray flux (2–10 keV) as proxy for neutrino emission



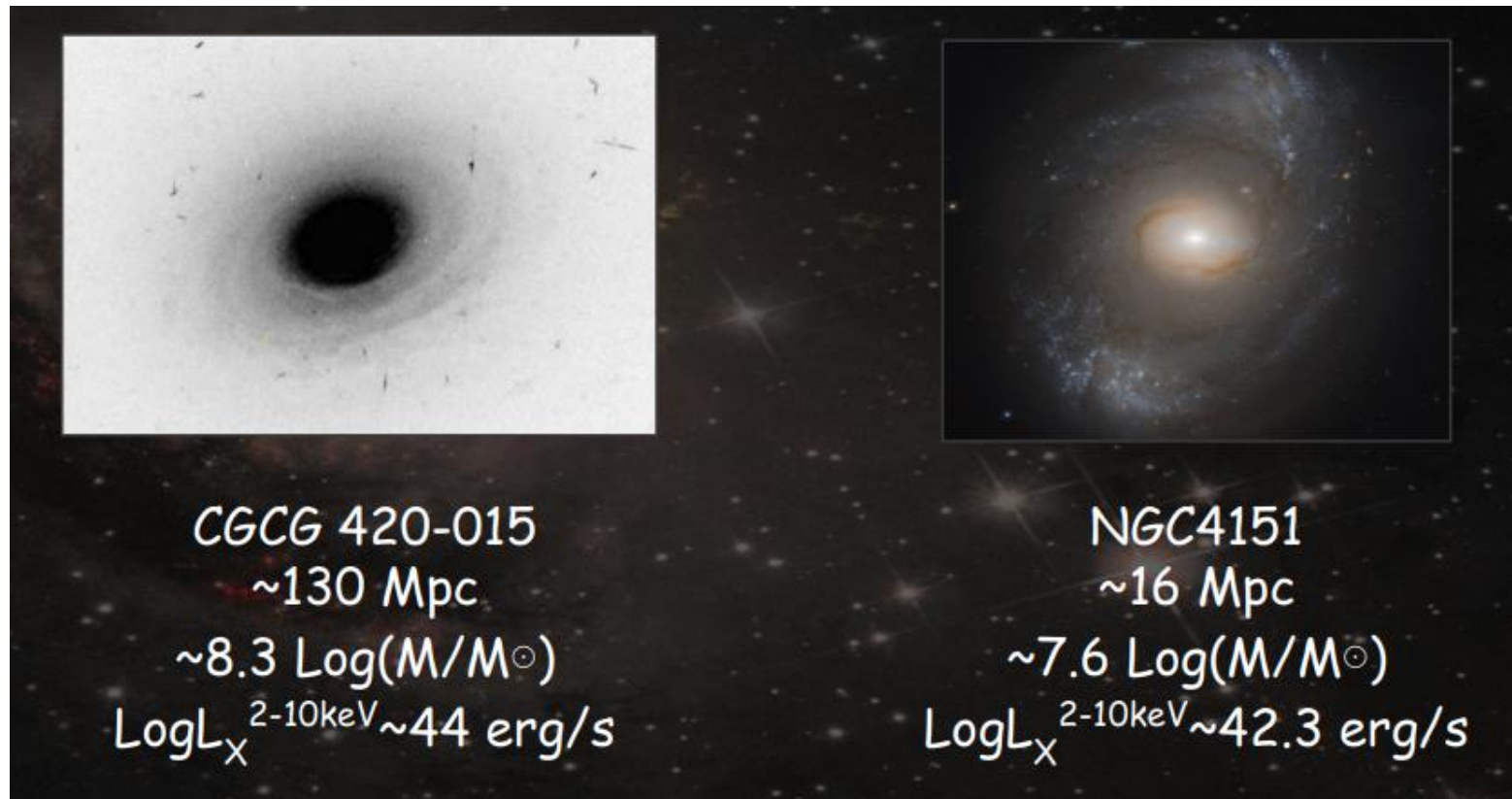
credit: D. Wilkins



No significant emission is found in the stacking search excluding NGC 1068.

More Seyferts: Individual sources

No assumption about neutrino emission model



Two more source candidates at 2.5σ and 2.1σ level

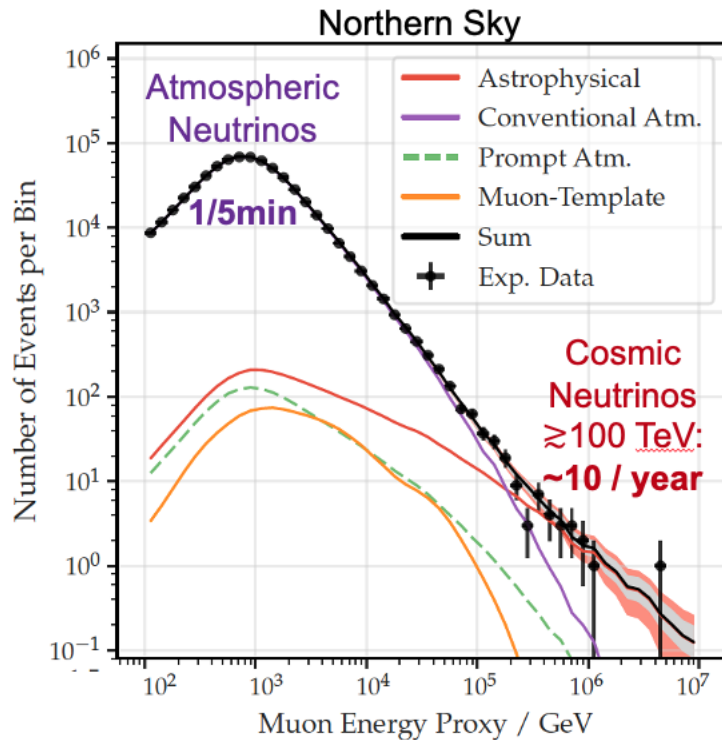


Take home message

Seyferts are currently the most promising neutrino source class.

However, it seems we do not fully understand their emission processes.

Search for Extragalactic Sources: Strategies



1. Look for hotspots in the neutrino sky $\rightarrow$ identify source candidates
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3. **Focus on high-energy neutrinos with high signal probability $\rightarrow$ look for EM counterparts**

Neutrinos as Triggers

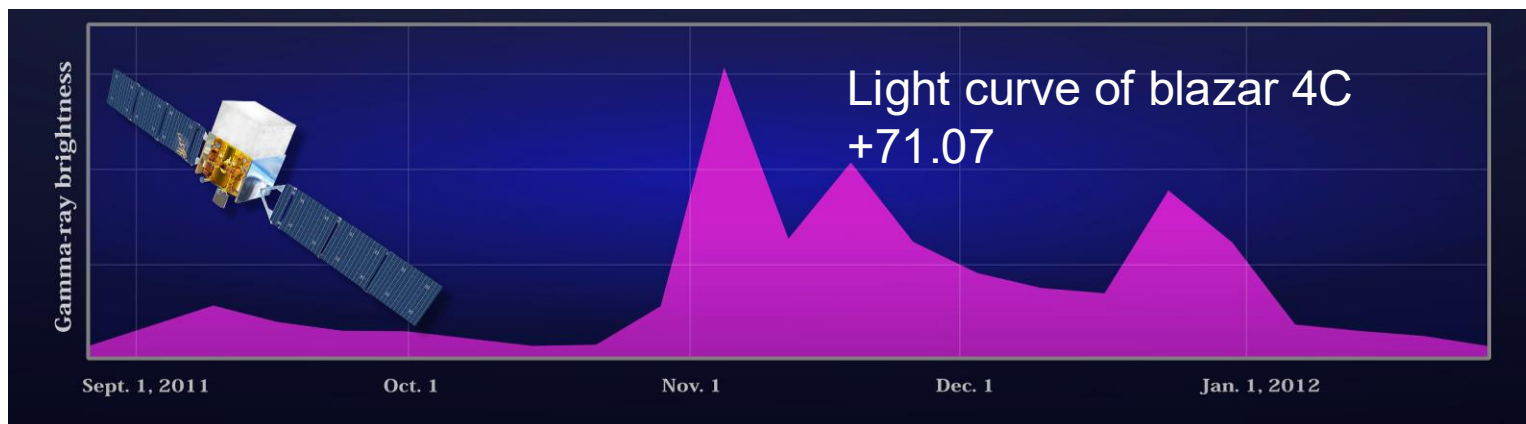
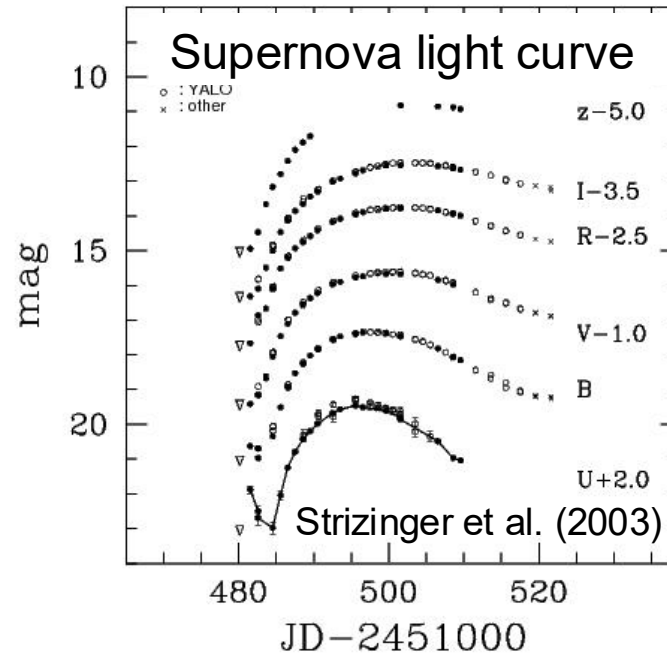
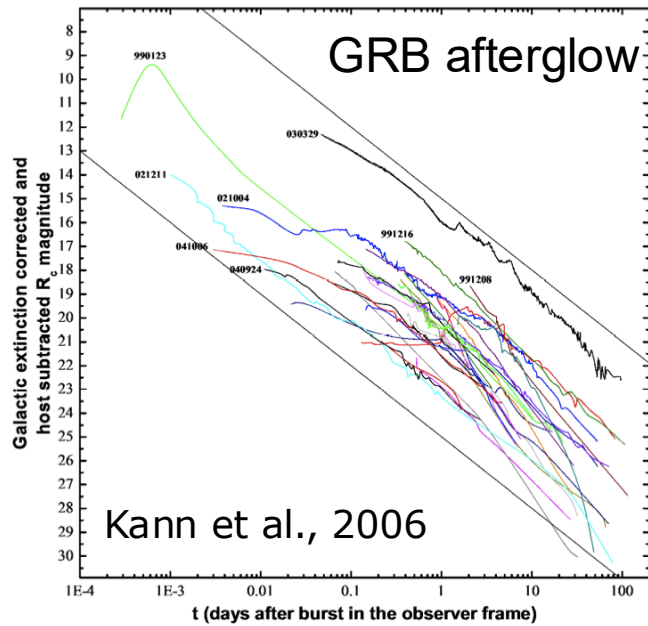
Public alerts since April 2016

- Single high-energy muon track events ($> \sim 100\text{TeV}$)
- “Gold” alert stream: 10 / yr, ~ 5 / yr of cosmic origin
- Median latency: 30 sec

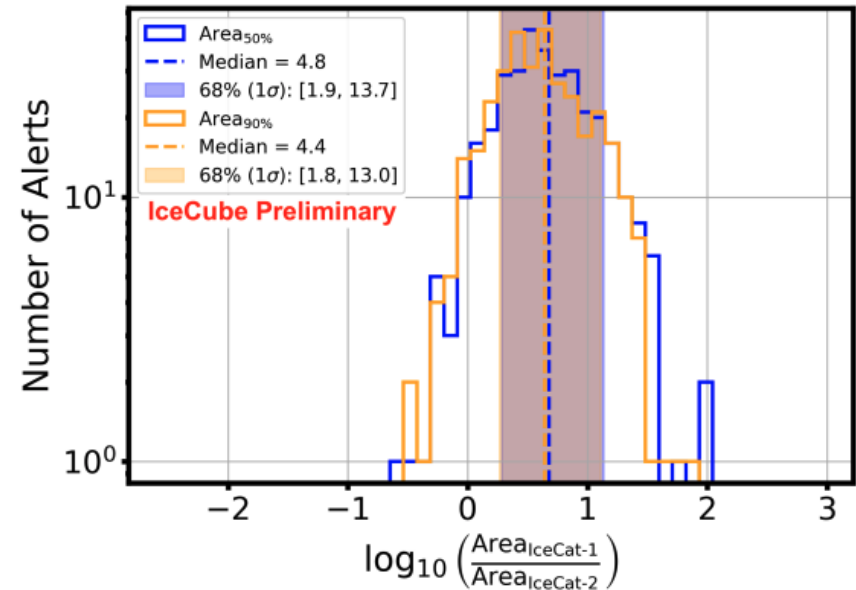
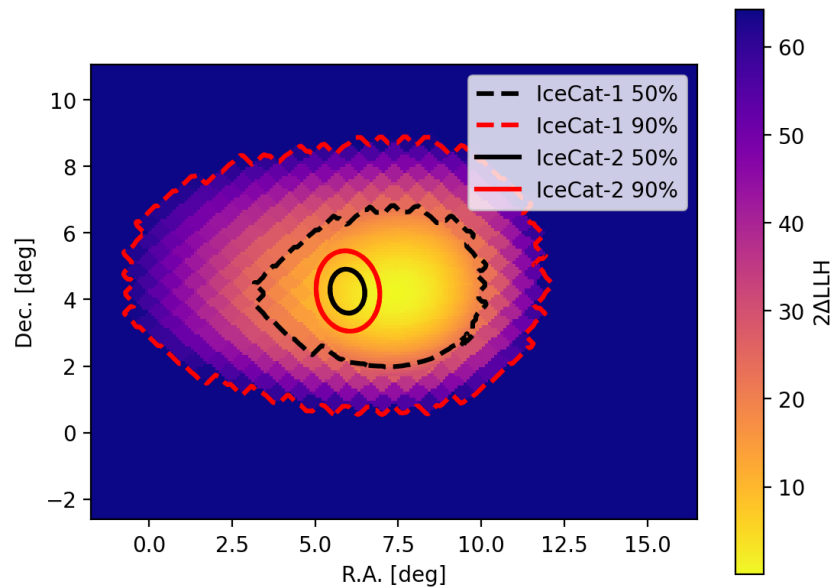
Goal: Find electromagnetic counterpart



Expected EM Counterparts



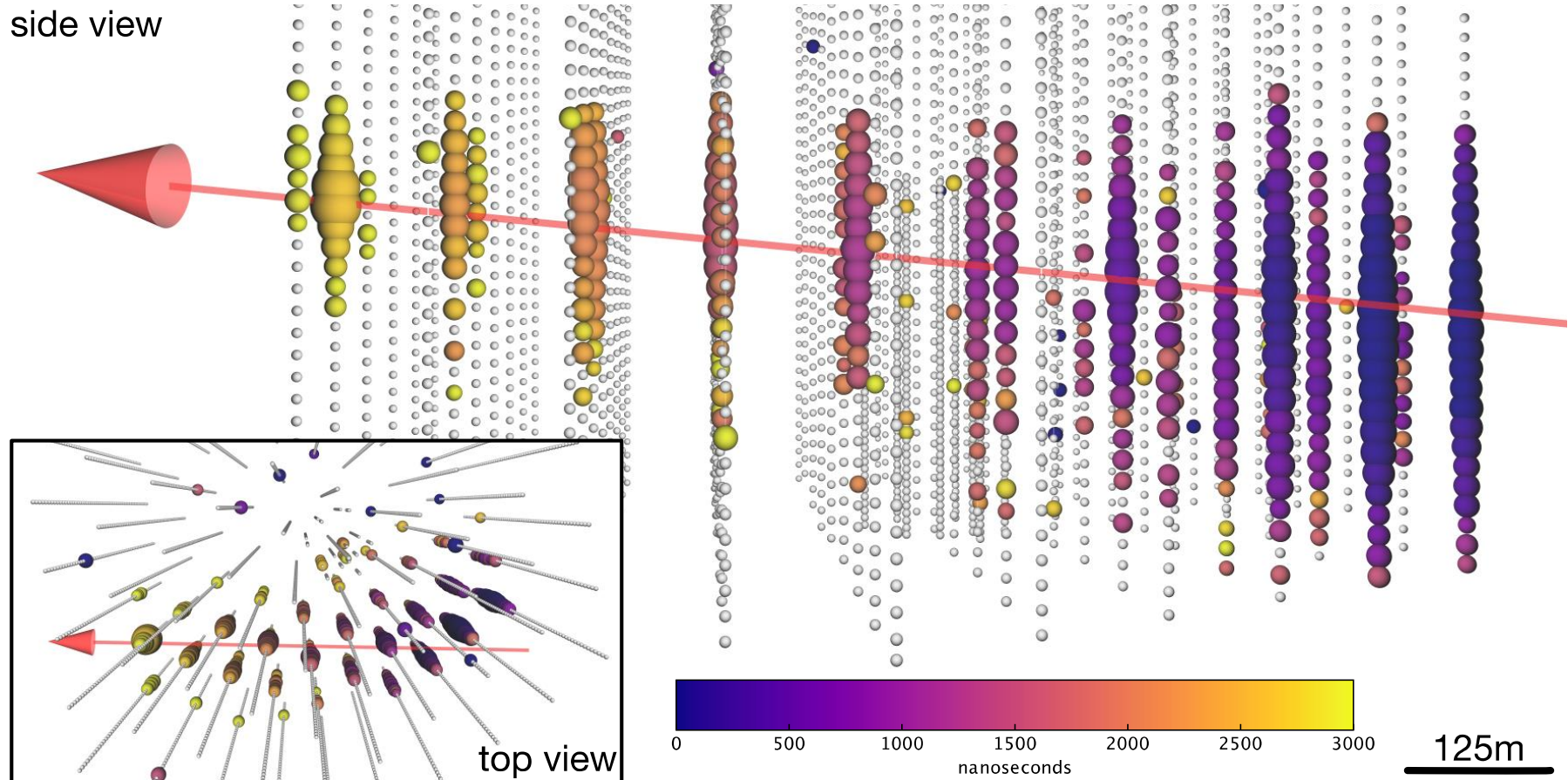
IceCat-2



Improvement in 90% uncertainty area of factor 4.4

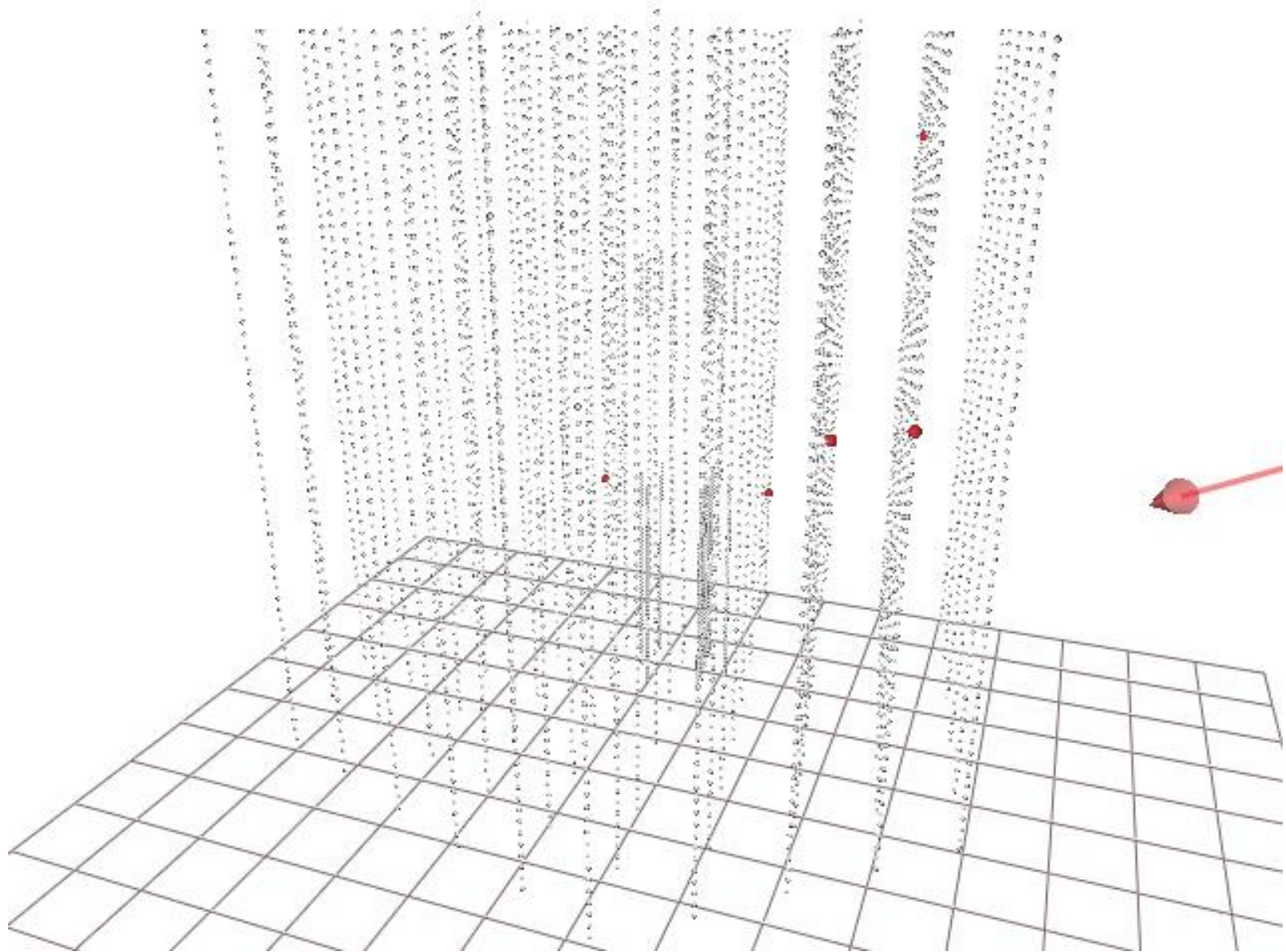
Archival alert catalog (IceCat-2) under development

First example: IC-170922A – a 290 TeV Neutrino

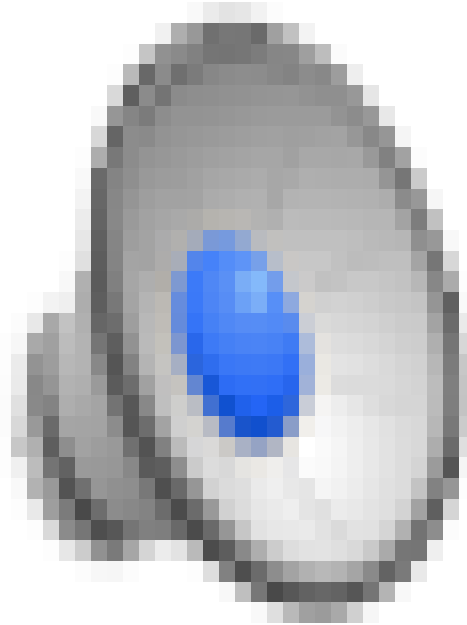


Signalness: 56.5%

IC-170922A – a 290 TeV Neutrino

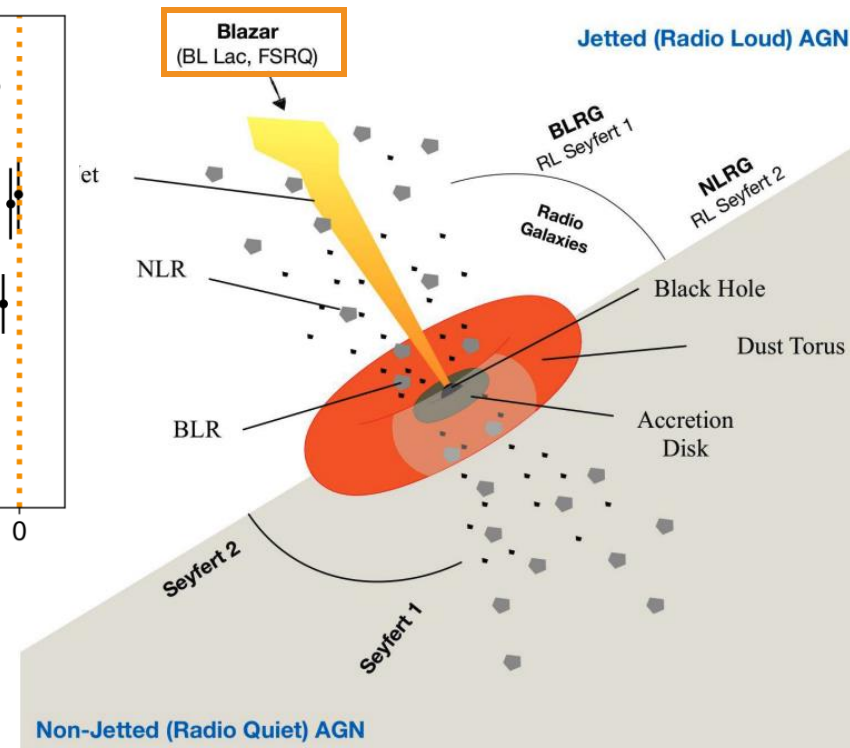
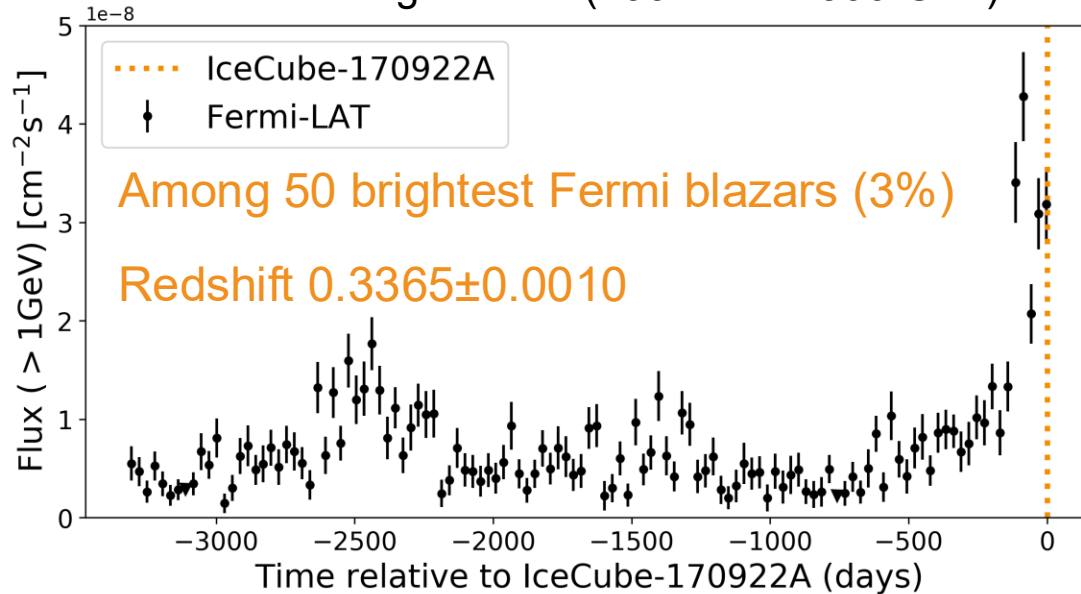


Fermi-LAT finds Flaring Source



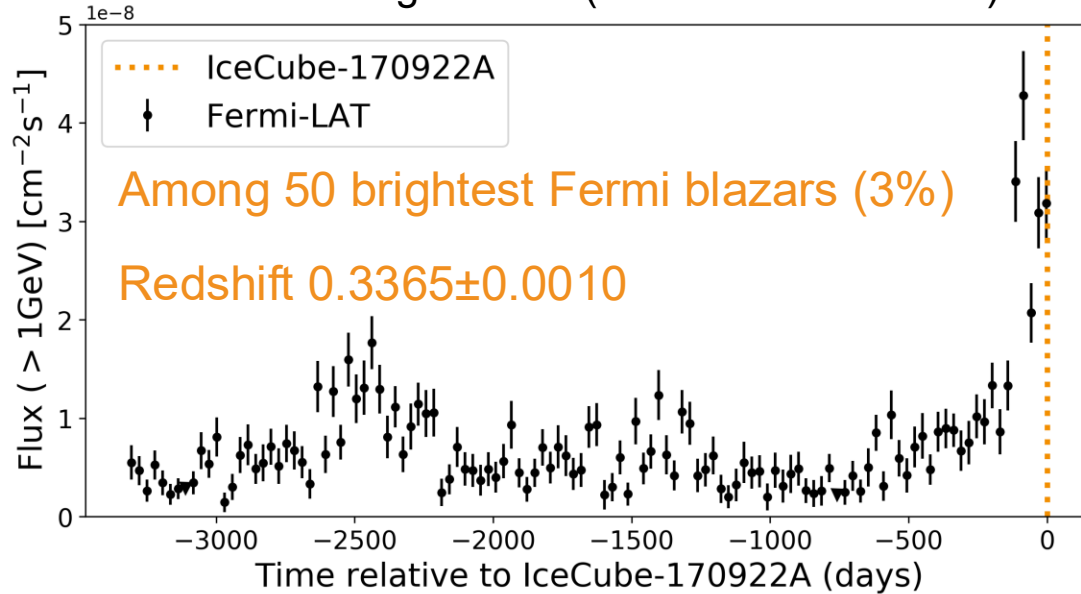
Fermi-LAT finds Flaring Blazar, TXS 0506+056

Fermi-LAT light curve (100 MeV – 300 GeV)

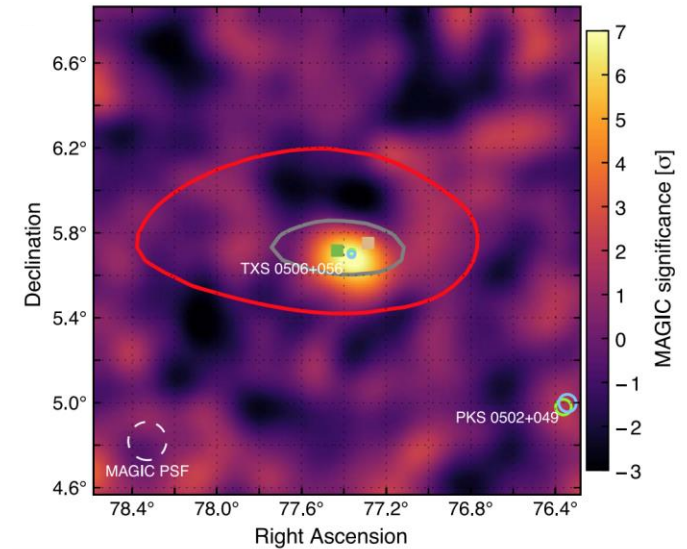


Fermi-LAT finds Flaring Blazar, TXS 0506+056

Fermi-LAT light curve (100 MeV – 300 GeV)

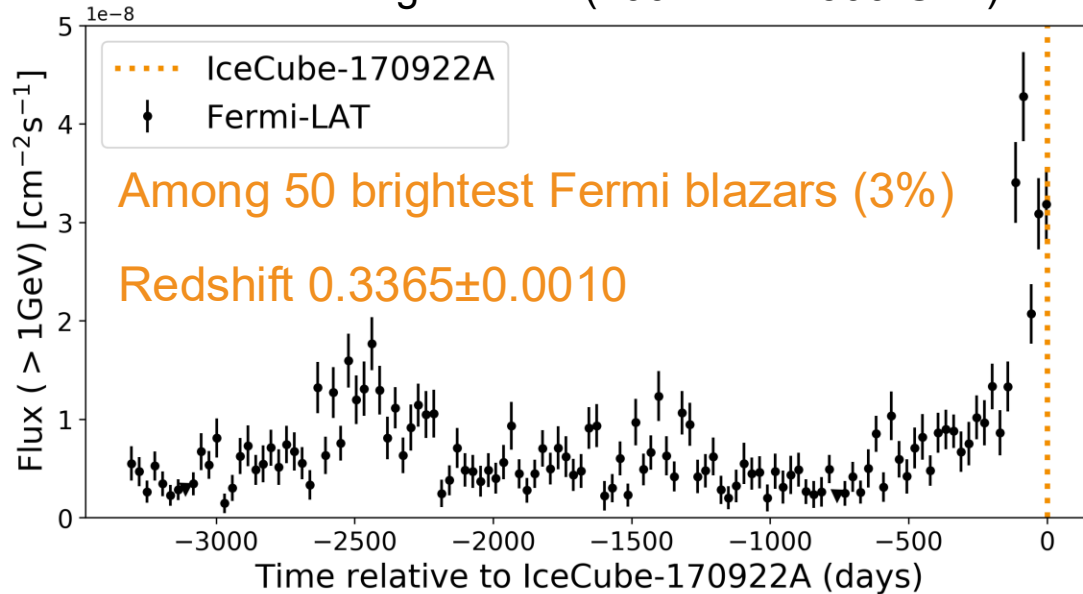


MAGIC Significance map:
TeV gamma rays

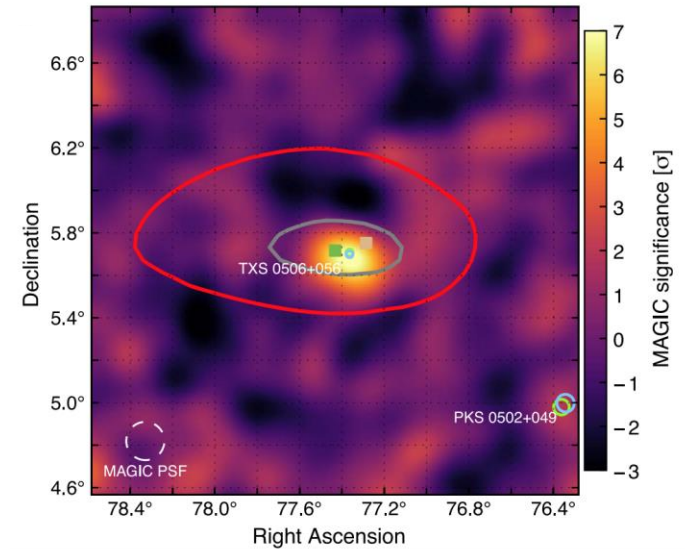


Fermi-LAT finds Flaring Blazar, TXS 0506+056

Fermi-LAT light curve (100 MeV – 300 GeV)

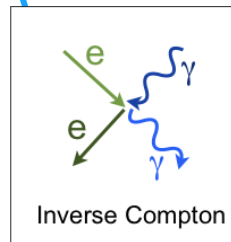
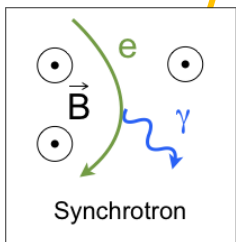
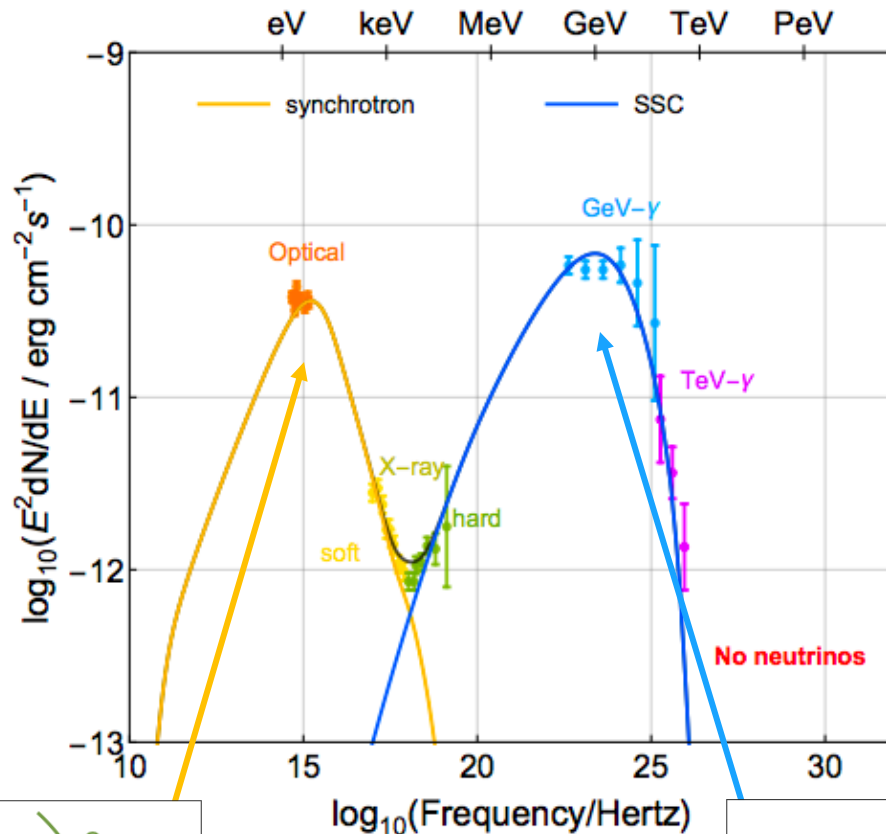


MAGIC Significance map:
TeV gamma rays



3 sigma significance including trials
> 6 PeV protons accelerated in the source

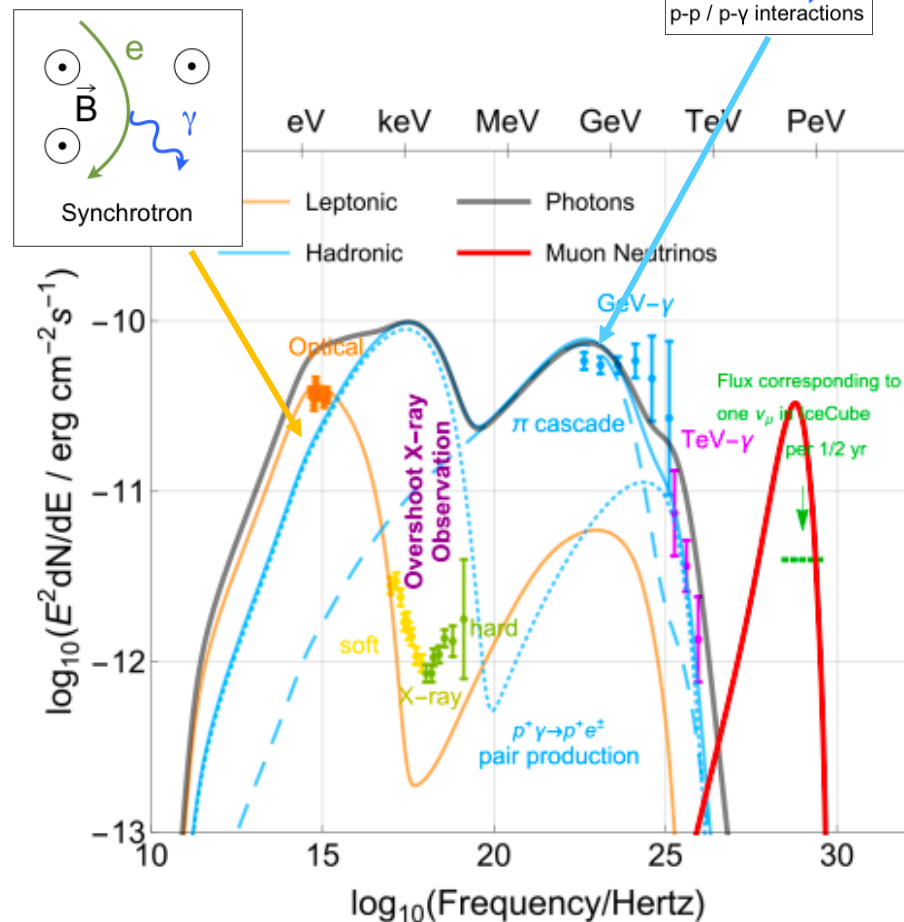
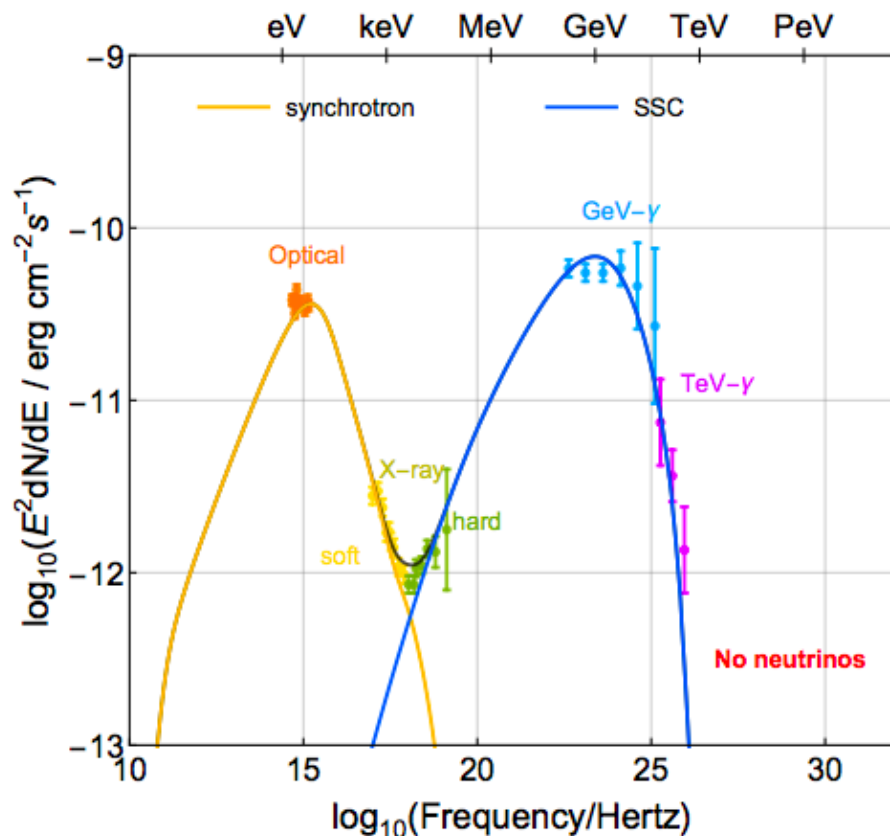
Modeling – leptonic



Gao et al., Nature Astronomy 2018,

Keivani et al., ApJ, 2018, MAGIC Coll., ApJ, 2018, Cerruti et al. MNRAS 2018, ...

Modeling – leptonic, hadronic

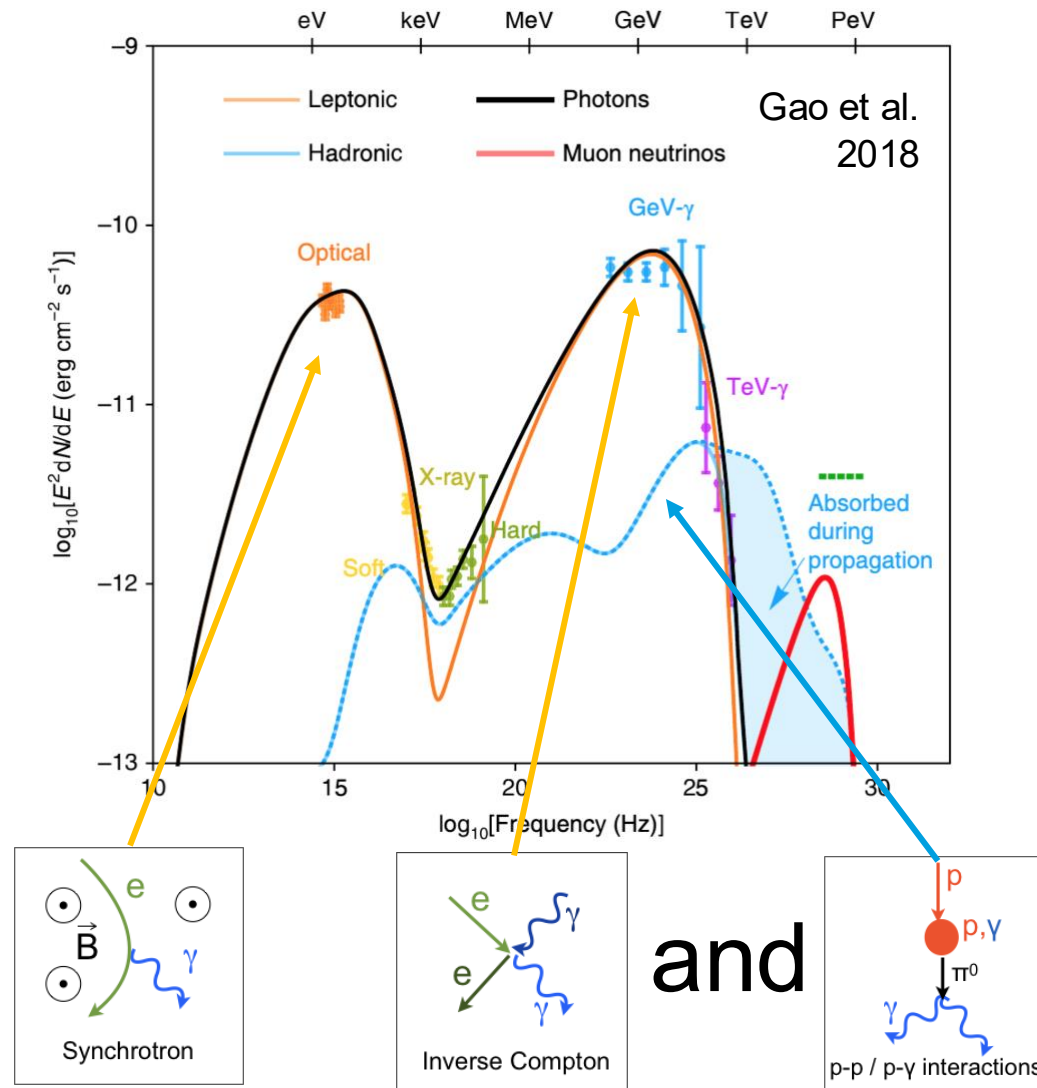


**Simple one-zone hadronic models violate X-ray constraints
→ More complex models needed**

Gao et al., Nature Astronomy 2018,

Keivani et al., ApJ, 2018, MAGIC Coll., ApJ, 2018, Cerruti et al. MNRAS 2018, ...

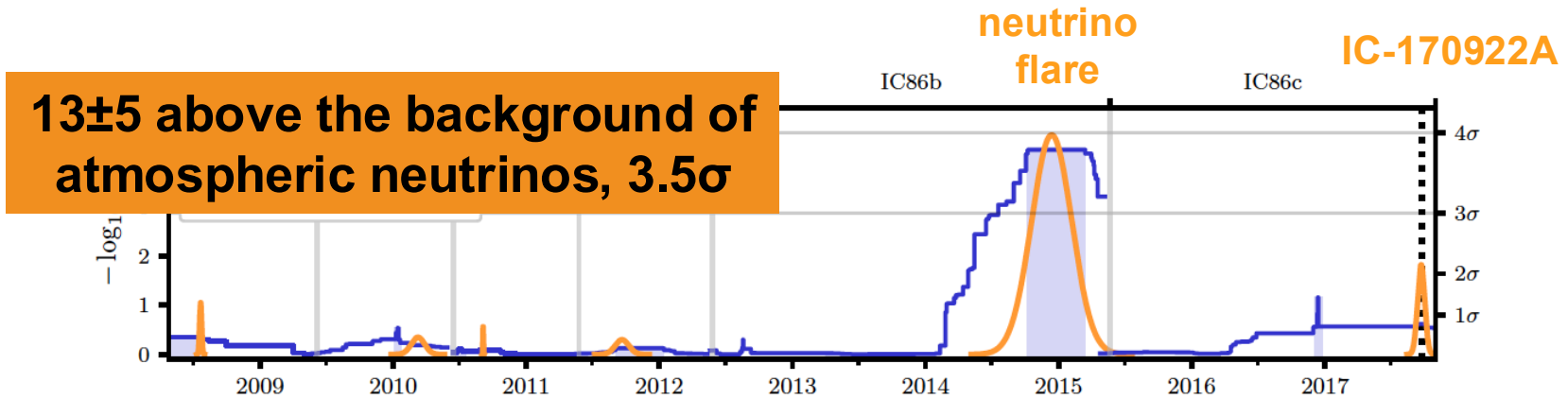
Modeling – lepto-hadronic



Gao et al., Nature Astronomy 2018,

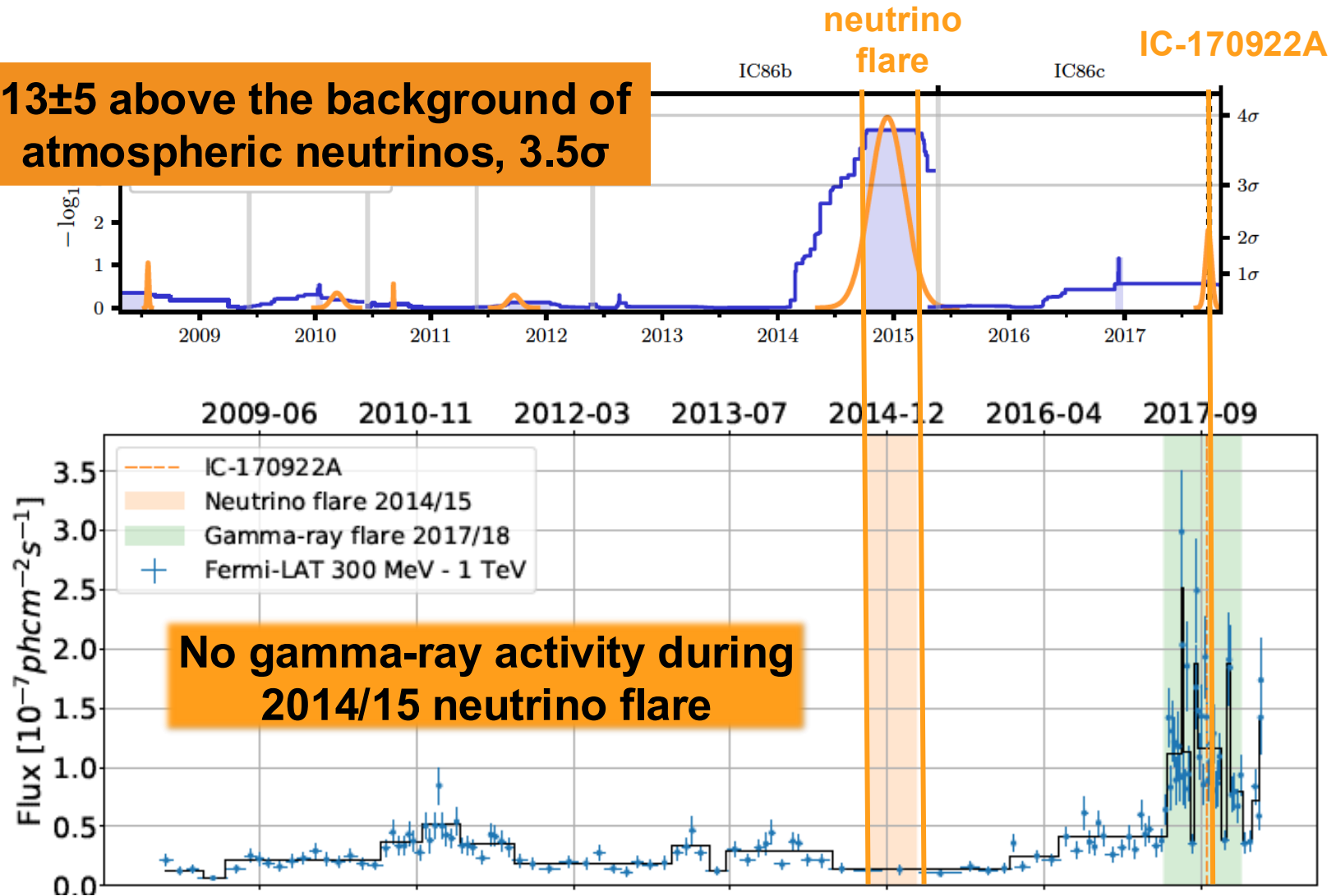
Keivani et al., ApJ, 2018, MAGIC Coll., ApJ, 2018, Cerruti et al. MNRAS 2018, ...

Are there more Neutrinos from this Source?



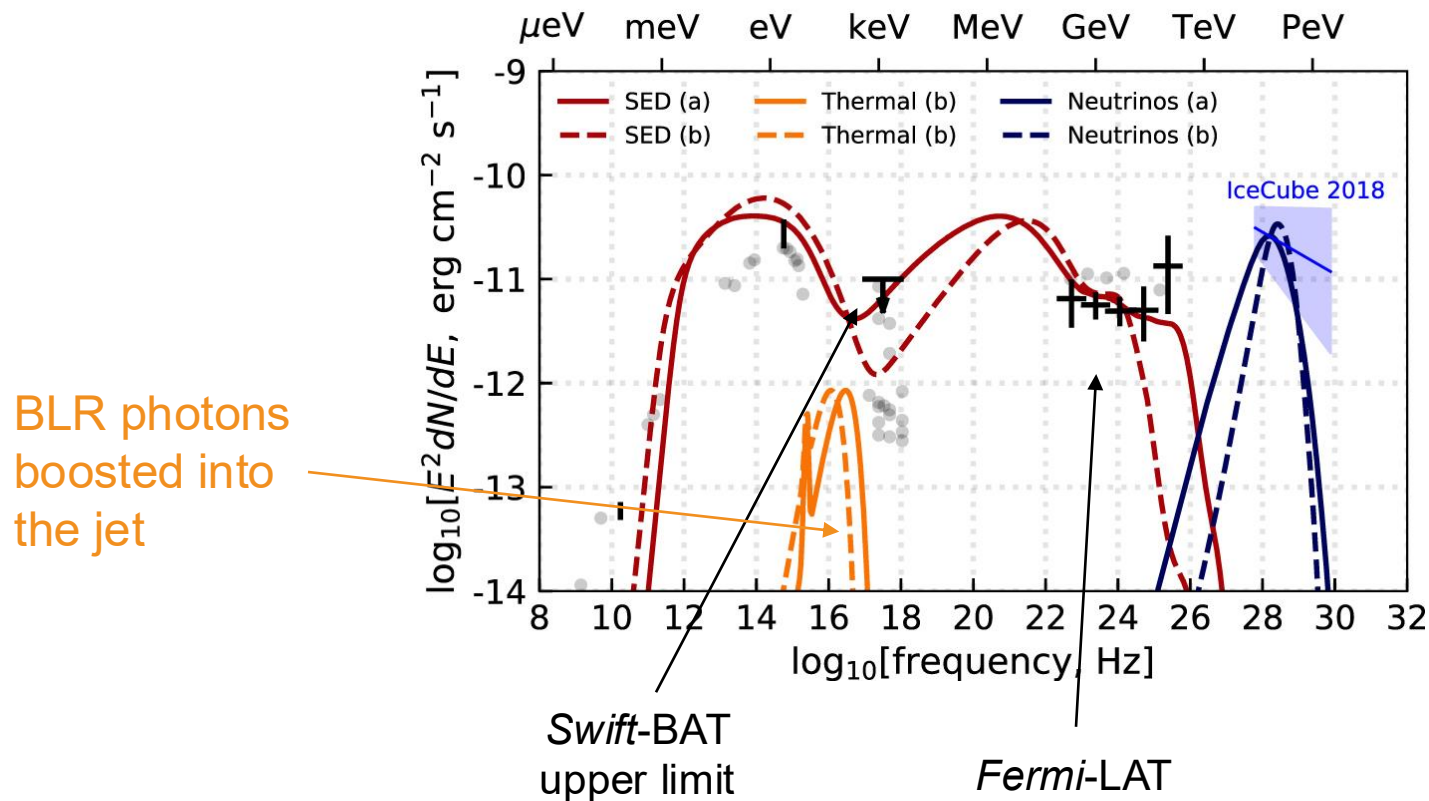
Is there also a Gamma-ray Flare?

13 ± 5 above the background of atmospheric neutrinos, 3.5σ



Modeling of 2014/15 neutrino flare

neutrino luminosity is ~ 4 times higher than gamma-ray luminosity
 $\rightarrow$ **challenge for models**

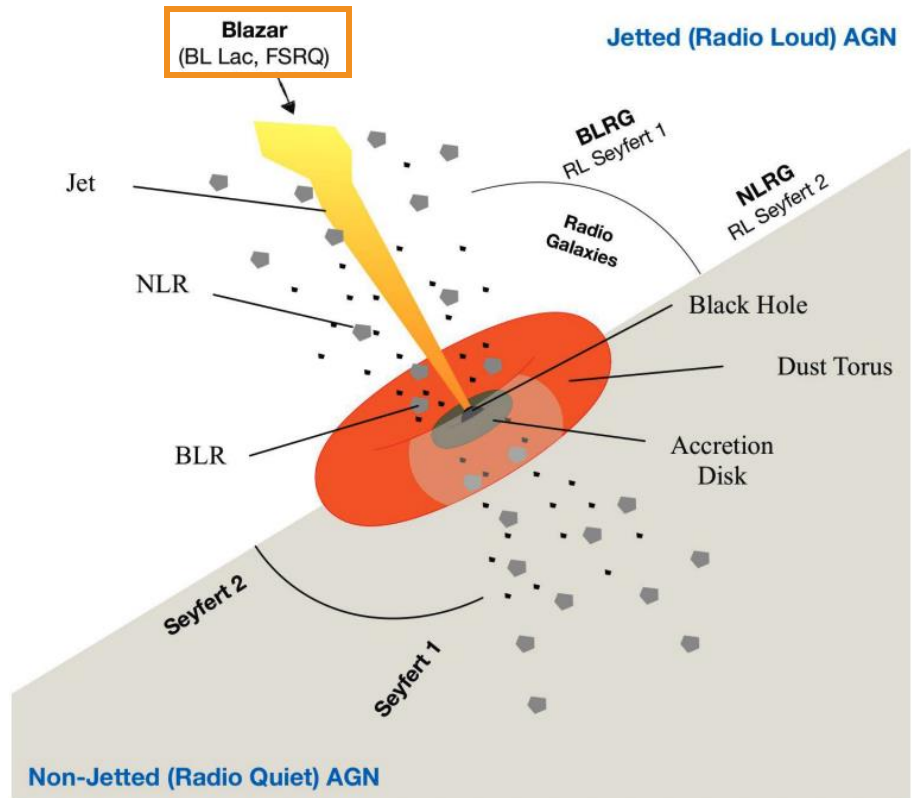
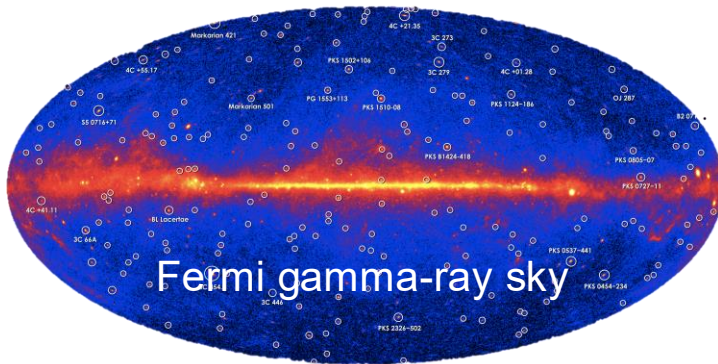


see e.g. Rodrigues et al. ApJL 874 2019, A. Reimer et al. ApJ 881 2019,
F. Halzen et al. ApJL 874 2019

**Do gamma-ray blazars
produce all diffuse neutrinos?**

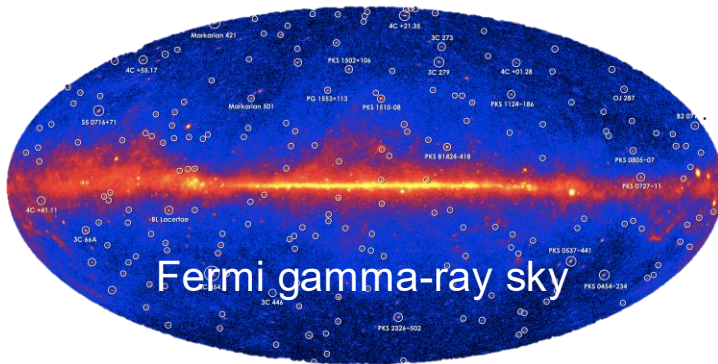
Fermi Blazars

Gamma rays tell us **where** to look for neutrinos

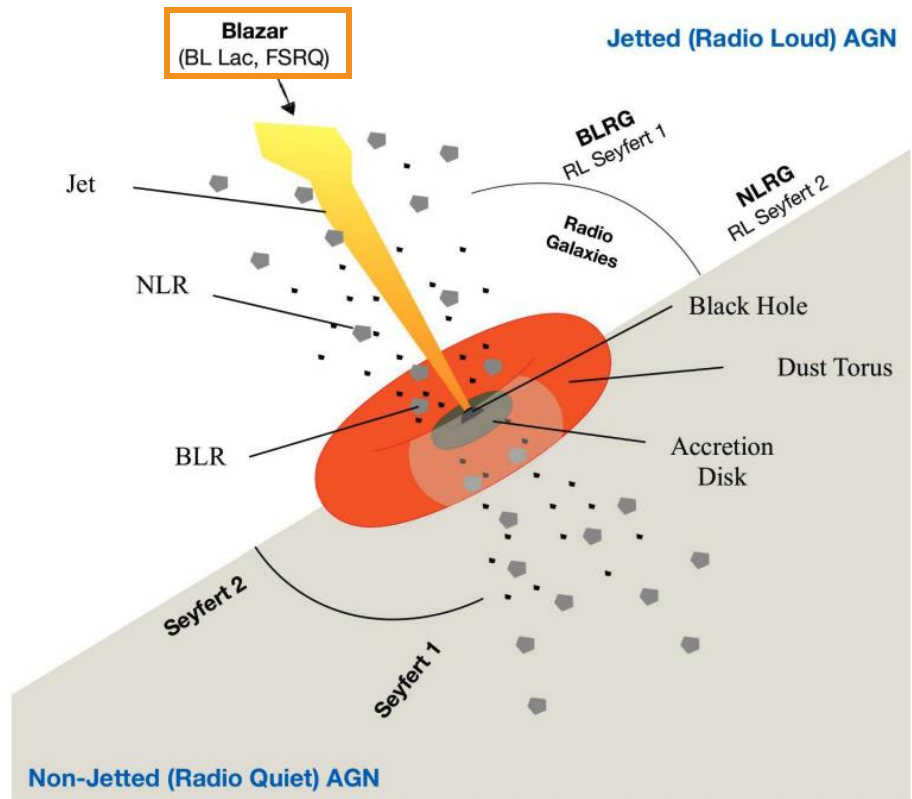


Fermi Blazars

Gamma rays tell us **where** to look for neutrinos



Correlation study of 12 years of IceCube data and 2089 **Fermi-LAT blazars**

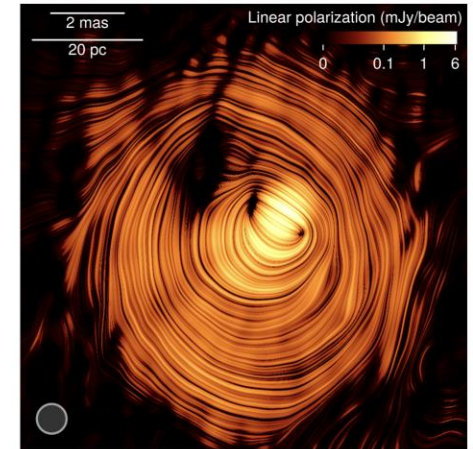
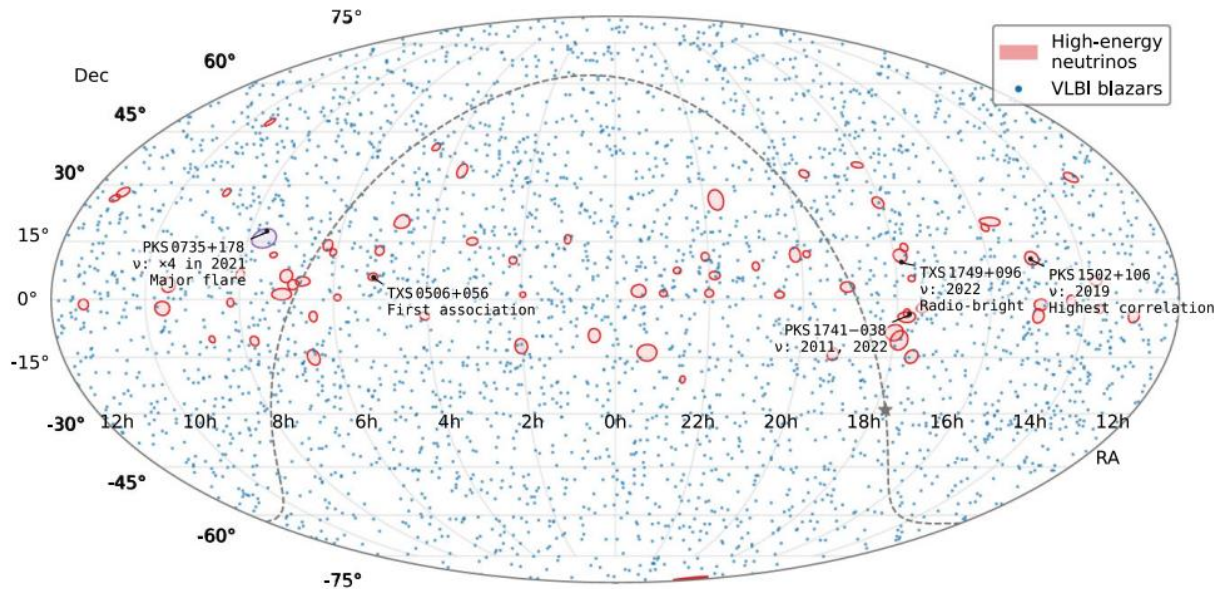


Fermi-LAT blazars can only be responsible for a **small fraction** of the observed ν 's.

**Other wavelength as tracer for
neutrinos from blazars?**

Radio Core Brightness as Tracer

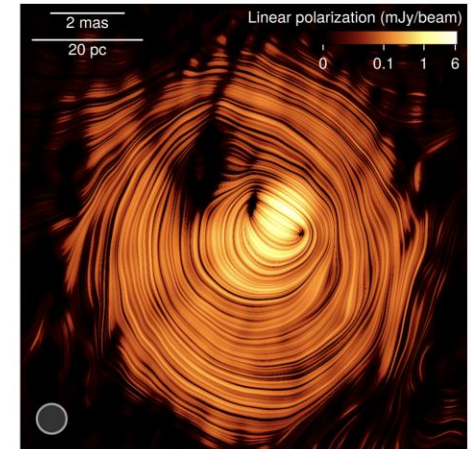
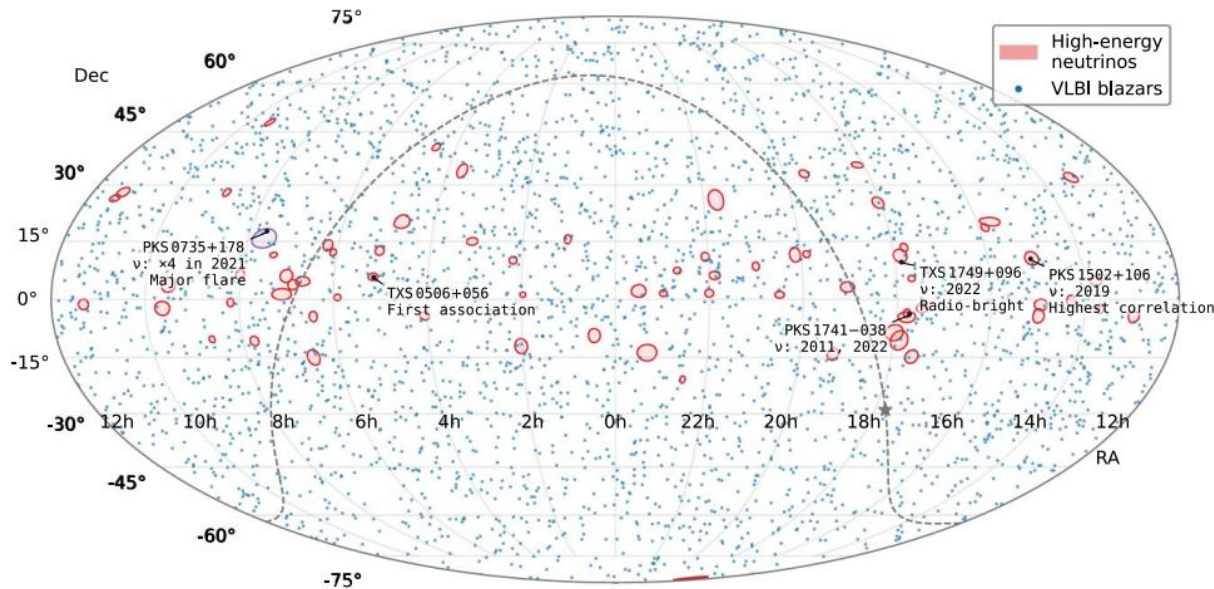
Correlation with VLBI-flux-density limited sample of blazars



Correlation of radio-bright blazars with
IceCube neutrino alerts at chance
coincidence of $3 \cdot 10^{-4}$

Radio Core Brightness as Tracer

Correlation with VLBI-flux-density limited sample of blazars



Correlation of radio-bright blazars with IceCube neutrino alerts at chance coincidence of $3 \cdot 10^{-4}$

Correlation not confirmed by slightly different analyses

Zhou et al. 2021 PRD 103 (2021),
IceCube ApJ 954 (2023)

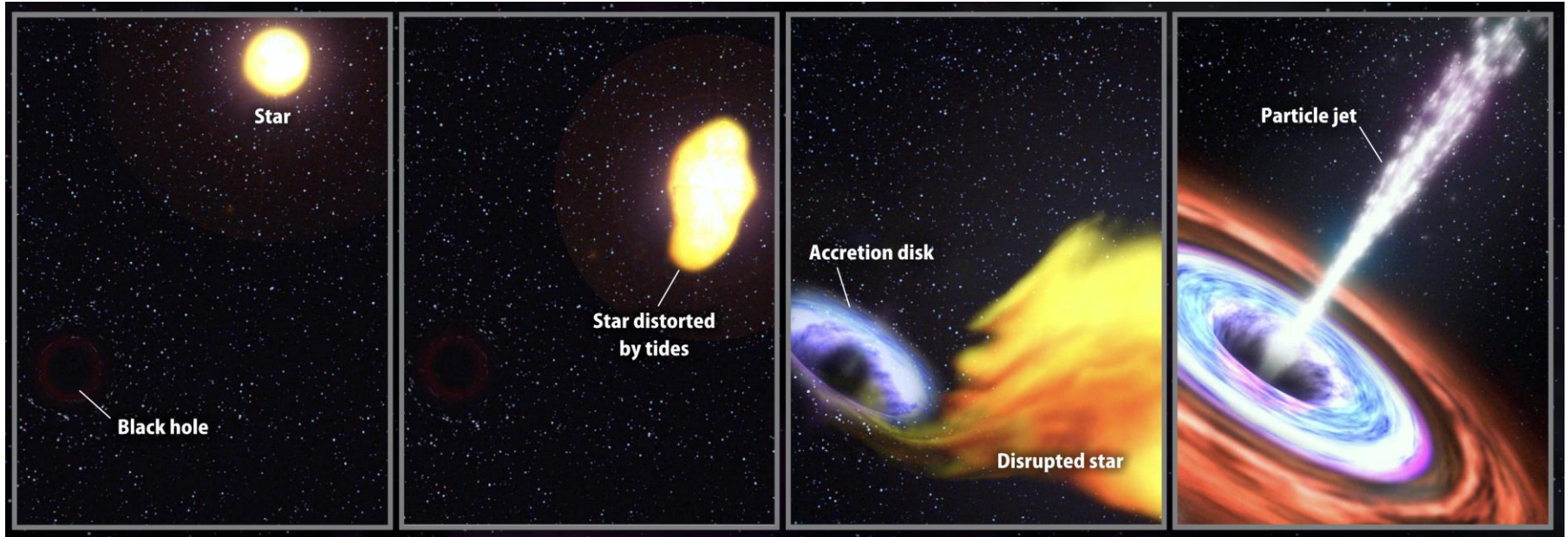


Take home message

Gamma-ray emission in blazars could not be established as a universal tracer for neutrino emission, but blazars are not ruled out as neutrino sources.

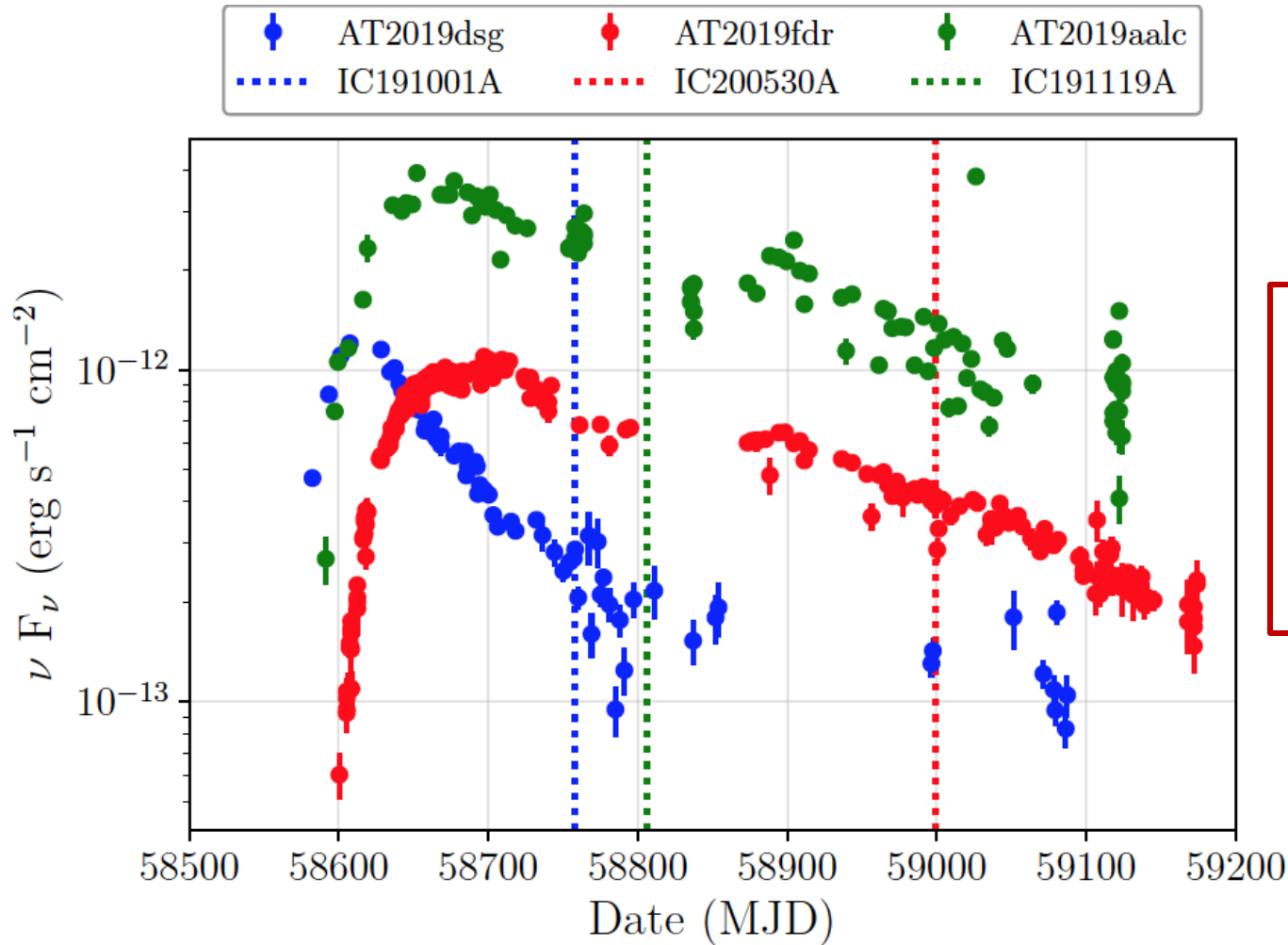
Other Sources?

What about quiet black holes? Tidal Disruption Events



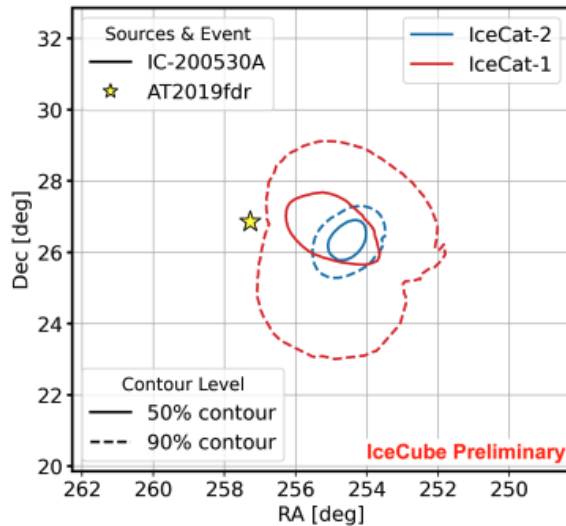
~50 TDEs identified, 4 jetted TDEs
bright in optical

Tidal Disruption Events as Neutrino Candidates

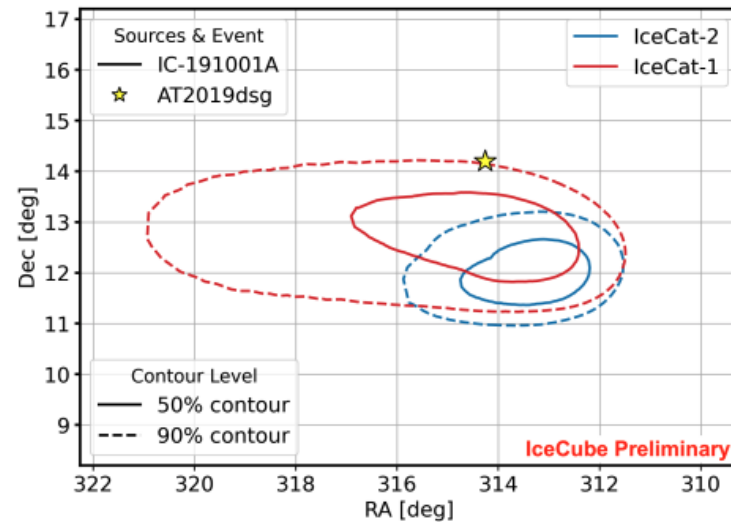


First hint of neutrino production in TDEs
→ **Very efficient** neutrino production in TDEs compared to AGN?

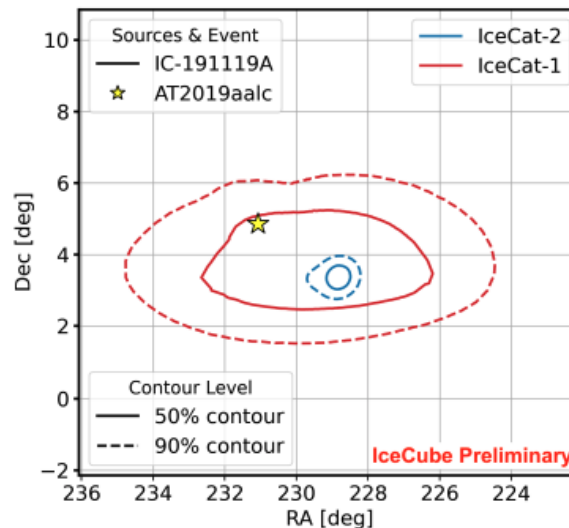
Neutrino TDEs disfavored with IceCat-2



(a) AT2019fdr

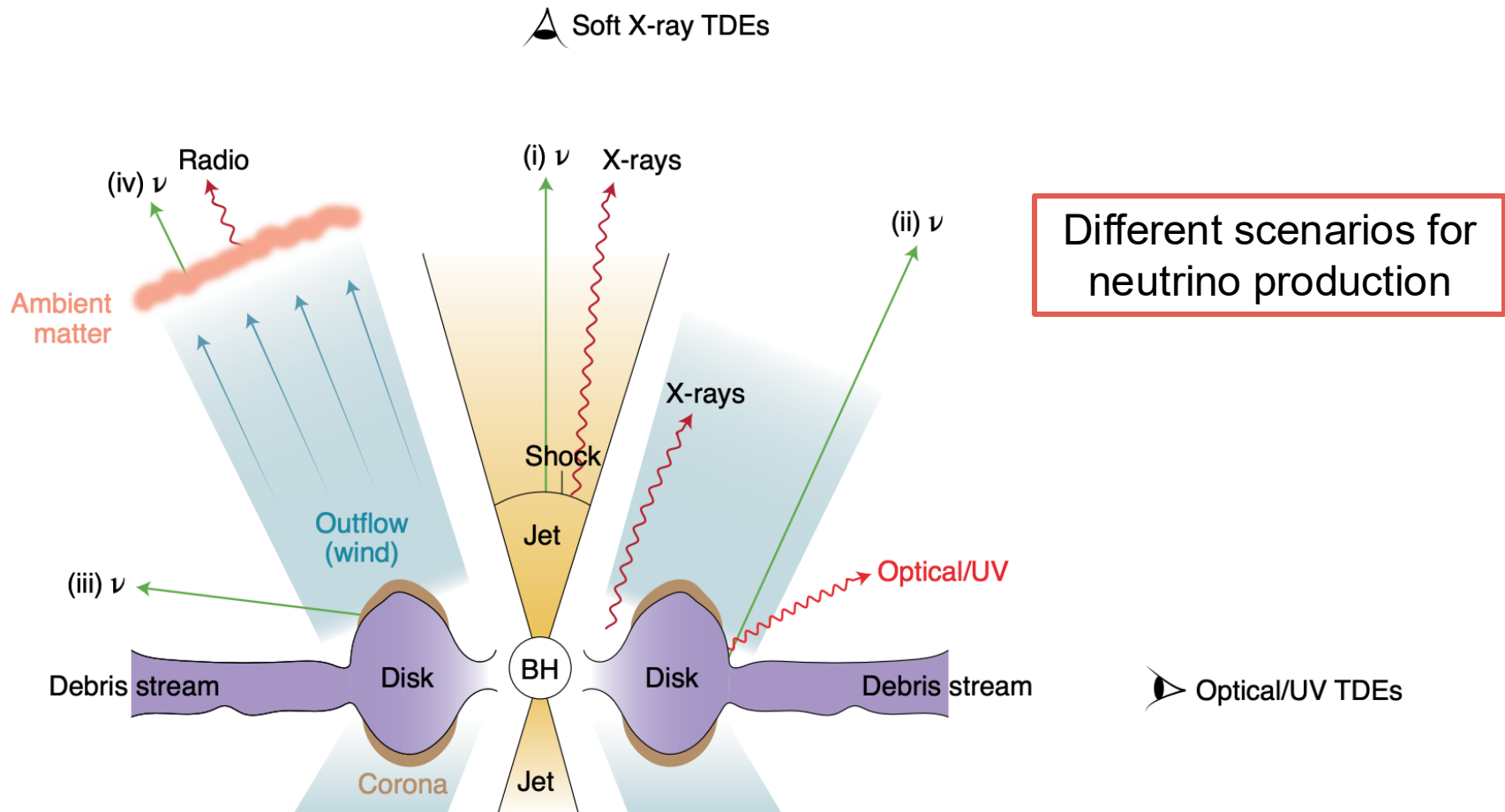


(b) AT2019dsg



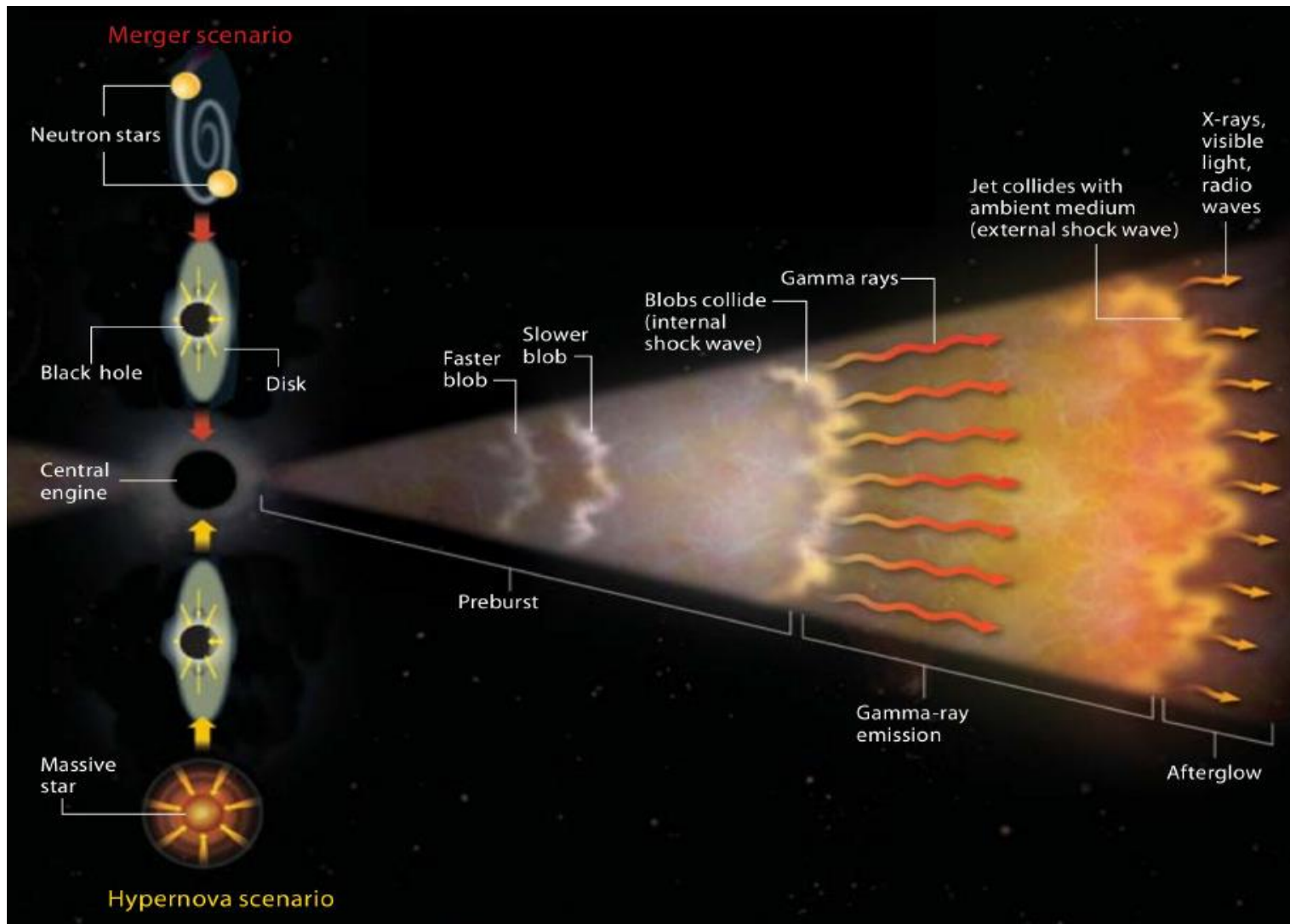
(c) AT2019aalc

Neutrino Production in TDEs

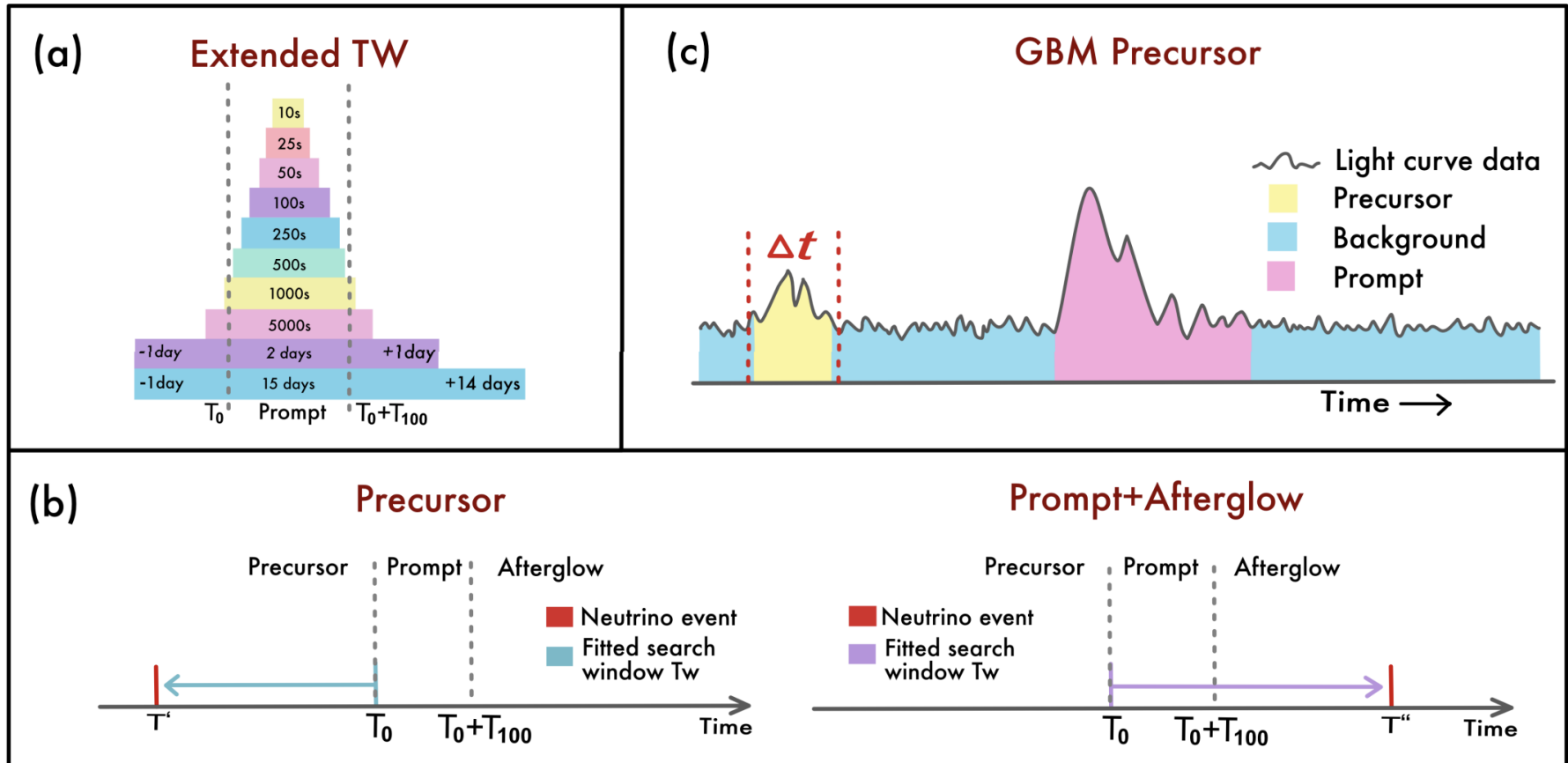


Stellar collapse as neutrino sources?

Gamma-ray Bursts

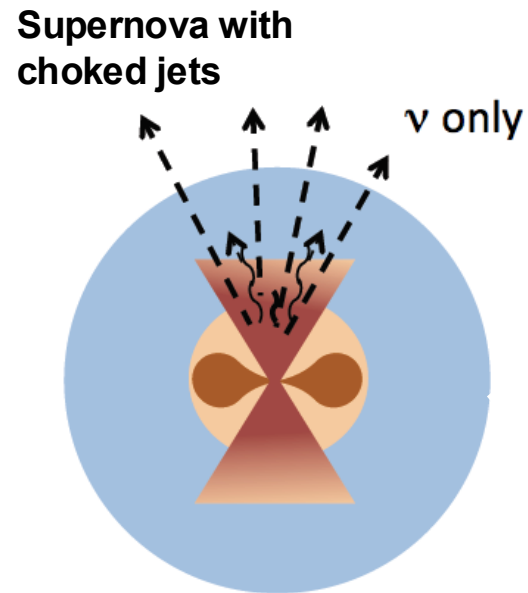
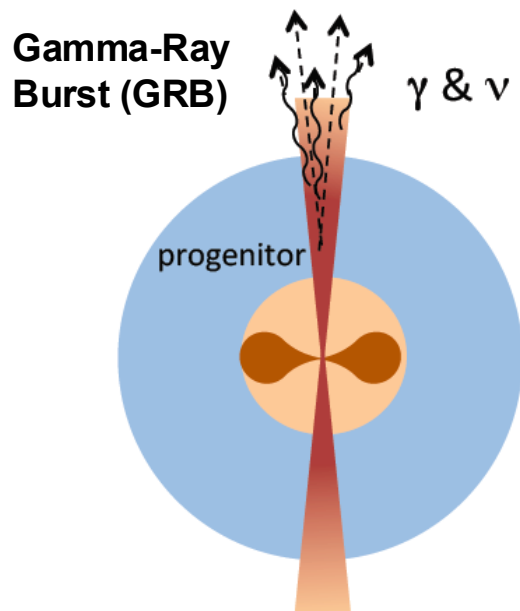


Gamma-ray Bursts: Neutrino Searches

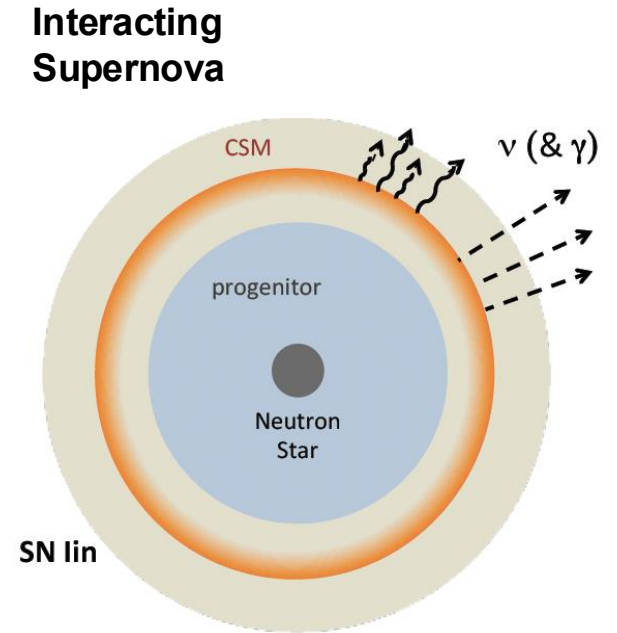
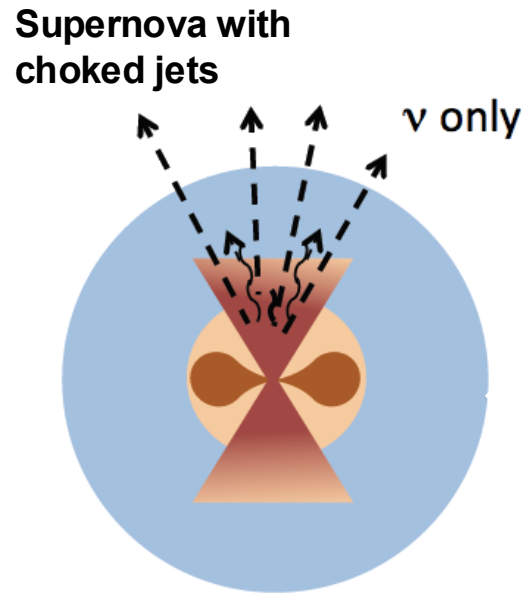
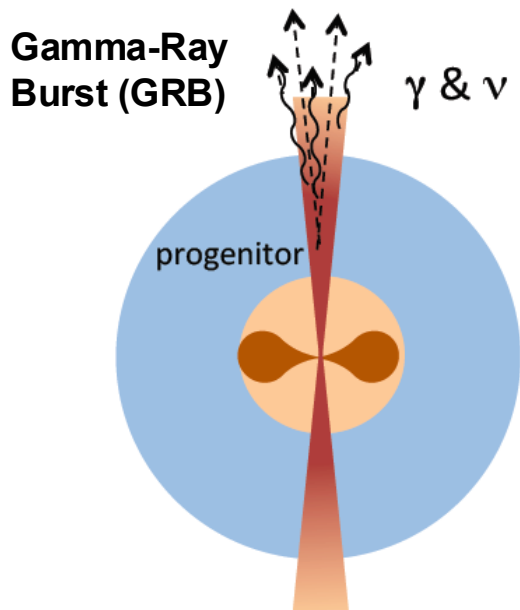


No excess found $\rightarrow$ Prompt neutrinos contribute $<1\%$ to diffuse neutrino flux

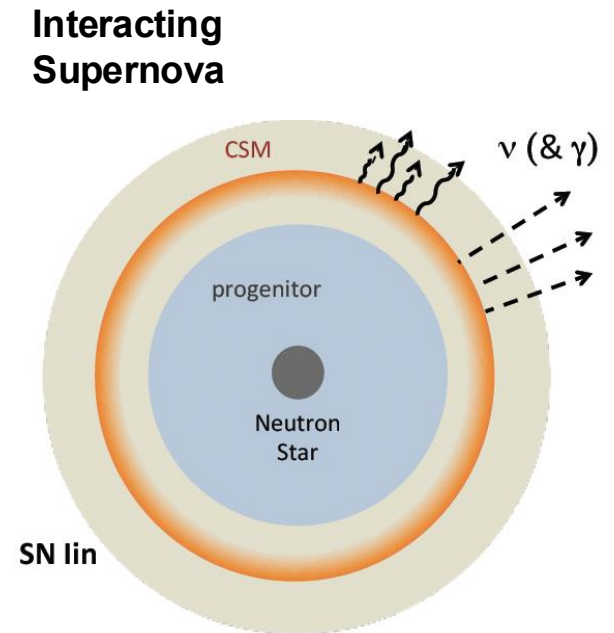
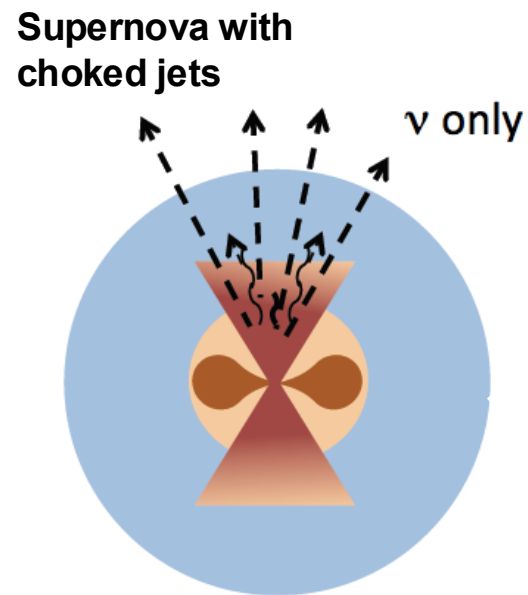
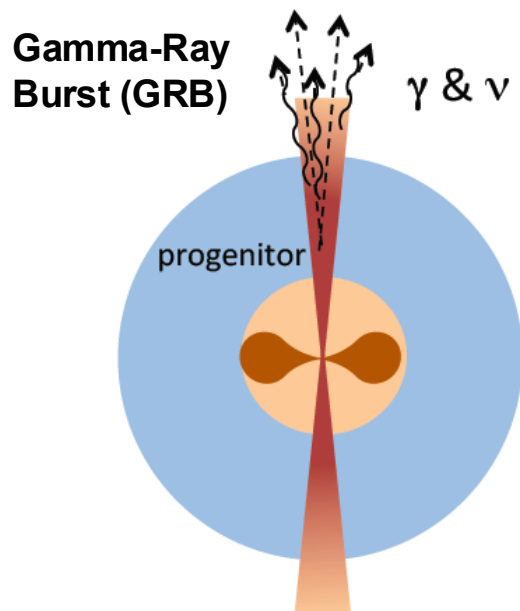
Alternatives to GRBs



Alternatives to GRBs



Alternatives to GRBs



“hidden” sources, best discovered in the optical

Optical Counterpart to a high-energy neutrino: SN2023uqf

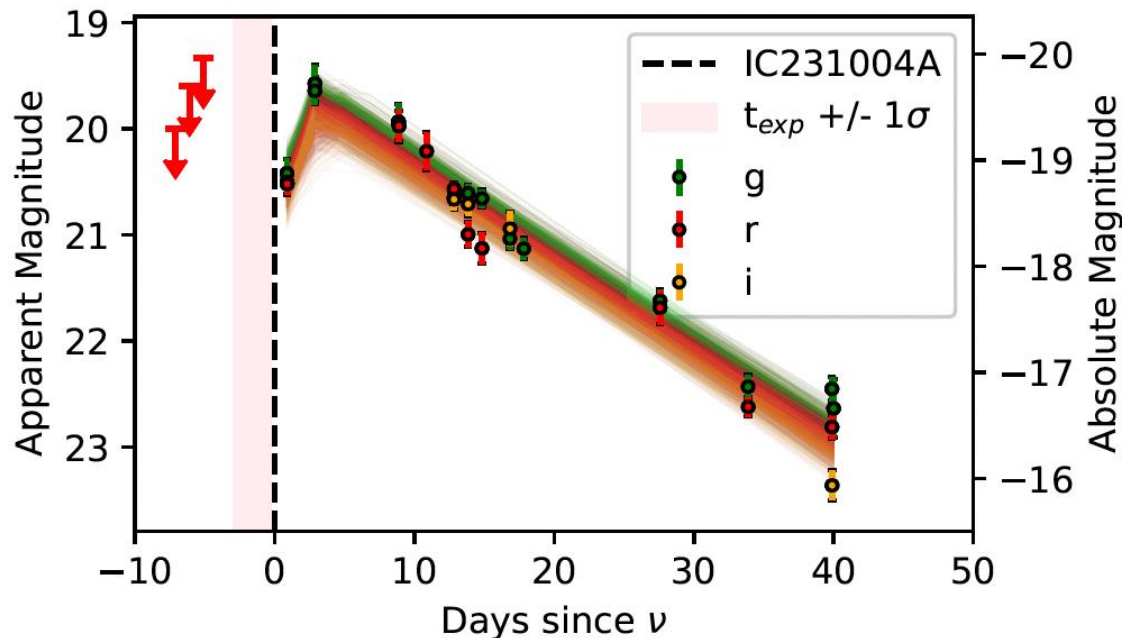
IC-231004A

440 TeV Neutrino

84% signalness

3σ association

Distance: 700 Mpc



Optical Counterpart to a high-energy neutrino: SN2023uqf

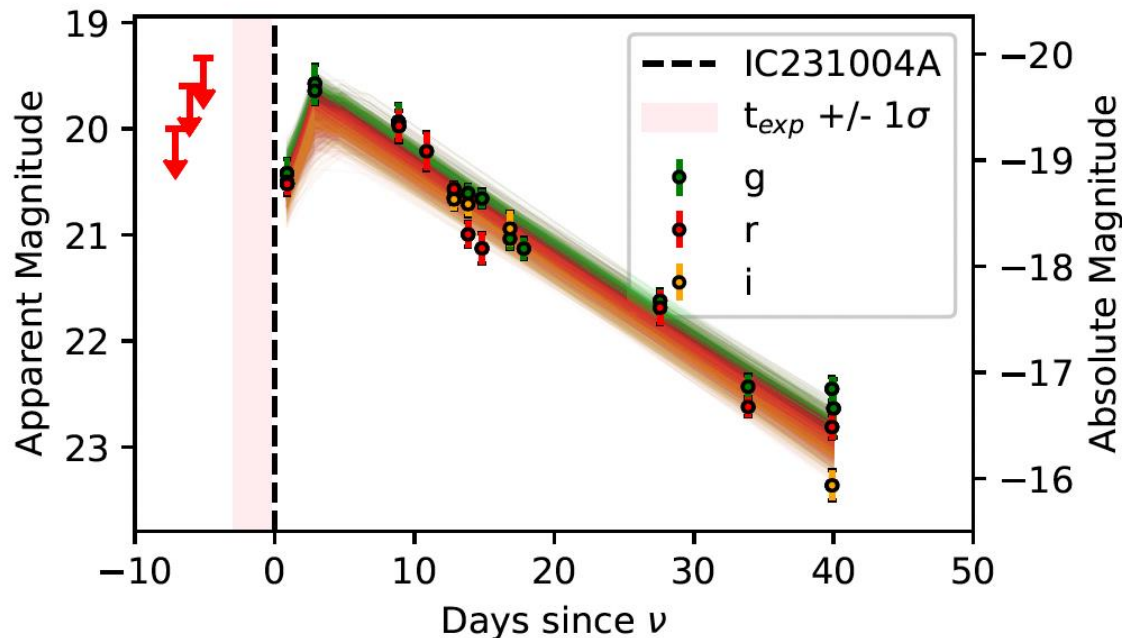
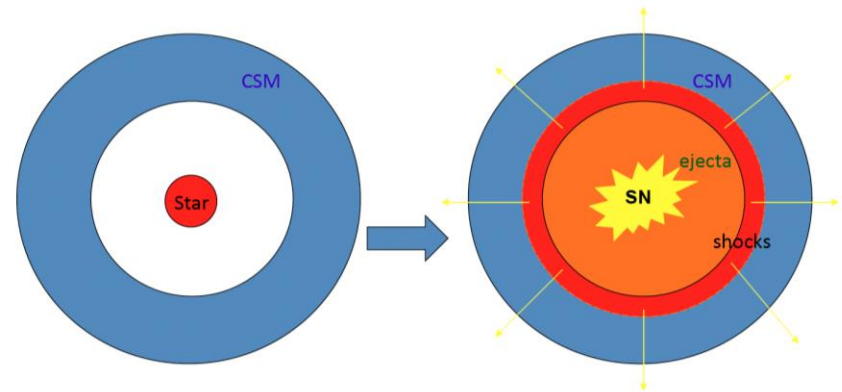
IC-231004A

440 TeV Neutrino

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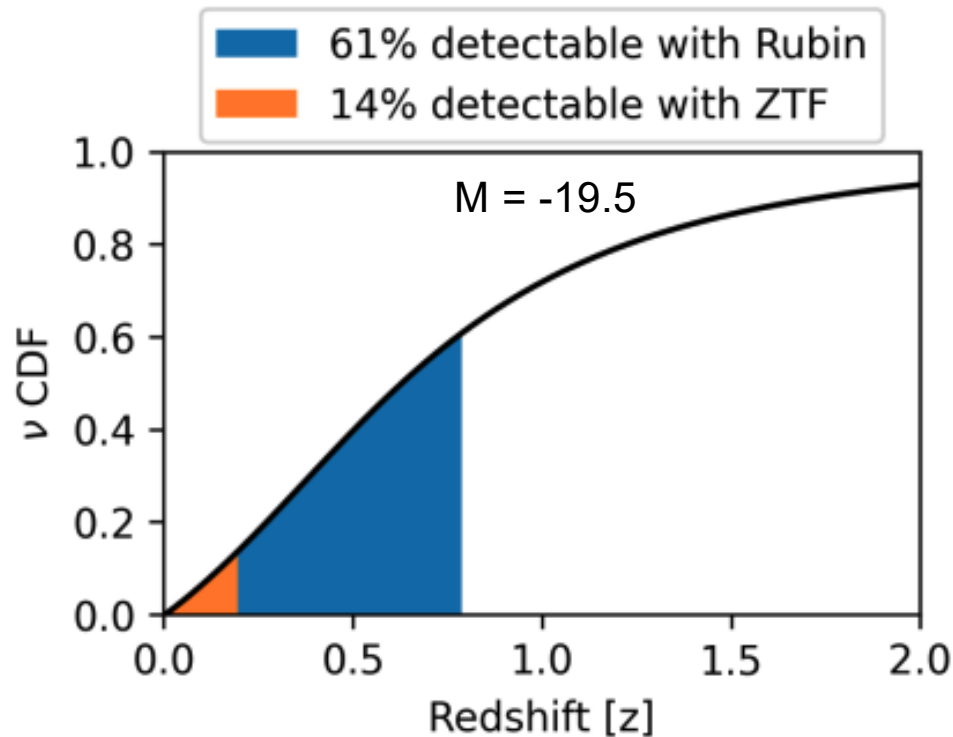
Distance: 700 Mpc



Murase MNRAS 440 (2013)

Evidence for hadronic acceleration in core-collapse supernova explosions

Improved follow-up



First light: July 2025, Alert stream 4-7 months later
Rubin ToO report: 2-12 triggers/year

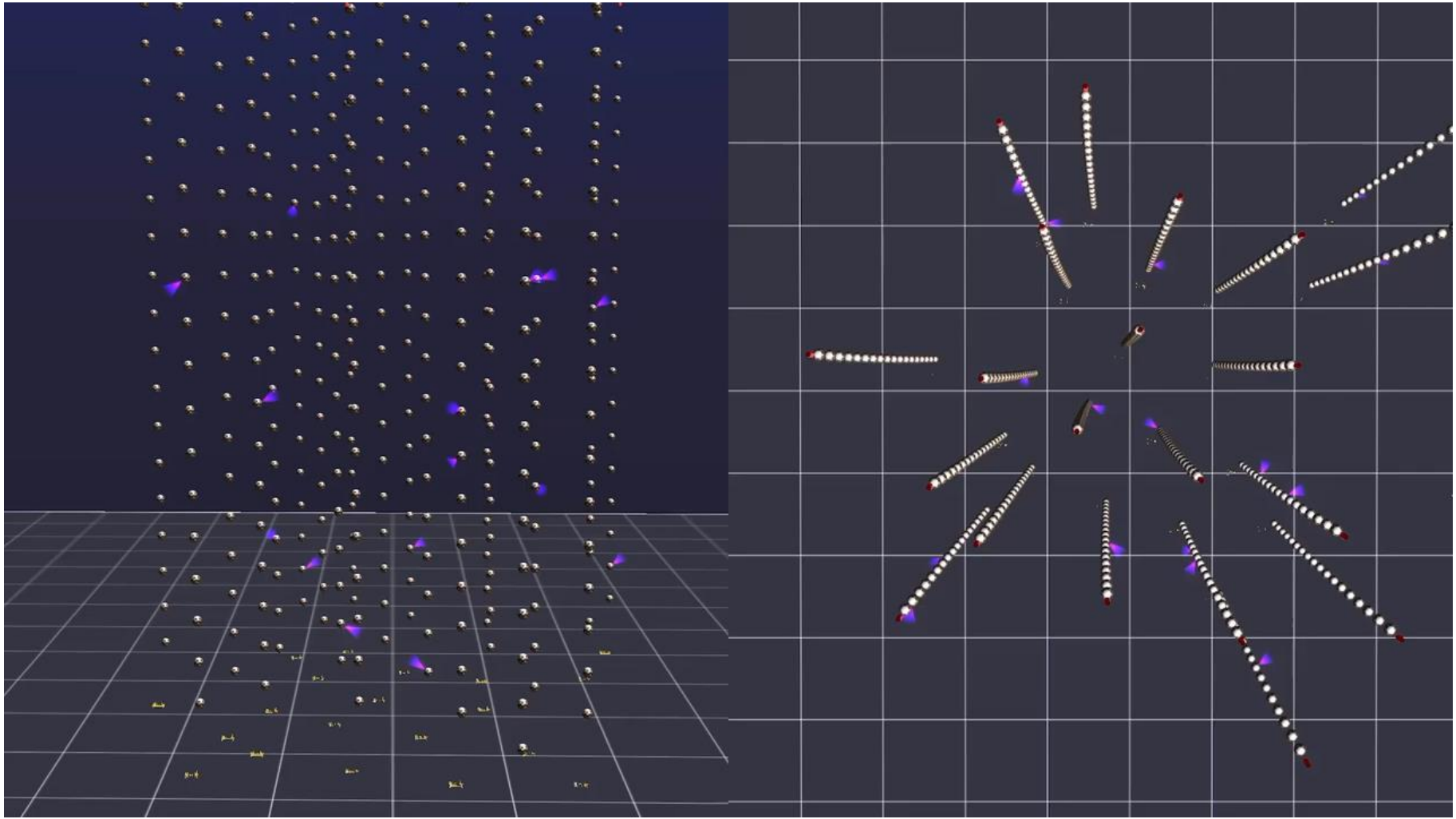


Take home message

While GRBs are disfavored, choked-jet or interacting Supernovae are promising source classes, which we'll be able to probe with up-coming instruments

The highest energy neutrino event ever detected

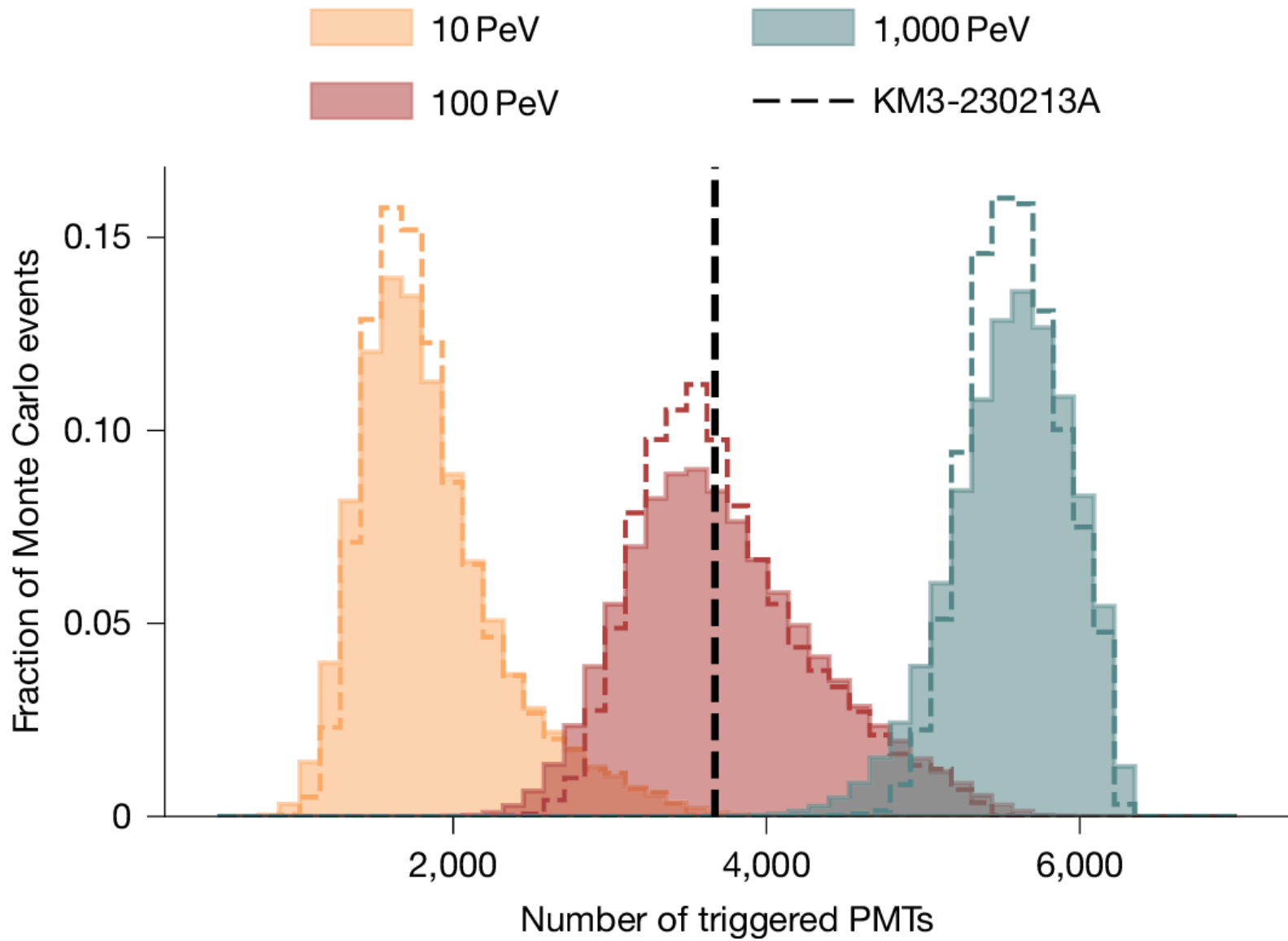
220 PeV Event from KM3NeT



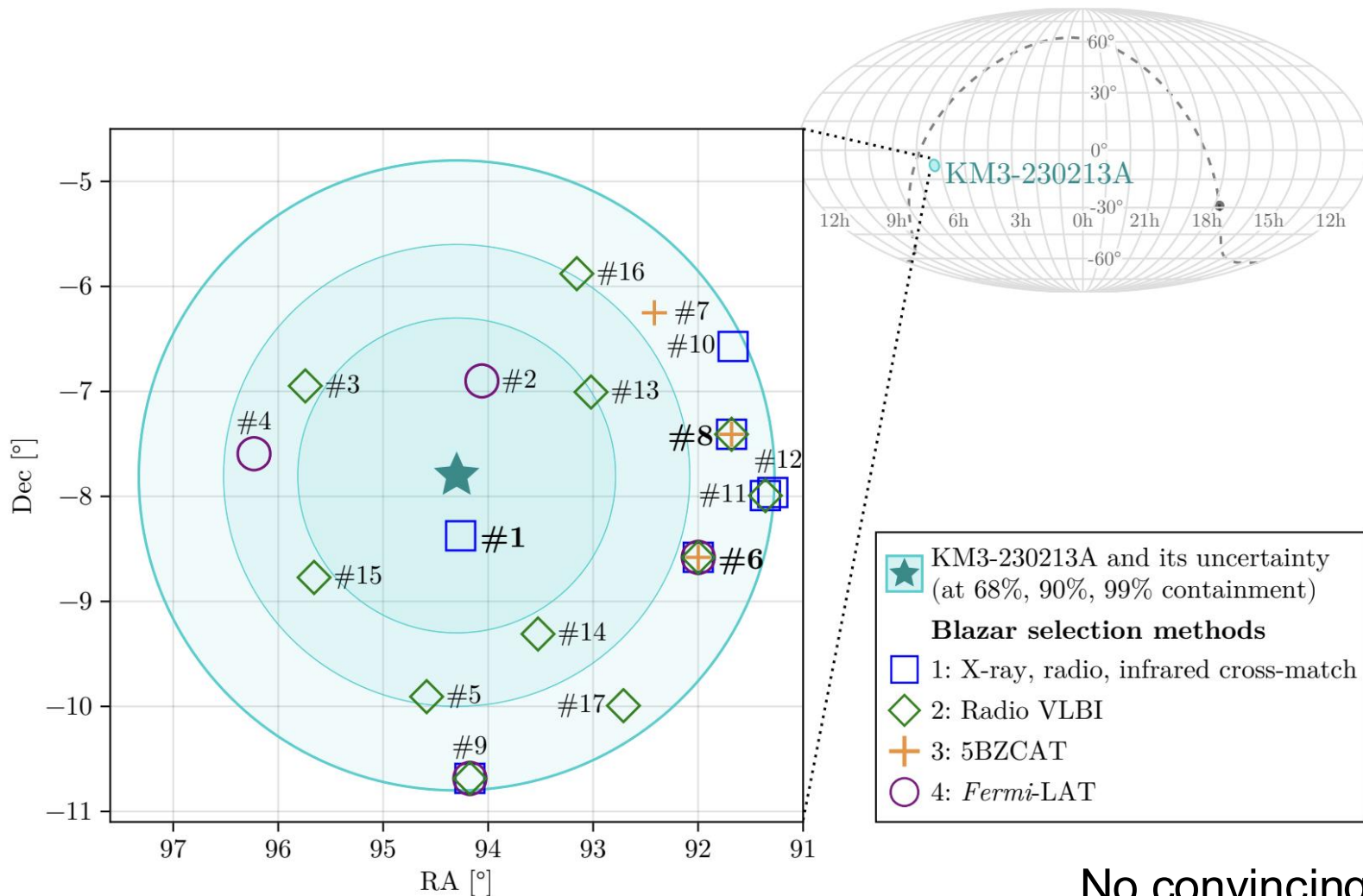
220 PeV Event from KM3NeT

- Huge amount of light seen in the ARCA detector
- Horizontal event (1deg above horizon) as expected since Earth opaque to neutrinos at PeV scale
- 3672 PMTs (35%) were triggered
- Muon simulation indicates that 100 PeV muons are needed to reproduce this signature

Muon energy

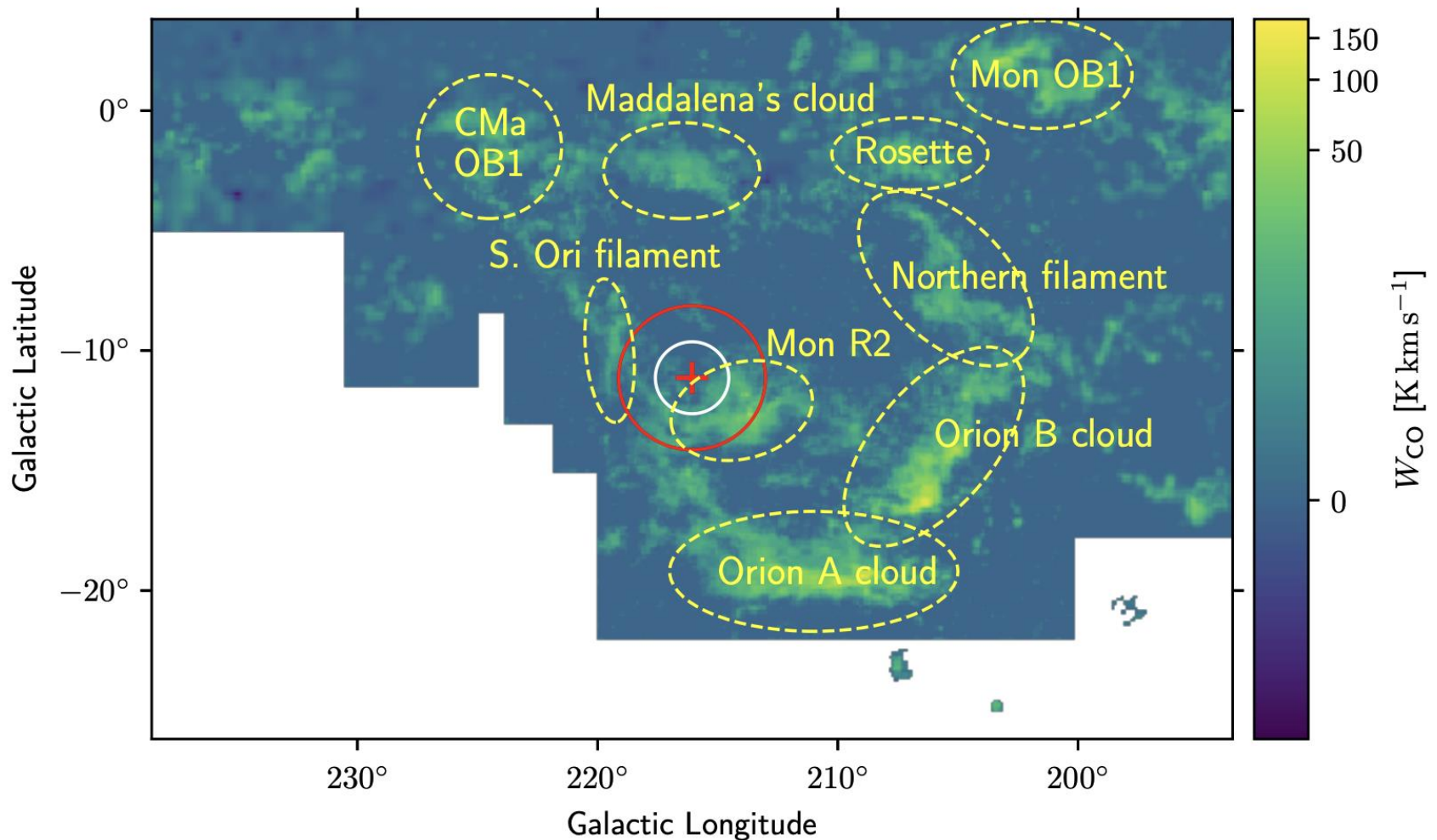


Potential sources?



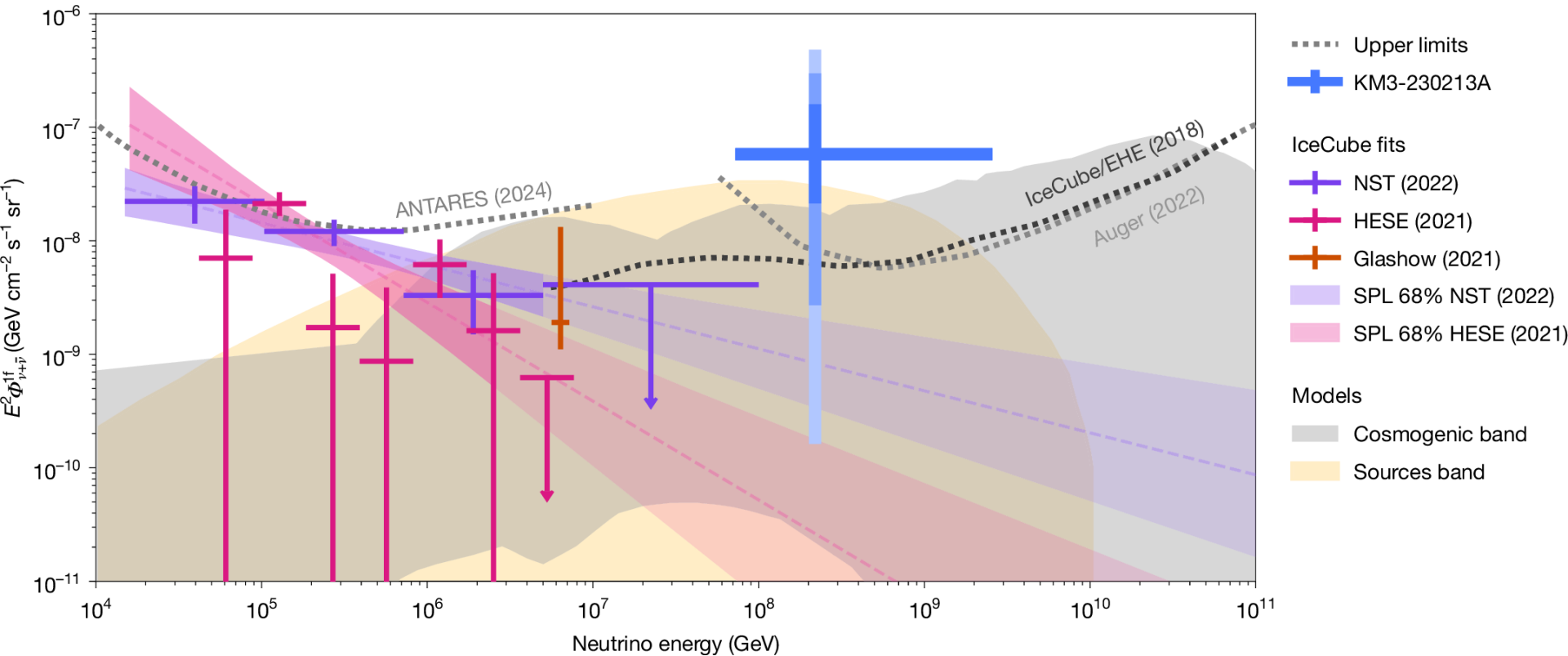
No convincing blazar counterpart found

Potential sources?



Galactic origin unlikely

Tension with IceCube and Auger upper limits?





Take home message

KM3NeT detected the highest energy neutrino event with $\sim 220\text{PeV}$
This is in mild tension with IceCube upper limits.
It's origin is unknown.

Results from water / ice Cherenkov Telescopes

- **Astronomy Summary**

- Diffuse Flux
 - First detected in 2013, now deviation from simple power law
- Milky Way: $\sim 10\%$ of diffuse flux
- Source Candidates
 - Neutrino map alone does not reveal sources
 - Using electromagnetic information
 - Seyfert galaxy NGC 1068 at 4 sigma in 10 year data + other potential Seyfert candidates
 - Blazar TXS 0506+056 at 3 sigma through realtime observations
 - Interacting Supernova SN2023uqf ?

Astro Quiz

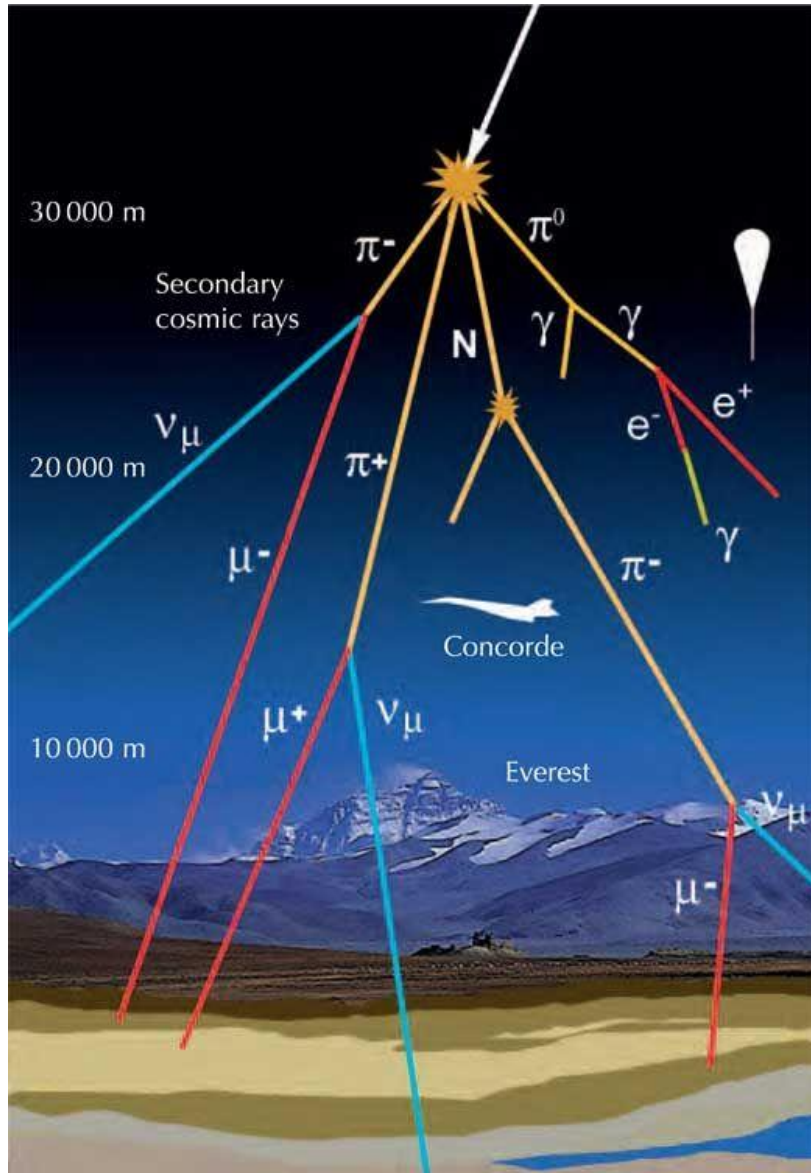


<https://forms.gle/UHKjRUvZPEJZqJqy5>

Results from water / ice Cherenkov Telescopes

- Astronomy
 - Diffuse Flux
 - Milky Way
 - Source Candidates
- **Neutrino as a particle**
 - Neutrino Oscillations
 - Neutrino Cross Section
- New physics
 - Sterile Neutrinos
 - Dark Matter

Atmospheric Neutrinos



What is neutrino oscillation?

- Three neutrino flavor states: one per lepton flavor
- Mixing between these flavor states and the neutrino mass states



- Characterised by the PMNS matrix

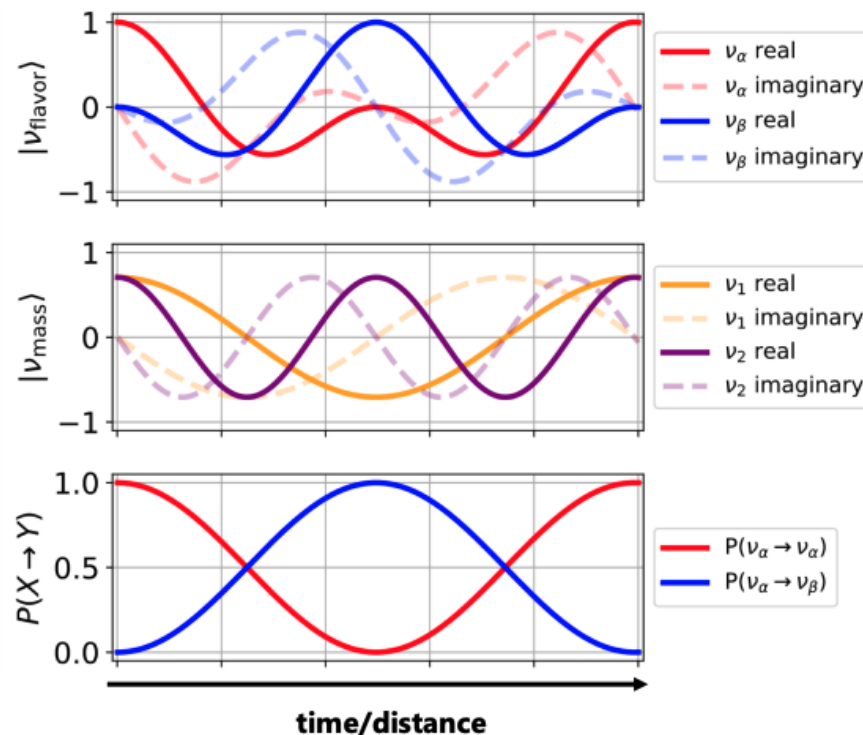
$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

- A neutrino produced as a given flavor is thus a superposition of all three mass states
- The wavefunction of each mass state evolves with a different frequency (defined by its mass) as they propagate
- The superposition of the flavor states therefore changes over time → time-dependent flavour composition
- A neutrino produced as one flavor can therefore be detected later as another → this is neutrino oscillations

A simple example (credit: T. Stuttard)

Visualise the superposition effects and resulting oscillations using simplified model...

- 2 neutrino states (flavour = ν_α, ν_β , mass = ν_1, ν_2) $\rightarrow$ flavor states are 50:50 mix of mass states ($\theta = 45^\circ$)
- $m_2 = \sqrt{2} m_1$



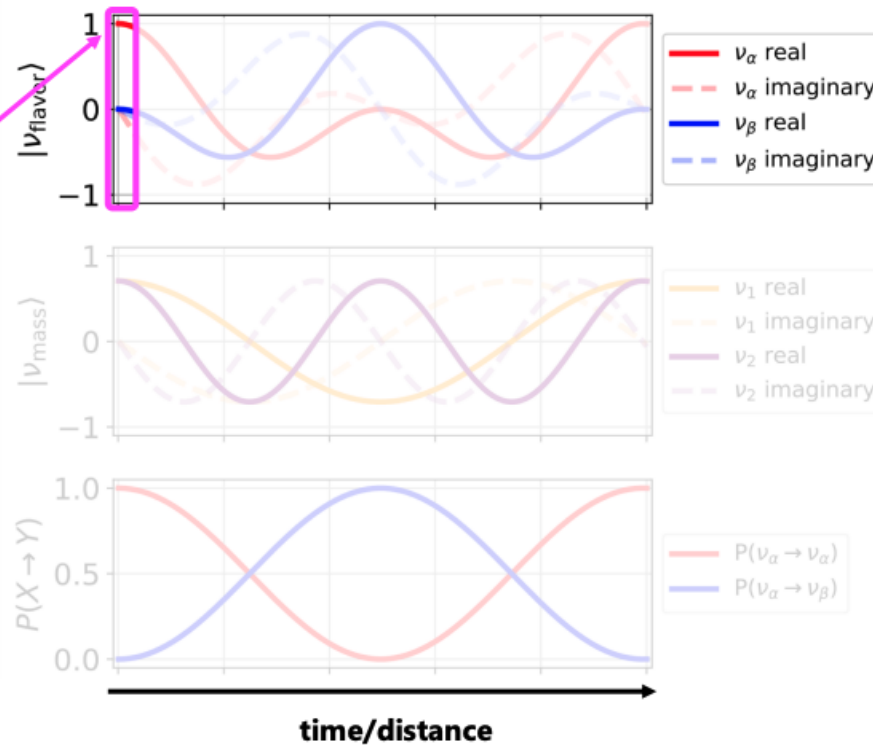
$$|\nu_j(L)\rangle = \exp\left\{-i\frac{m_j^2 L}{2E}\right\} |\nu_j(0)\rangle$$

A simple example (credit: T. Stuttard)

Visualise the superposition effects and resulting oscillations using simplified model...

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- $m_2 = \sqrt{2} m_1$

1) Neutrino produced in pure (single) flavor state
In this case ν_α



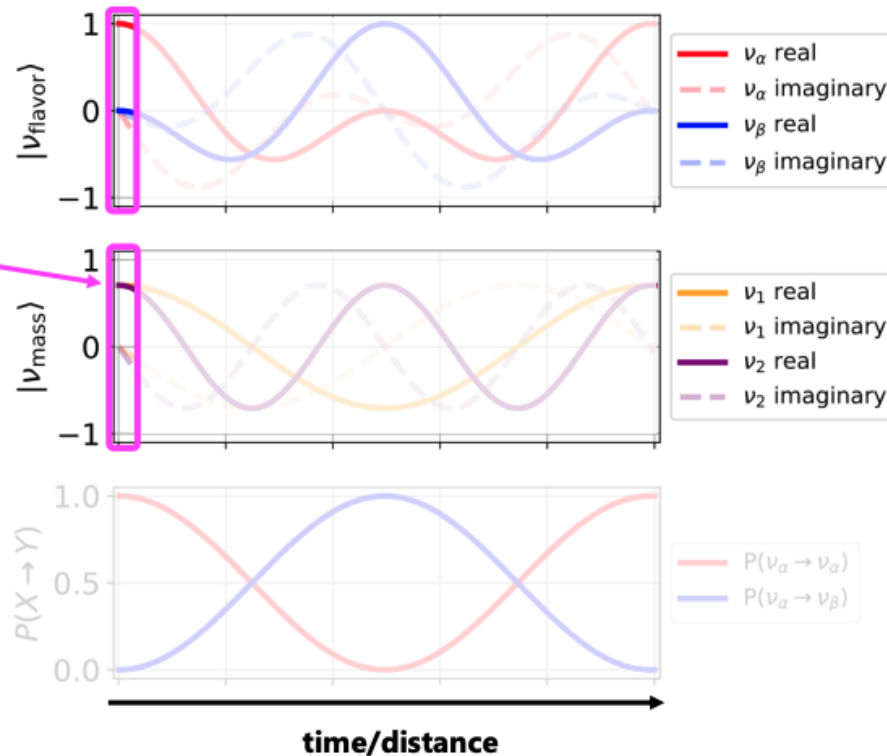
A simple example (credit: T. Stuttard)

Visualise the superposition effects and resulting oscillations using simplified model...

- 2 neutrino states (flavour = ν_α, ν_β , mass = ν_1, ν_2) $\rightarrow$ flavor states are 50:50 mix of mass states ($\theta = 45^\circ$)
- $m_2 = \sqrt{2} m_1$

2) This means the neutrino is initially an equal mix of the two mass states

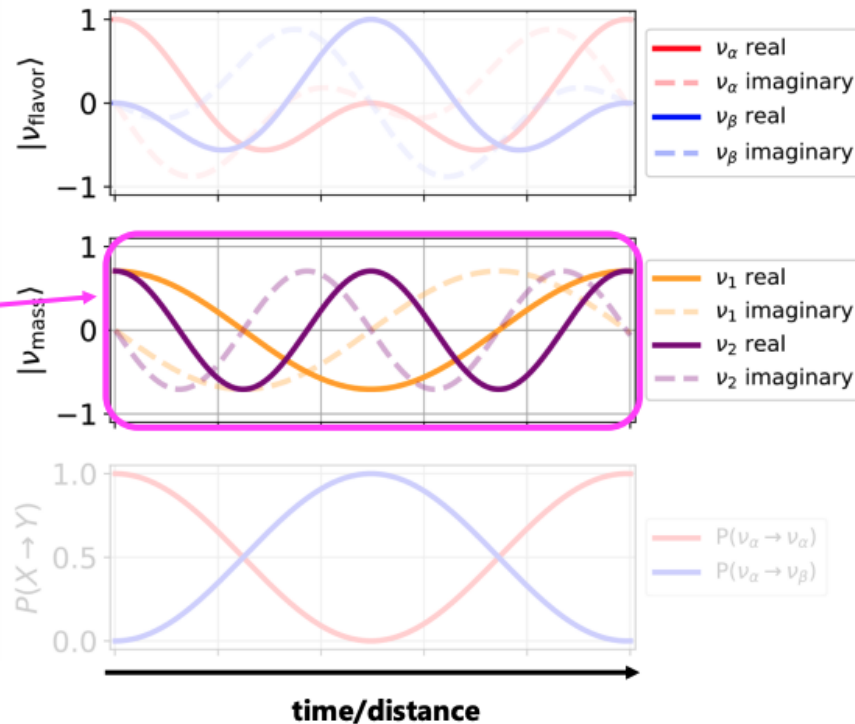
Because of 50:50 mixing in this example



A simple example (credit: T. Stuttard)

Visualise the superposition effects and resulting oscillations using simplified model...

- 2 neutrino states (flavour = ν_α, ν_β , mass = ν_1, ν_2) $\rightarrow$ flavor states are 50:50 mix of mass states ($\theta = 45^\circ$)
- $m_2 = \sqrt{2} m_1$



$$|\nu_j(L)\rangle = \exp\left\{-i \frac{m_j^2 L}{2E}\right\} |\nu_j(0)\rangle$$

Mass state evolution
(Plane wave)

3) The wavefunctions of the mass states evolve over time

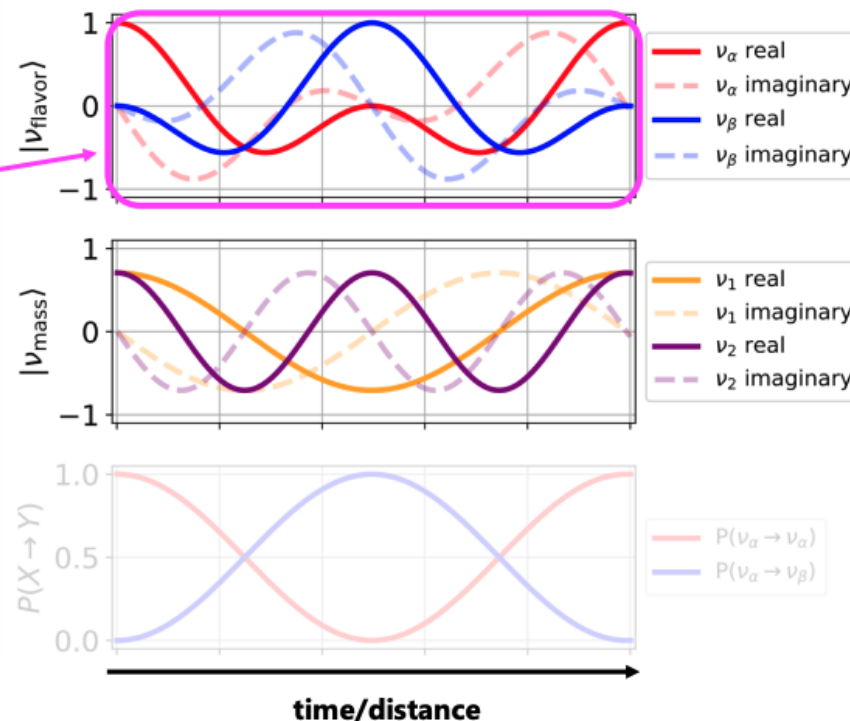
The frequency of this evolution depends on the mass state mass
The frequency of ν_2 is $2 \times \nu_1$ in this example

A simple example (credit: T. Stuttard)

Visualise the superposition effects and resulting oscillations using simplified model...

- 2 neutrino states (flavour = ν_α, ν_β , mass = ν_1, ν_2) $\rightarrow$ flavor states are 50:50 mix of mass states ($\theta = 45^\circ$)
- $m_2 = \sqrt{2} m_1$

4) The evolution of the mass states modified the flavor state superposition over time



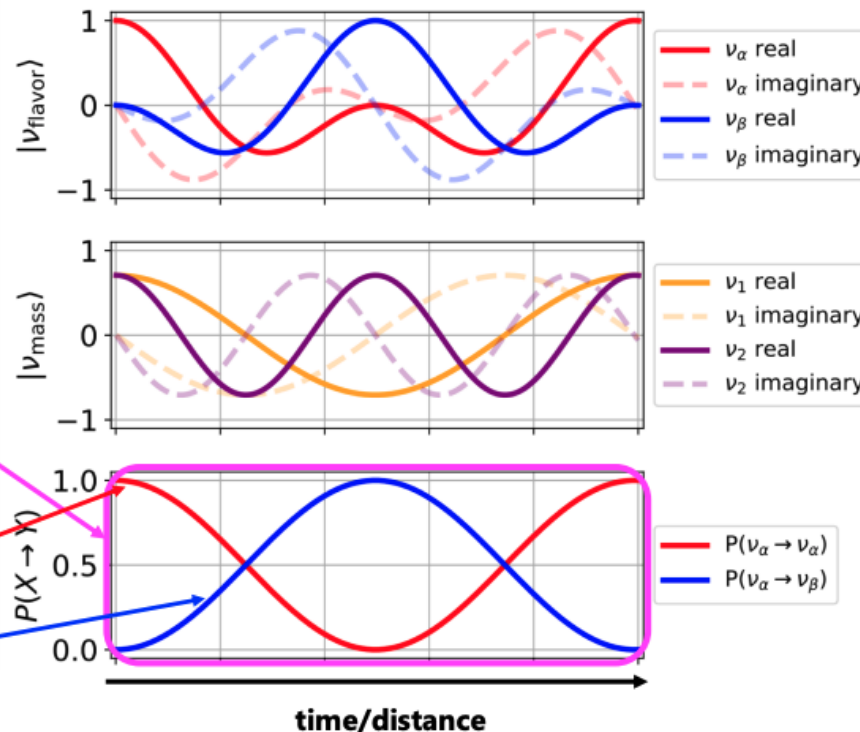
$$|\nu_j(L)\rangle = \exp\left\{-i\frac{m_j^2 L}{2E}\right\} |\nu_j(0)\rangle$$

Mass state evolution
(Plane wave)

A simple example (credit: T. Stuttard)

Visualise the superposition effects and resulting oscillations using simplified model...

- 2 neutrino states (flavour = ν_α, ν_β , mass = ν_1, ν_2) $\rightarrow$ flavor states are 50:50 mix of mass states ($\theta = 45^\circ$)
- $m_2 = \sqrt{2} m_1$



$$|\nu_j(L)\rangle = \exp\left\{-i\frac{m_j^2 L}{2E}\right\} |\nu_j(0)\rangle$$

Mass state evolution
(Plane wave)

5) The probability to be detected as a given flavor thus changes over time

Initially the probability to detect as the production flavor is 100%

Probability to detect as another flavor (*oscillate*) rises and falls over time

Neutrino Oscillations

- Neutrino mixing characterised by the (complex) PMNS matrix

$$\begin{bmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{bmatrix} = \begin{bmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} \end{bmatrix} \begin{bmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{bmatrix}$$

- If unitary, can be parameterised as 3 mixing angles and a CP-violating phase

$$\theta_{12}, \theta_{13}, \theta_{23}, \delta_{CP}$$

- Oscillation frequency depends on the mass-difference between mass states

$$\Delta m_{21}^2, \Delta m_{31}^2$$

→ 6 oscillation parameters to measure in experiments

How can we access those parameters?

- Probability of oscillating between flavors (for simplified 2-flavour case)

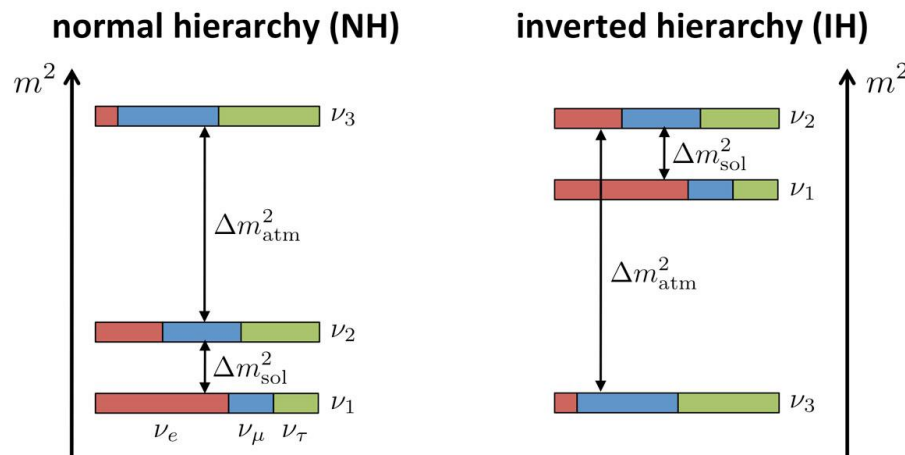
$$P_{\alpha\beta} = \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2 L}{4E}\right)$$

Mixing defines oscillation amplitude

Mass-difference defines frequency

Frequency also depends on L/E ratio for a given neutrino

- Insensitive to the sign of $\Delta m^2 \rightarrow$ neutrino mass ordering problem



Modified oscillation effects for neutrinos passing through matter

Experiments

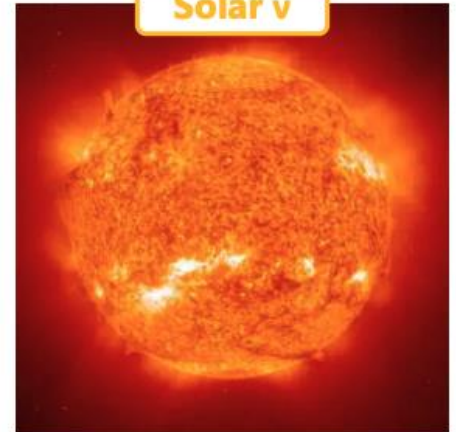
Accelerators



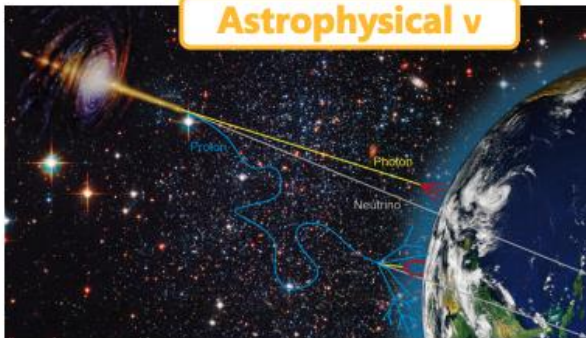
Nuclear reactors



Solar ν



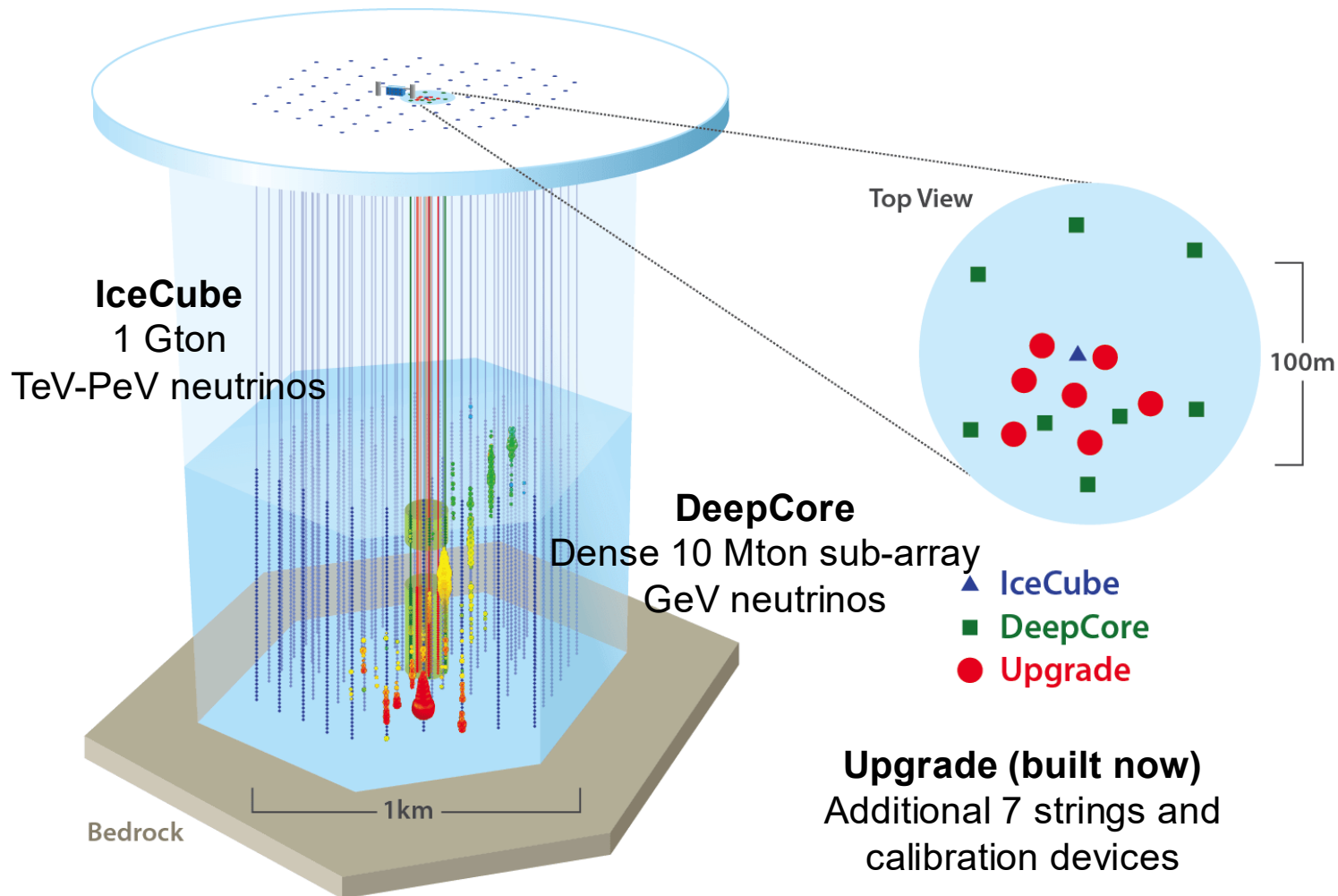
Astrophysical ν



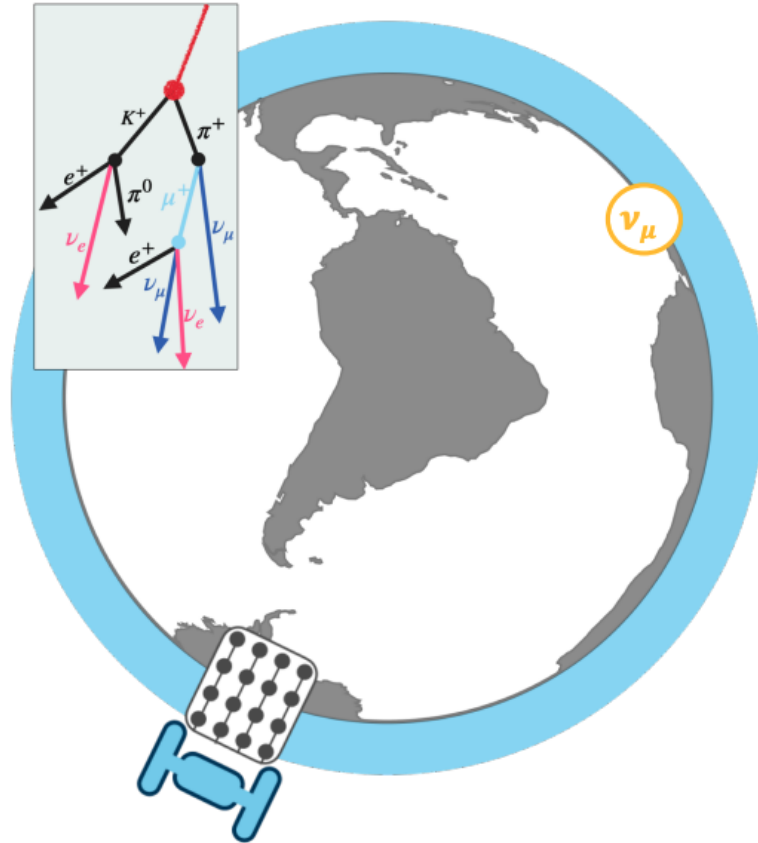
Atmospheric ν



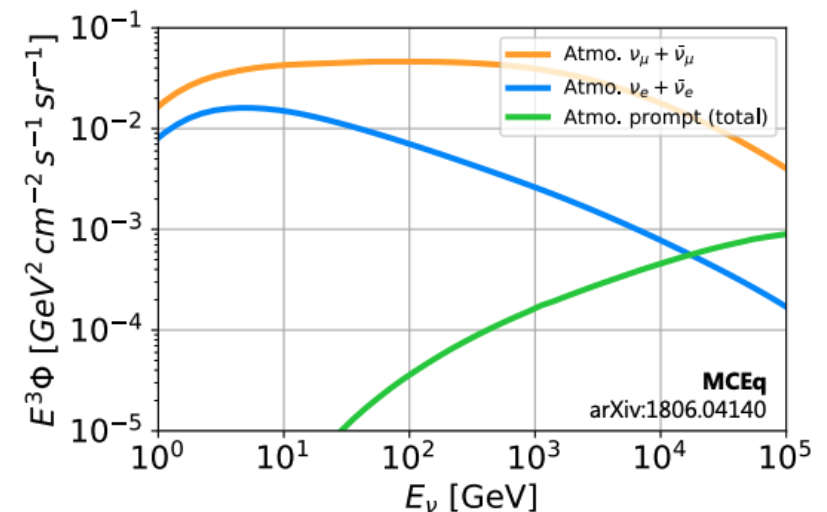
IceCube DeepCore



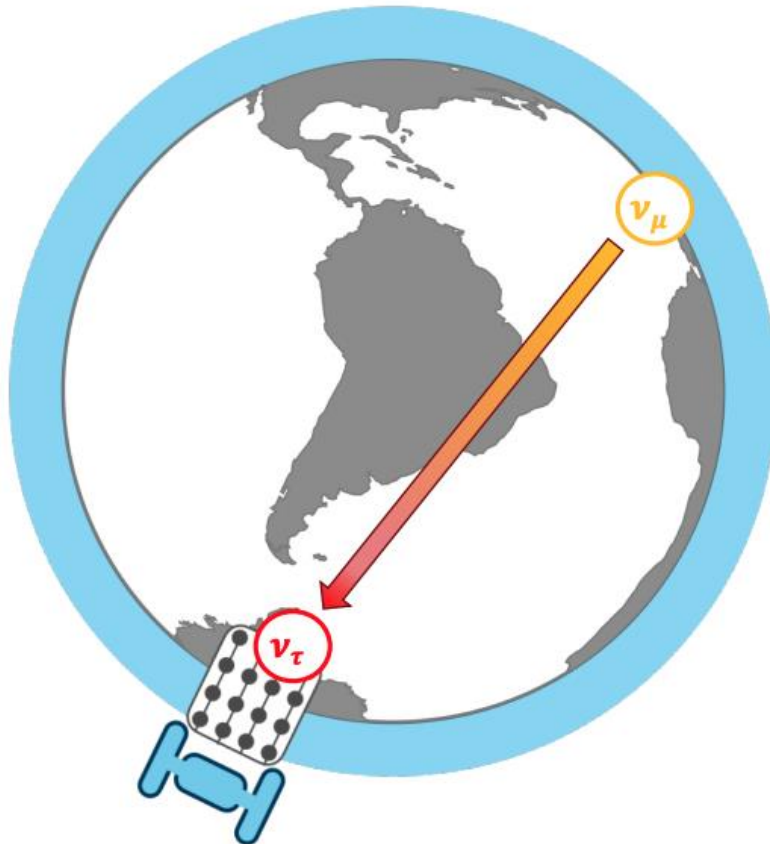
Atmospheric Neutrinos: Background to one, signal to another



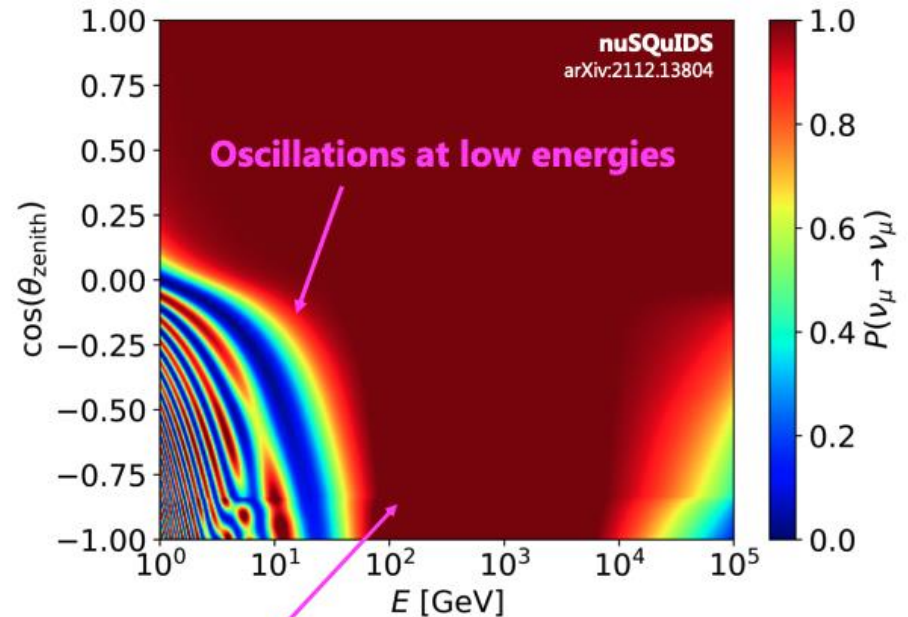
(1) Cosmic rays interact in the atmosphere and produce air showers
→ Large flux of high energy neutrinos



Atmospheric Neutrinos - Oscillations



(2) Neutrinos propagate across the Earth

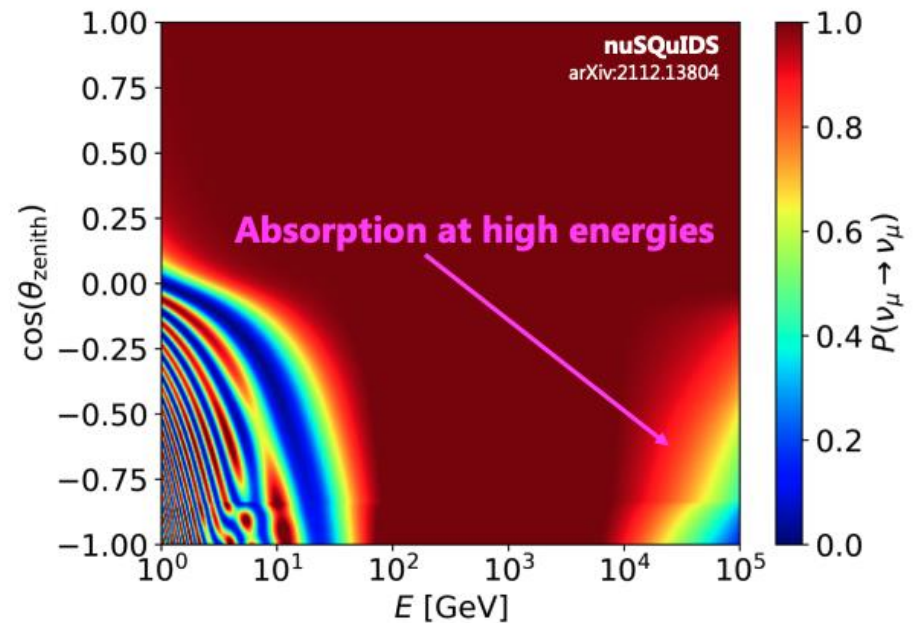


When E is $\gtrsim 100$ GeV, oscillation baseline is larger than the Earth's diameter

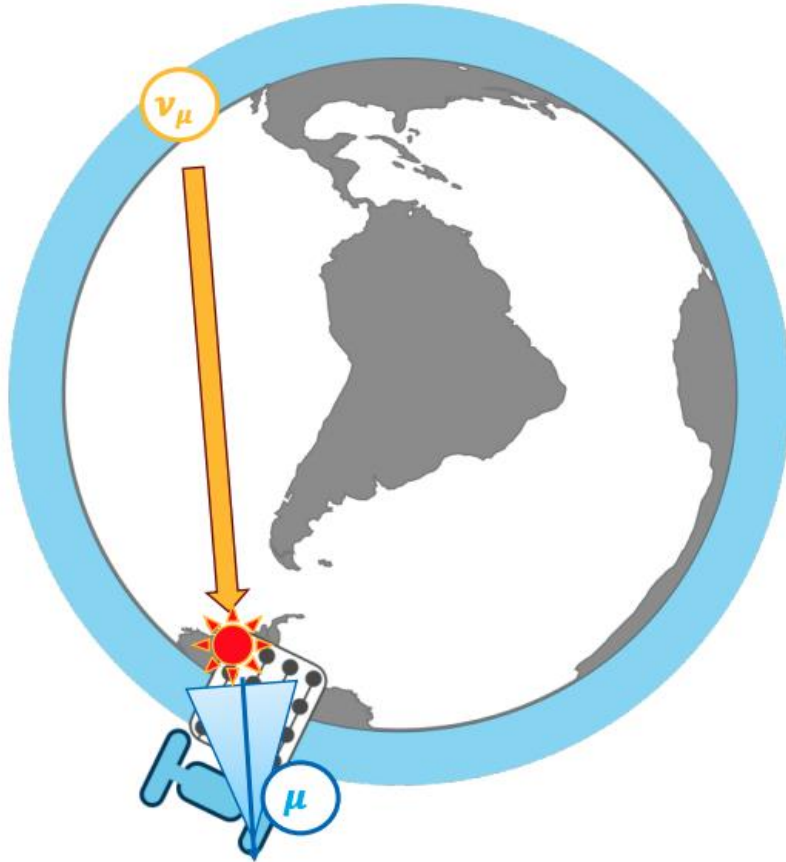
Atmospheric Neutrinos - Oscillations



(2) Neutrinos propagate across the Earth



Atmospheric Neutrinos - Oscillations

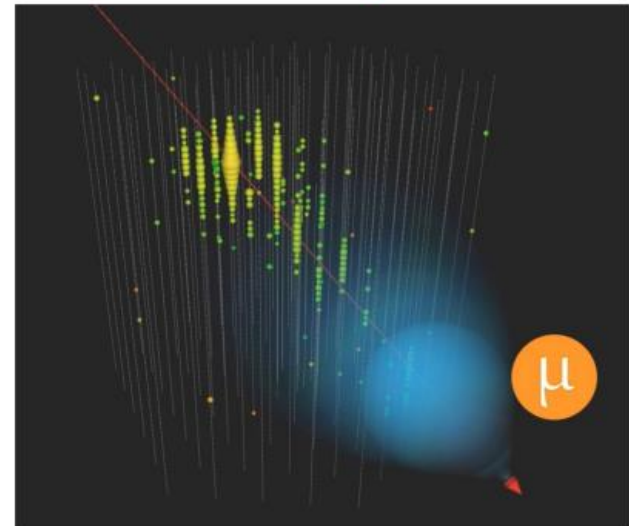


(3) Detection via Cherenkov emission from products of $\nu - N$ interactions

Predominantly Deep Inelastic Scattering (DIS)

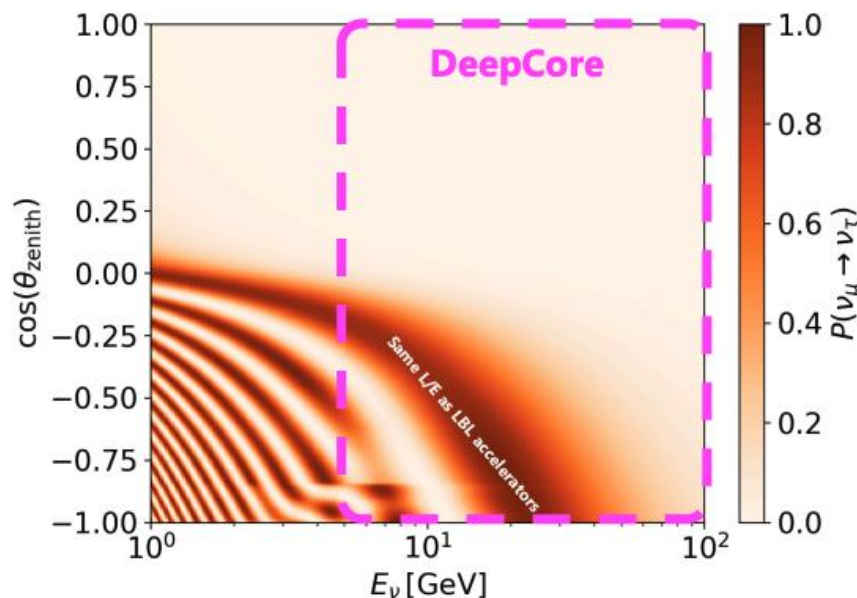
“Tracks” from secondary μ

“Cascades” from secondary e , τ and hadrons



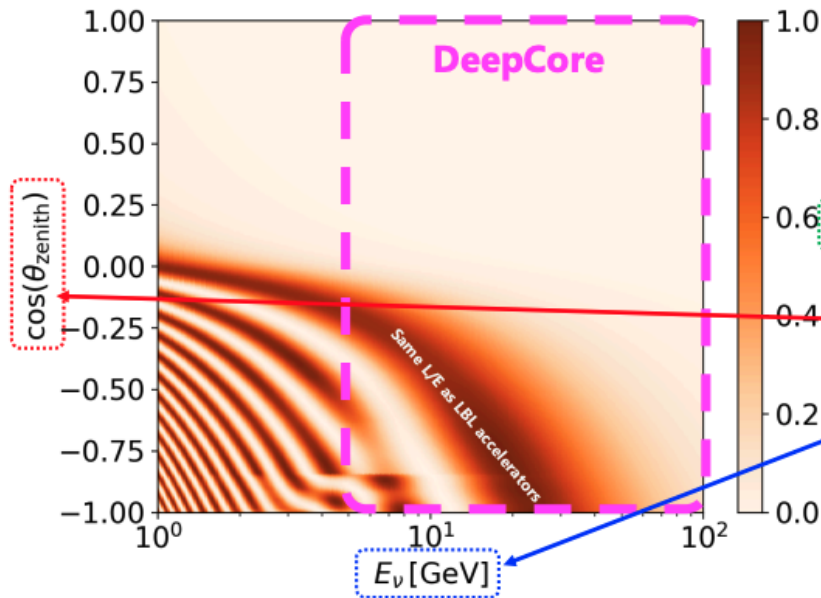
Measuring Neutrino Oscillations

- The DeepCore sub-array of IceCube can measure atmospheric neutrino oscillations in the 5 – 100 GeV energy range
- Earth-crossing ν_μ **near maximally oscillate to ν_τ** $\rightarrow$ measures θ_{23} and Δm_{32}^2



$$P_{\alpha\beta} = \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2 L}{4E}\right)$$

Measuring Neutrino Oscillations



To measure these effects, for each neutrino must reconstruct:

Flavor – also known as Particle ID (PID)

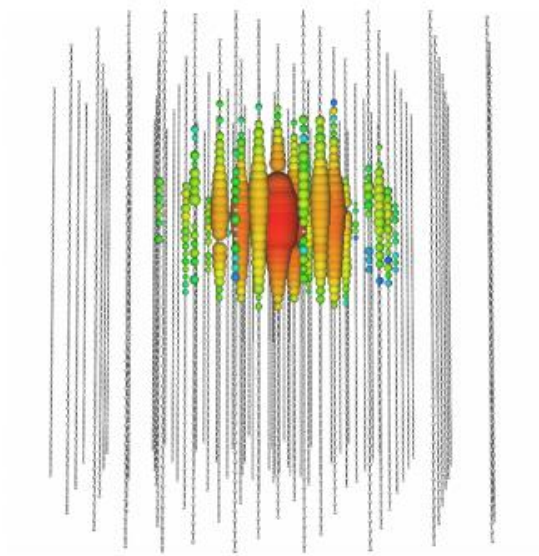
Baseline - determined from arrival (zenith) angle

Energy

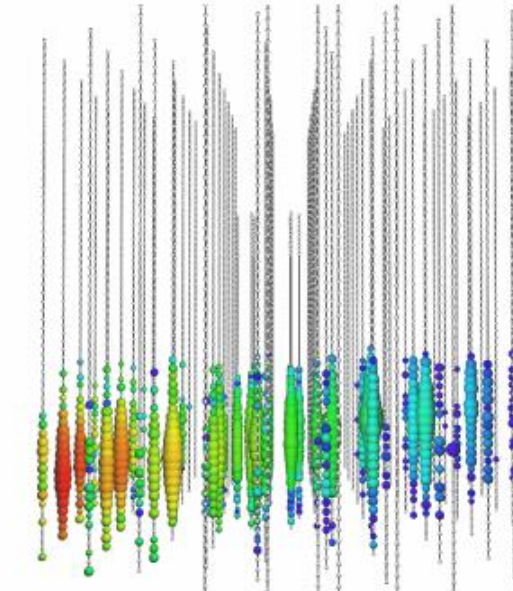
$$P_{\alpha\beta} = \sin^2(2\theta) \sin^2\left(\frac{\Delta m^2 L}{4E}\right)$$

Measuring Neutrino Oscillations

- Two event topologies at high energies



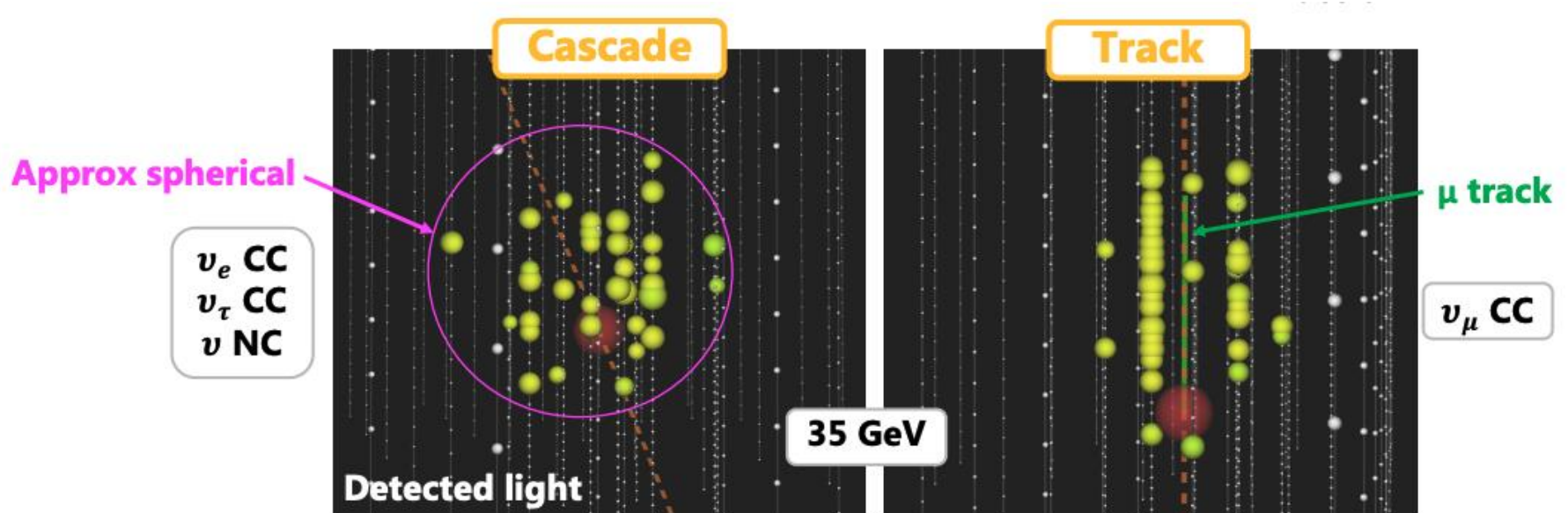
Cascade
 $E \sim 1 \text{ PeV}$



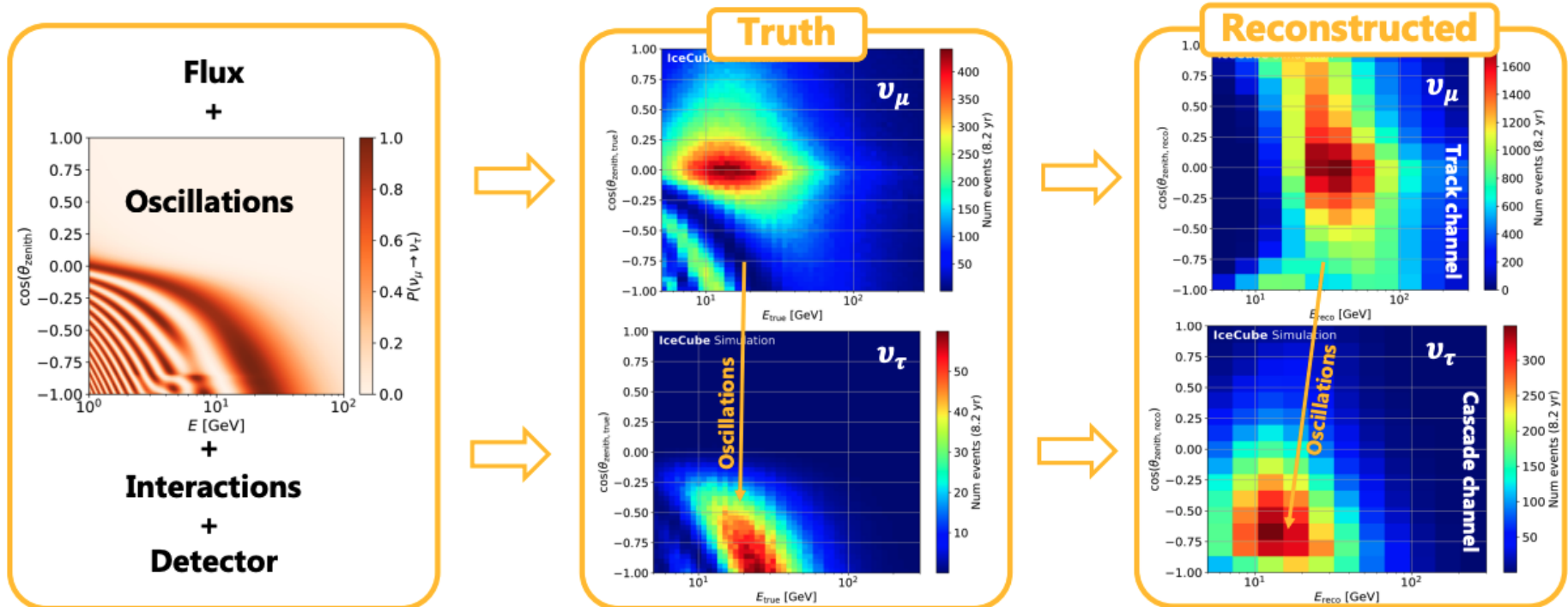
Muon track
 $E \sim 140 \text{ TeV}$

Measuring Neutrino Oscillations

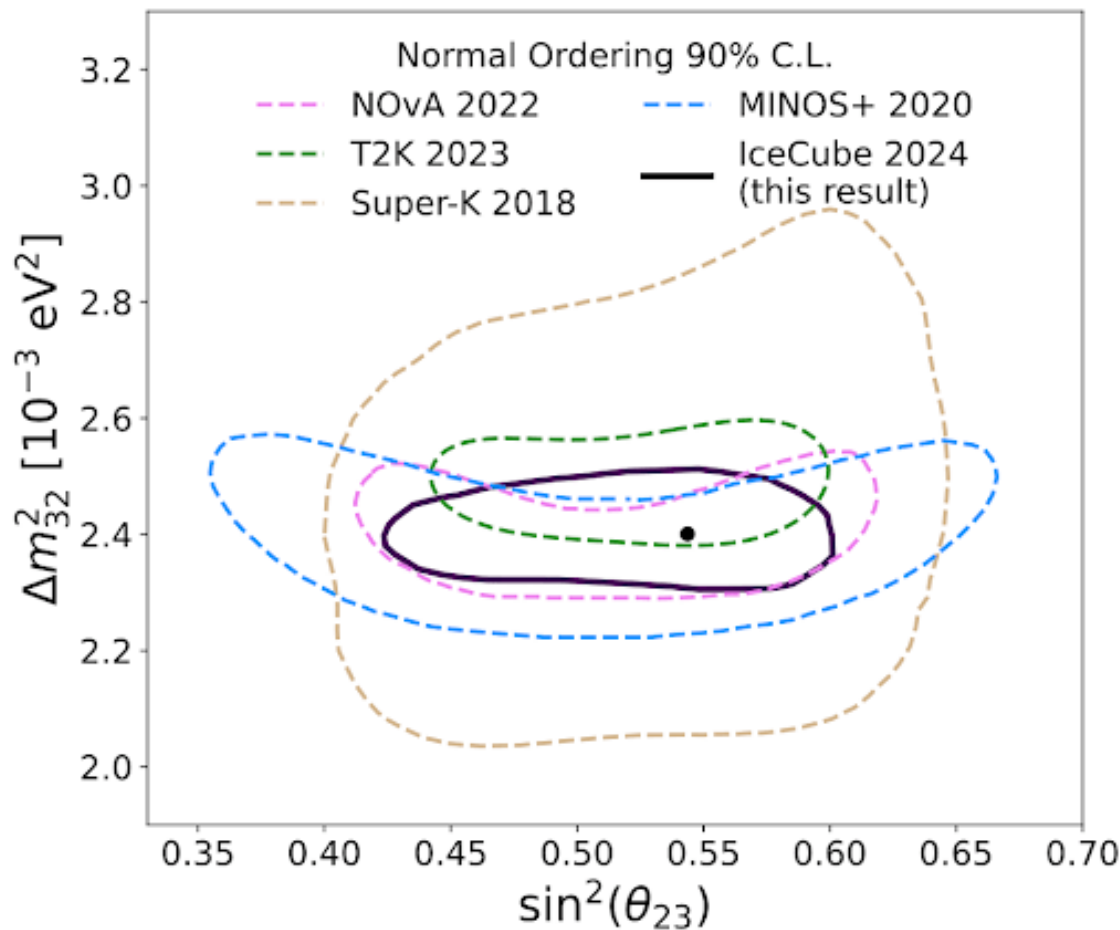
- Two event topologies at oscillation energies:



Measuring Neutrino Oscillations



Neutrino Oscillations: Results



DeepCore data taken
between 2012-2021

150,000 candidate
neutrino events between
5 GeV and 100 GeV



Take home message

Neutrino telescopes can do competitive measurements of neutrino oscillation parameters. The $O(10\text{GeV})$ range is crucial.